### **Lecture Series Buenos Aires 18-3-2024 until 22-3-2024**

Lecture M8 – Attosecond solid state physics



Marc Vrakking

[marc.vrakking@mbi-berlin.de](mailto:marc.vrakking@mbi-berlin.de)

Max-Born-Institut

## Contents

- High-harmonic generation in solid samples
- Attosecond and femtosecond time-resolved pump-probe spectroscopy in solid samples
	- Strong field-driven electron dynamics
	- Perturbative electron dynamics

## **Discovery of HHG in solids**



High-order harmonic generation in ZnO crystals: (left) measured spectra for two pulse energies; (right) high-energy cutoff as a function of the peak field, along with a linear fit. The intercept along the energy axis is near the bandgap (3.2 eV)

**Ghimire, Nat. Phys. 7, 138 (2011)**

### **Mechanism: intraband vs interband**

#### **Atomic HHG**



**Ghimire, Nat. Phys. 7, 138 (2011)**

### Attosecond and femtosecond time-resolved pump-probe spectroscopy in solid samples

Strong field-driven electron dynamics

## **Strong field interactions: atoms vs. solids**



- When  $\gamma$ <<1, the ionization proceeds by adiabatic tunneling – horizontal process
- When  $\gamma$ >>1, the ionization proceeds by multiphoton absorption – vertical process
- When  $\gamma \approx 1$ , non-adiabatic tunneling

#### **M. Y. Ivanov et al. J. Modern Optics 52, 165 (2005)**

#### **Atom in strong field Periodic system in strong field**



- When γ<<1, the ionization *intraband* motion leads to adiabatic tunneling and transfer from the valence (VB) band to the conduction (CB) band
- When γ>>1, (multi)-photon *interband* transitions from the valence (VB) band to the conduction (CB) band

#### **S. Y. Kruchinin, Rev. Mod. Phys. 90, 021002 (2018)**

### **First: an attosecond experiment without attosecond pulses**



Excitation of a  $SiO<sub>2</sub>$  sample (dielectric, negligible conductivity) with an intense, few-cycle near-IR laser

Measurement, using two electrodes, of a photo-induced current

N.B. Bandgap 9 eV, photon energy 1.7 eV  $\rightarrow$  falls within adiabatic (tunneling) picture

**A. Schiffrin et al., Nature 493, 70 (2013)**

### **First: an attosecond experiment without attosecond pulses**



Observation of a current when the laser polarization is perpendicular to the electrodes, with the direction controlled by the **CEP** 

Laser electric field dependence of the current, revealing the high non-linearity involved in the process

#### **A. Schiffrin et al., Nature 493, 70 (2013)**



Two-pulse experiment with a strong pulse (2 V/cm) polarized parallel to the electrodes creating the conduction band population and a weak pulse (0.2 V/cm) driving the current + comparison of the field of the latter pulse to attosecond streaking

### **Next: using attosecond transient absorption to probe the injection process**



**M. Schultze et al., Nature 493, 75 (2013)**

streaking (in Ne) and attosecond transient absorption (in  $SiO<sub>2</sub>$ )

### **Next: using attosecond transient absorption to probe the injection process**

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IR driver pulse (determined by streaking experiment

Experimental transient absorption into conduction band (blue) and comparison to theory (red)

Transient lineshifts in the experiment (blue) in theory (red)

Evidence of a reversible, nearinstantaneus response – production of virtual carriers in the CB by field-driven tunneling

### Attosecond and femtosecond time-resolved pump-probe spectroscopy in solid samples

Perturbative electron dynamics



IR/VIS excitation from the valence (VB) to the conduction (CB) band

What are the relaxation times and the relaxation mechanisms of the CB electrons and VB holes that are produced?

Approach: excited  $3d_{3/2}$  and  $3d_{5/2}$  core electrons into available states (i.e. 4p character)

**M. Zürch et al., Nat. Comm. 8, 15734 (2017)**





#### **M. Zürch et al., Nat. Comm. 8, 15734 (2017)**



#### **Hole dynamics (lower photon energies)**

- Separate the measured response into two parts using singular value decomposition (SVD)
- Slower component (~1 ps): electron-hole recombination
- Faster component (~140 fs): scattering between different hole states
- Additional fast component (~170 fs): hole relaxation

#### **M. Zürch et al., Nat. Comm. 8, 15734 (2017)**

### Correlated electronic/nuclear motion in LiBH<sub>A</sub>



 $\cdot\,10^4$ 

Steady-state XUV absorption around Li K-edge Transient XUV absorption around Li K-edge



*Time-resolved X-ray diffraction and XUV absorption reveal correlated electronic and nuclear dynamics!!*

B A  $A \cdot 10^4$  $100$ 150 200 250  $-100 - 50$  $\overline{0}$  $50$  $-100 - 50$ 100  $150^{-1}$  $\frac{200}{ }$  $\overline{250}$ Delay [fs] Delay [fs] 3.0  $B_{3g}$  $A_{\mathcal{Q}}$ Raman spectrum 2.5  $\frac{\text{S}}{\text{S}} = \frac{1.5}{1.5}$ <br>Spectrum 1.5  $\Delta A \cdot 10^3$ FT of transient



#### J. Weisshaupt et al., Phys. Rev. B 95, 081101(R) (2017)

### **Time-resolved XMCD using HHG**



Willems et al., Phys. Rev. B 92, 220405(R) (2015)



When absorption edges from deeper-lying core levels are exploited, the theoretical description and – thus – interpretation of XAS experiments becomes easier  $\rightarrow$  drive towards implementation of ATAS in the water window



• Probing of conduction band electrons using  $2p \rightarrow 3d$  Ti L-edge (~460 eV)

#### **B. Buades et al., Appl Phys. Rev. 8, 011408 (2021)**



**B. Buades et al., Appl Phys. Rev. 8, 011408 (2021)**

• Plot the relative change in the absorption (compared to the static absorption, without pump laser)

- The experiment shows that the L-edge absorption increases and decreases as a function of IR-XUV delay, with two oscillations occurring per driver laser period
- These observations are well reproduced by theoretical calculations (both TDDFT and corestate resolved Bloch Equation (cBE) model



• Build-up of charge in the conduction band in TDDFT and cBE calculations

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### What do the oscillations mean?

The authors conclude that as soon as the charge arrives in the conduction band, it starts to perform an intraband motion, with – particular – an increase of the electron density on the Ti atom



#### **B. Buades et al., Appl Phys. Rev. 8, 011408 (2021)**

## **Conclusion**

# **The plan for this week**



Marc Vrakking Federico Furch

# **Key ingredients**

- $\triangleright$  We will discuss how materials can be studied by means of pumpprobe experiments
- $\triangleright$  We will discuss the interaction of materials with intense laser light, and see how this can give rise to the formation of extreme ultraviolet (XUV) attosecond laser pulses
- $\triangleright$  We will discuss how pump-probe experiments can be configured using attosecond laser pulses
- $\triangleright$  We will learn about key experimental techniques, i.e. how intense lasers are built, sample preparation, detection techniques, etc.
- $\triangleright$  To understand the lectures, we will understand a basic knowledge of optics and quantum mechanics – if something is not clear, please ask!

## **Selected take home messages that you will understand on Friday**

- *The proliferation of attosecond science became possible because of major developments in ultrafast laser technology, and attosecond science drives numerous new developments*
- *High-harmonic generation and the formation of attosecond pulses are best understood in a field picture of the laser-matter interaction*
- *Pump-probe experiments are interference experiments, where time-dependent signals arise through interference between multiple quantum paths that connect the initial state of a system to the final state detected in the experiment*
- *The techniques that were introduced in 2001 for the characterization of attosecond pulse trains and isolated attosecond pulses continue to be workhorses in experiments where attosecond pulses are used to probe electron dynamics*
- *The attosecond pulse structure of High-harmonic radiation is not its only important feature, HHG moreover permits core-level specific time-resolved spectroscopy with enormous impact in atomic, molecular and condensed phase physics*
- *The future of attosecond science looks extremely bright!*

*Thank you for your attention!*

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