

Lecture Series Buenos Aires

18-3-2024 until 22-3-2024

Lecture F6 – Optical Parametric Chirped Pulse Amplification

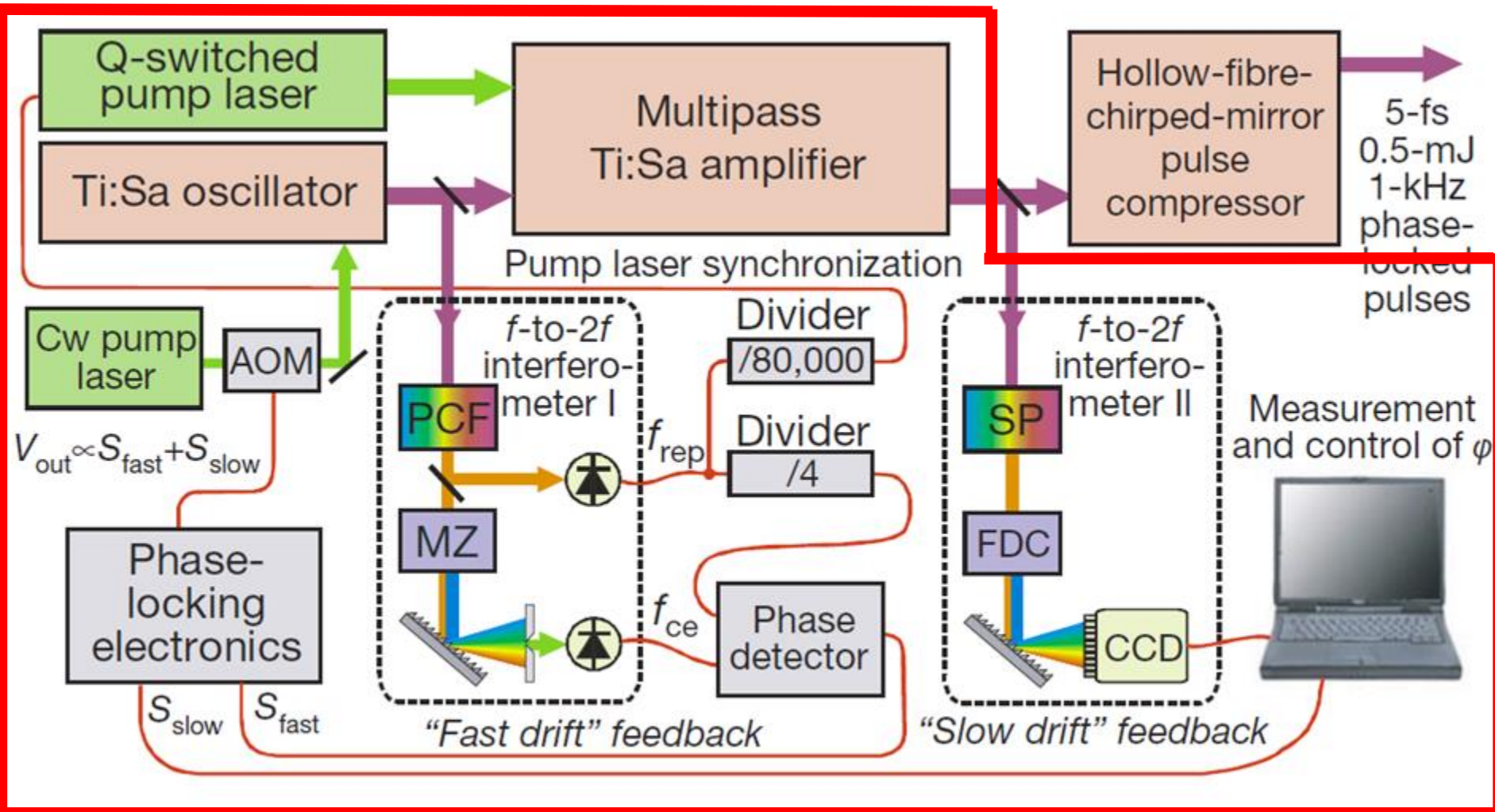


Max-Born-Institut

Federico Furch
furch@mbi-berlin.de

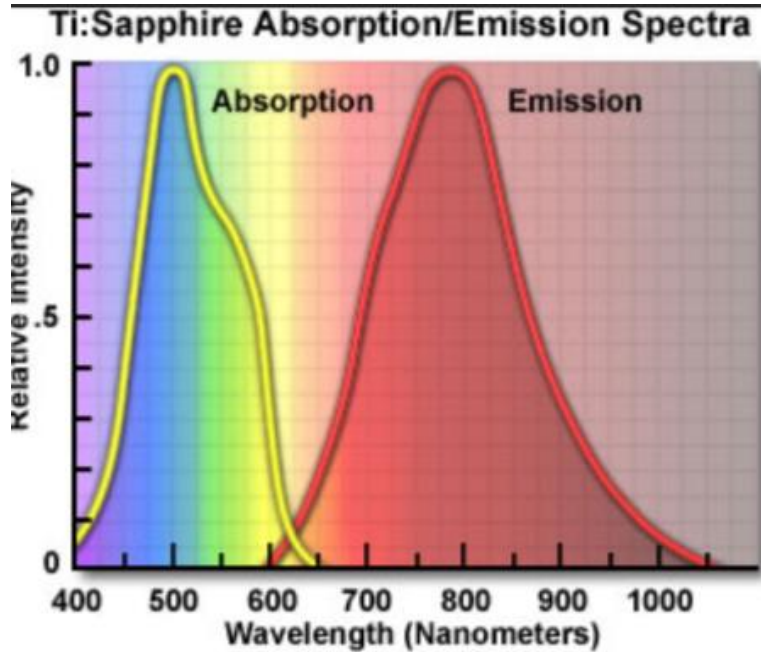
Optical Parametric Chirped Pulse Amplification (OPCPA)

A state-of-the-art laser system for attosecond science: an alternative

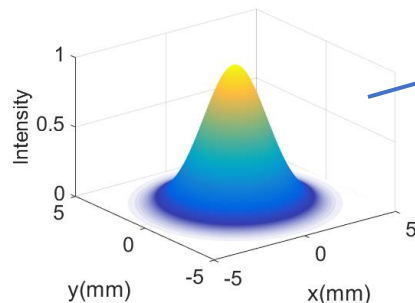


Limitations of Ti:Sapphire CPAs

Limitations to spectral range



$$P_{pump}(x, y) \propto e^{-2(x^2+y^2)/w^2}$$



Heat source profile: may originate

- Thermal lensing (dn/dT)
- Thermal induced birefringence
- Damage of material

Limitations to pulse duration

$$\text{gain} \sim e^{g(\omega)L}$$

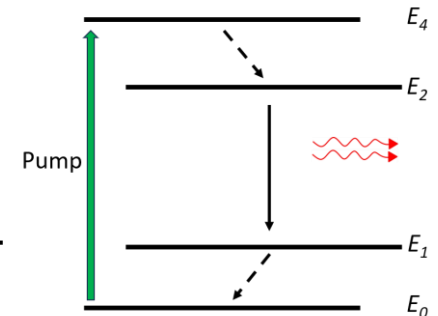
$$\text{gain narrowing} \quad \Delta\omega_{out} \ll \Delta\omega_{g(\omega)}$$

Limitations to power scaling

$$P_{avg} = \text{Energy}_{pulse} * \text{frep.rate}$$

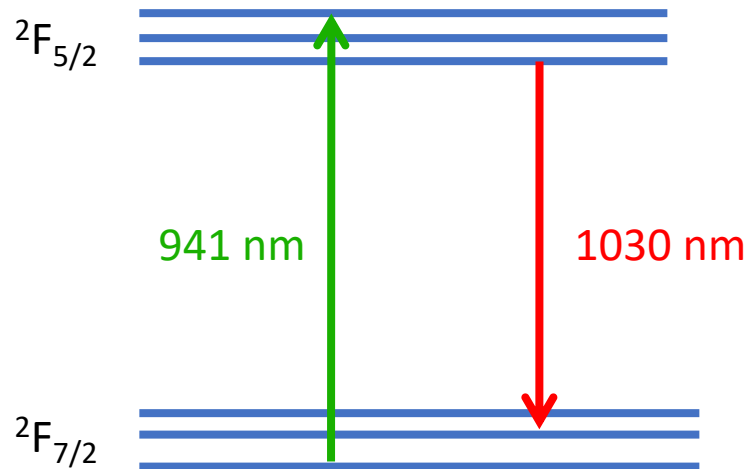
Fraction of pump
that turns into

$$\text{heat} \quad 1 - \frac{\hbar\omega_{laser}}{\hbar\omega_{pump}}$$



A laser material for high average power

Yb-doped Yttrium
Aluminum Garnet
Yb:YAG



- Absorption band at InGaAs wavelengths
 - High power laser diodes are commercially available
- Low quantum defect ($1 - \frac{\hbar\omega_{laser}}{\hbar\omega_{pump}} < 0.1$)
 - Potential for high average power operation
- Long upper level lifetime (~ 1 msec)
 - Efficiently store energy from low peak power pump
- High quality (large) crystals
 - Crystalline or ceramic form
- **BUT narrow gain bandwidth: post-compression, OPCAs**

Alternatives using Yb systems

Ti:Sapphire → ultrashort pulses

Yb-doped → high energy, high average power

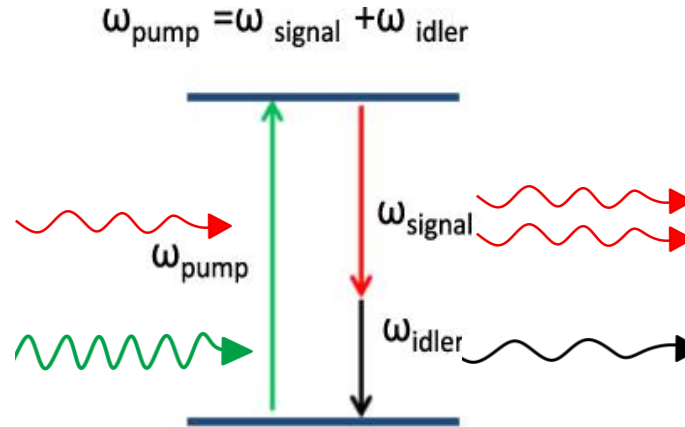
Energy transfer in Optical Parametric Amplifier

Yb-doped → high energy, high average power

Nonlinear pulse compression with large compression factors to reach sub-50 fs and even sub-10 fs pulses

Parametric Amplification

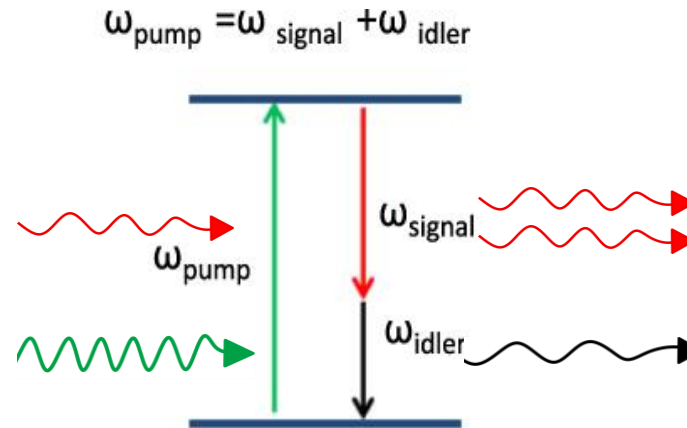
- No absorption



Instantaneous nonlinear process of second order in the E-field

Parametric Amplification

- No absorption



Instantaneous nonlinear process of second order in the E-field

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

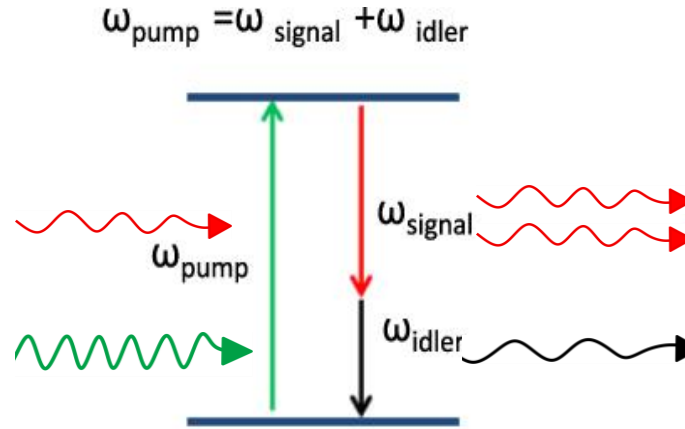
$$\nabla \times \mathbf{B} = \mu_0 \left(\mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \right)$$

$$\hat{\mathbf{D}}(\mathbf{r}, \omega, z) = \epsilon_0 \epsilon(\omega) \hat{\mathbf{E}}(\mathbf{r}, \omega, z) + \hat{\mathbf{P}}(\mathbf{r}, \omega, z),$$

$$\mathbf{P} = \mathbf{P}^{(1)} + \mathbf{P}^{(NL)}, \quad \mathbf{P}^{(NL)} \propto \chi^{(2)} \mathbf{E}^2$$

Parametric Amplification

- No absorption
- Gain bandwidth determined by phase-matching

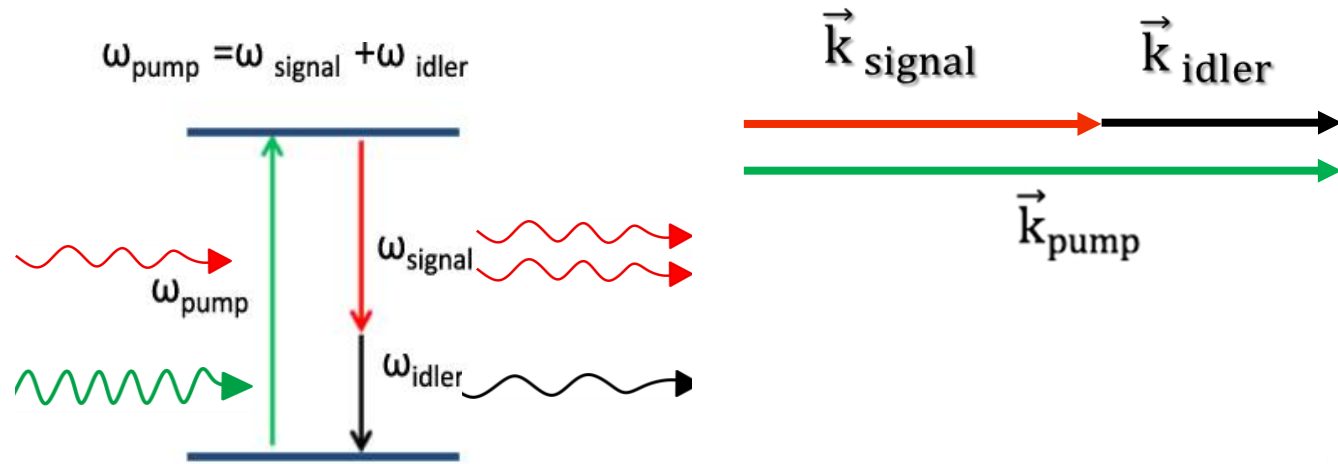


Phase-matching

$$\Delta \mathbf{k} = \mathbf{k}_{\text{signal}} + \mathbf{k}_{\text{idler}} - \mathbf{k}_{\text{pump}} = \mathbf{0}$$

Parametric Amplification

- No absorption
- Gain bandwidth determined by phase-matching



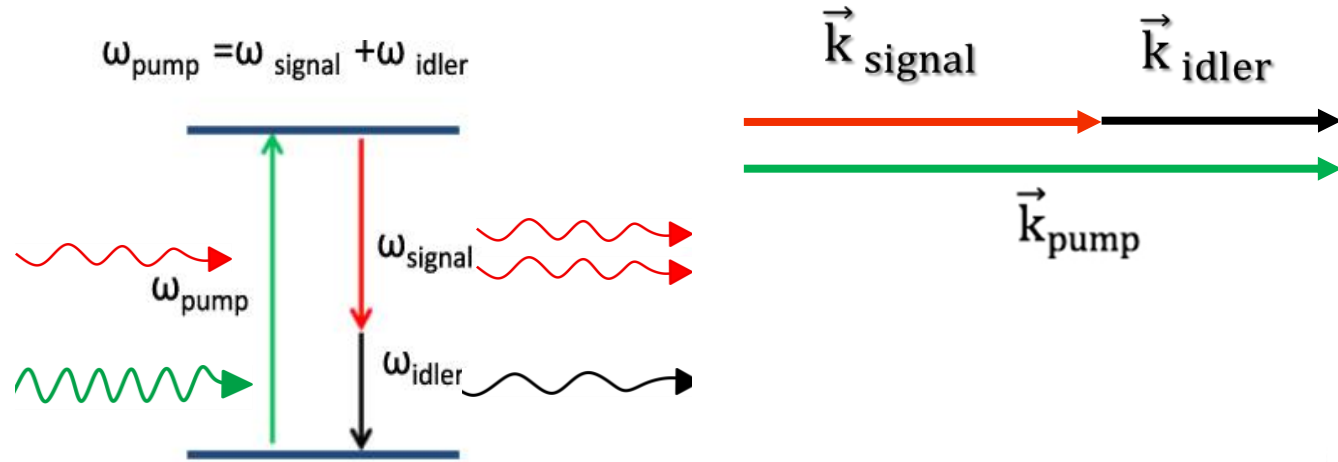
Phase-matching

$$\Delta \mathbf{k} = \mathbf{k}_{\text{signal}} + \mathbf{k}_{\text{idler}} - \mathbf{k}_{\text{pump}} = \mathbf{0}$$
$$k(\omega) = \frac{\omega n(\omega)}{c}$$

Use bi-refringent material

Parametric Amplification

- No absorption
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Phase-matching

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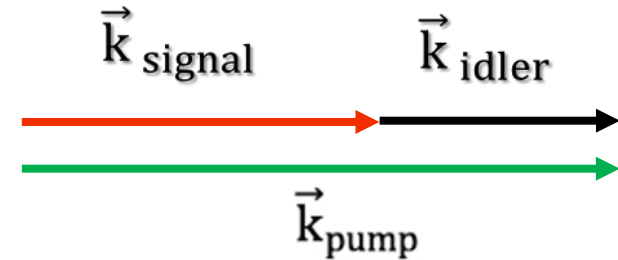
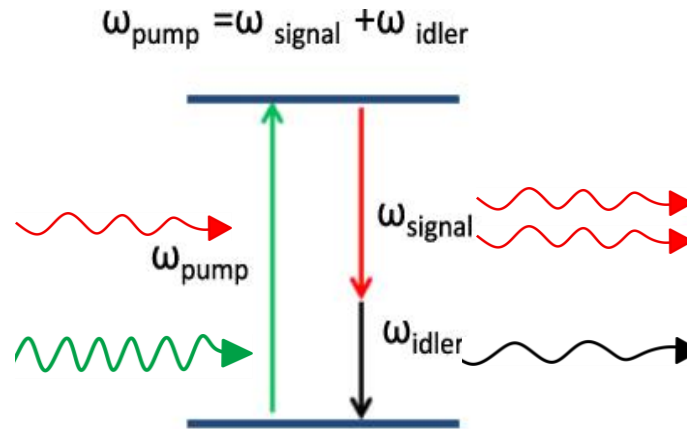
Type-I phase matching: $k_{\text{pump}}(\omega) = \frac{\omega n(\omega, \theta)}{c}$

θ : Angle between k-vector and crystal axis

$$\frac{1}{n_e(\theta)^2} = \frac{\sin^2 \theta}{n_{e,90^\circ}^2} + \frac{\cos^2 \theta}{n_o^2}$$

Parametric Amplification

- No absorption
- Gain bandwidth determined by phase-matching



Phase-matching

$$\Delta \mathbf{k} = \mathbf{k}_{\text{signal}} + \mathbf{k}_{\text{idler}} - \mathbf{k}_{\text{pump}} = \mathbf{0}$$

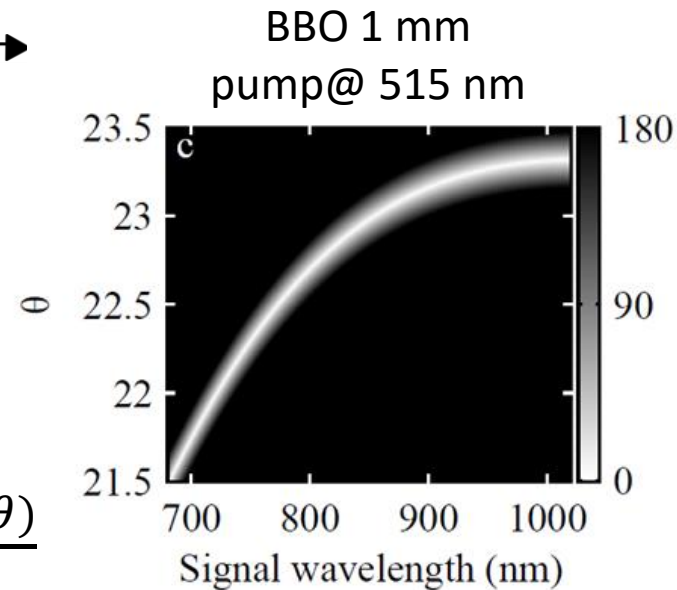
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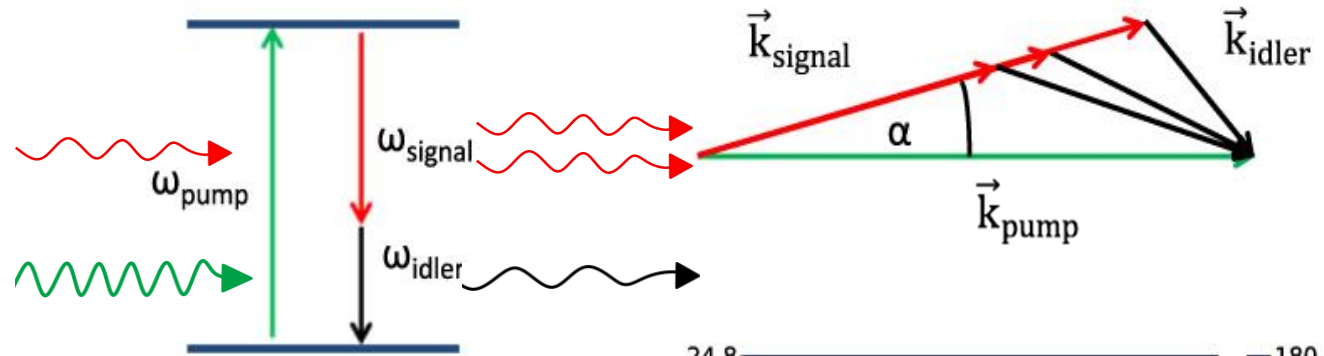
$$\frac{1}{n_e(\theta)^2} = \frac{\sin^2 \theta}{n_{e,90^\circ}^2} + \frac{\cos^2 \theta}{n_o^2}$$



Parametric Amplification

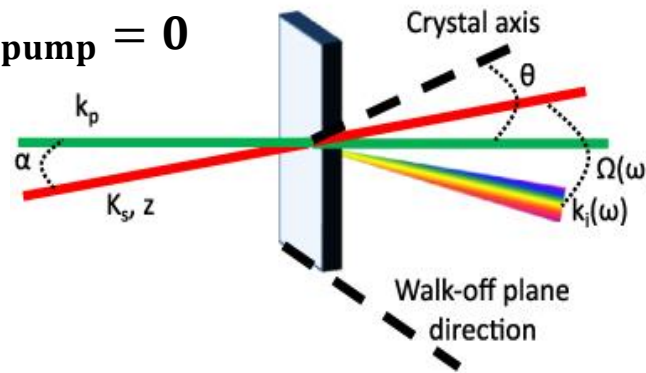
- No absorption
- Gain bandwidth determined by phase-matching

$$\omega_{\text{pump}} = \omega_{\text{signal}} + \omega_{\text{idler}}$$



Phase-matching

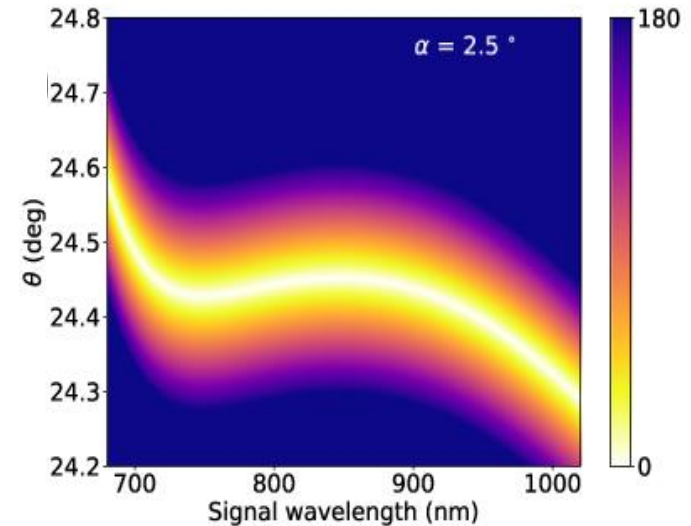
$$\Delta \mathbf{k} = \mathbf{k}_{\text{signal}} + \mathbf{k}_{\text{idler}} - \mathbf{k}_{\text{pump}} = \mathbf{0}$$



- Noncollinear geometry allows broadband amplification

Type-I phase matching:

$$k_{\text{pump}}(\omega) = \frac{\omega n(\omega, \theta)}{c}$$



Phase-mismatch ($\Delta \mathbf{k}$):
1mm BBO Crystal

Optical Parametric Amplification

OBSERVATION OF PARAMETRIC AMPLIFICATION IN THE OPTICAL RANGE

S. A. Akhmanov, A. I. Kovrigin, A. S. Piskarskas, V. V. Fadeev, and R. V. Khokhlov

Physics Faculty, Moscow State University

Submitted 23 July 1965

We report here the results of an experiment in which we observed directly parametric amplification of an optical signal with wavelength $\lambda_s = 1.06 \mu$ in a KDP crystal excited by an intense pump wave with $\lambda_p = 0.53 \mu$. The feasibility of such an effect in the optical band and its theory were detailed in [1-3]; results of experiments in which parametric amplification at wavelength $\lambda_s = 0.63 \mu$ has been indirectly registered are described in [4].

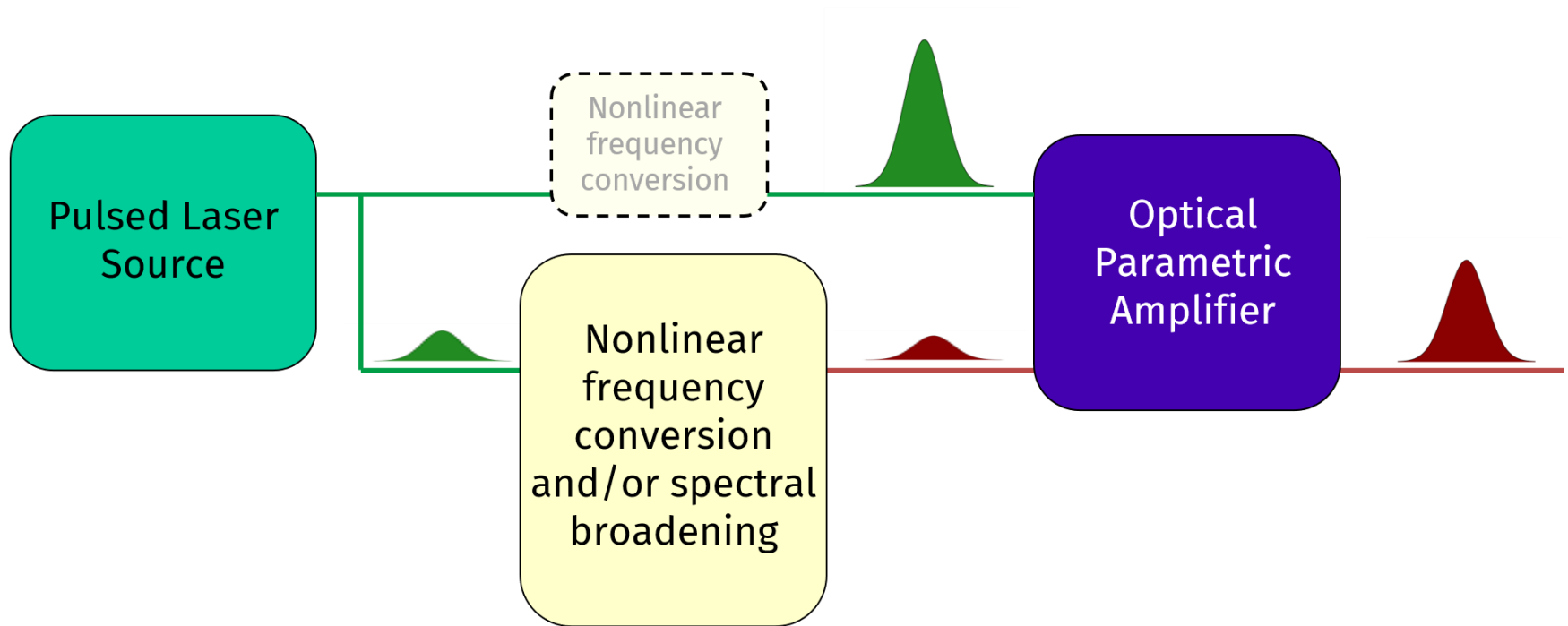
In a nonlinear medium with a polarization that depends quadratically on the magnetic field intensity, the energy of an intense pump wave (frequency ω_p) can be transferred to waves with frequencies ω_1 and ω_2 satisfying the relation $\omega_p = \omega_1 + \omega_2$. The energy transfer is most effective if the following relation is satisfied between the wave vectors of the interacting waves (the so-called synchronism condition):

$$\vec{k}_1 + \vec{k}_2 = \vec{k}_p.$$

(1)

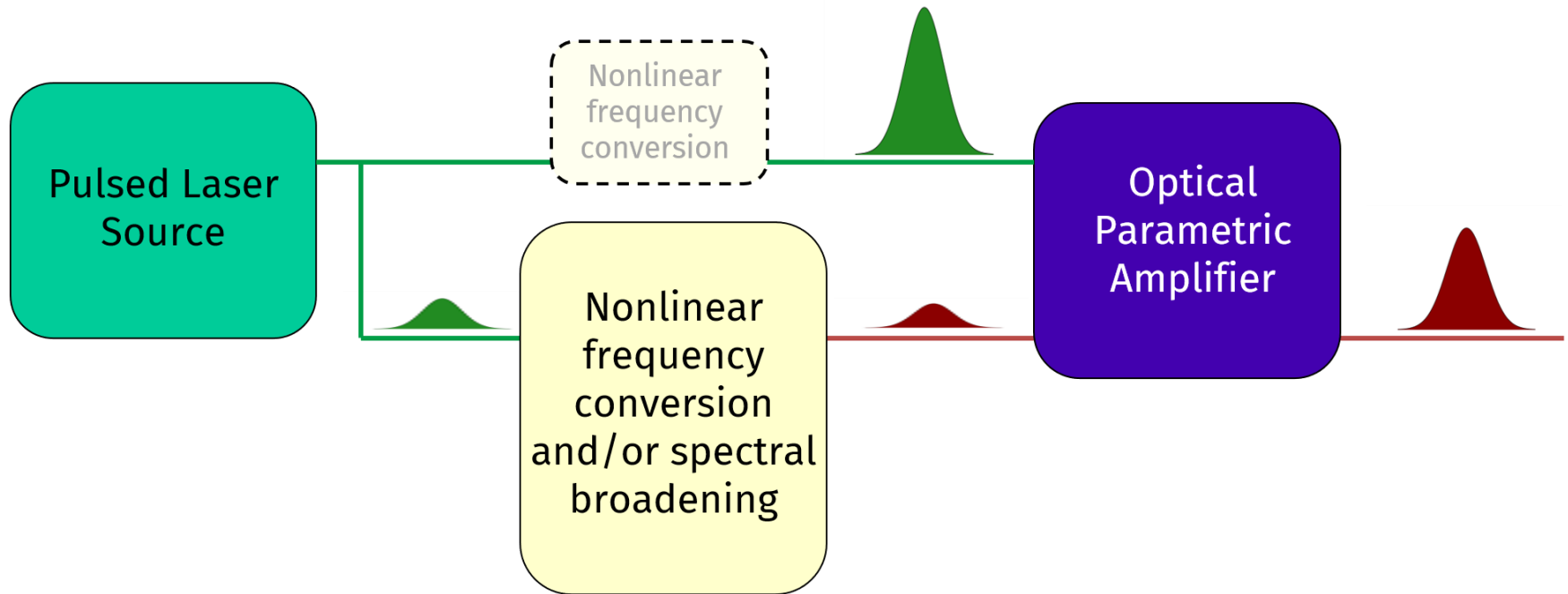
Optical Parametric Amplification

OPA: Great way to go to a different part of the spectrum
(if I already have short pulses with high energy)



Optical Parametric Amplification

OPA: Great way to go to a different part of the spectrum (if I already have short pulses with high energy)



Example: start from 800 nm and go to 1300 nm to generate higher frequencies during HHG, or IAP combining HHG with 800 nm + 1300 nm

Optical Parametric Amplification

1494 OPTICS LETTERS / Vol. 22, No. 19 / October 1, 1997

Sub-20-fs pulses tunable across the visible from a blue-pumped single-pass noncollinear parametric converter

T. Wilhelm, J. Piel, and E. Riedle

Institut für Medizinische Optik, Ludwig-Maximilians-Universität München, Barbarastrasse 16, D-80797 München, Germany

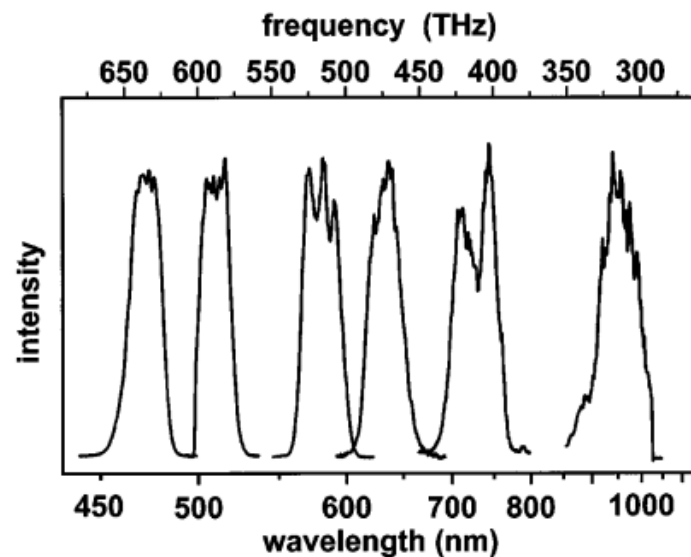
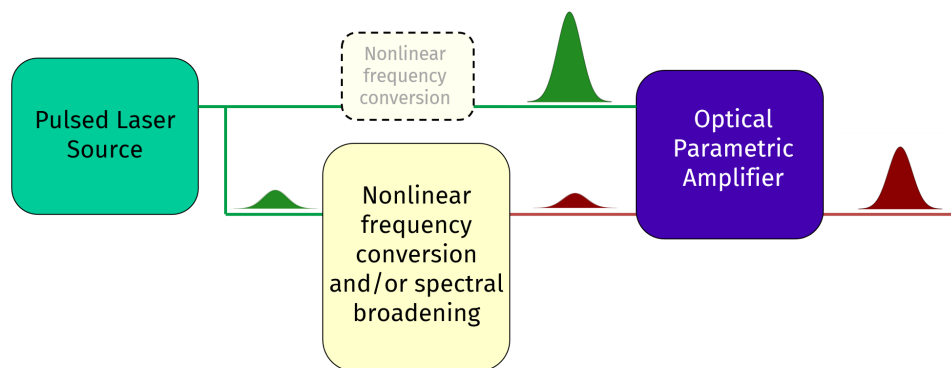
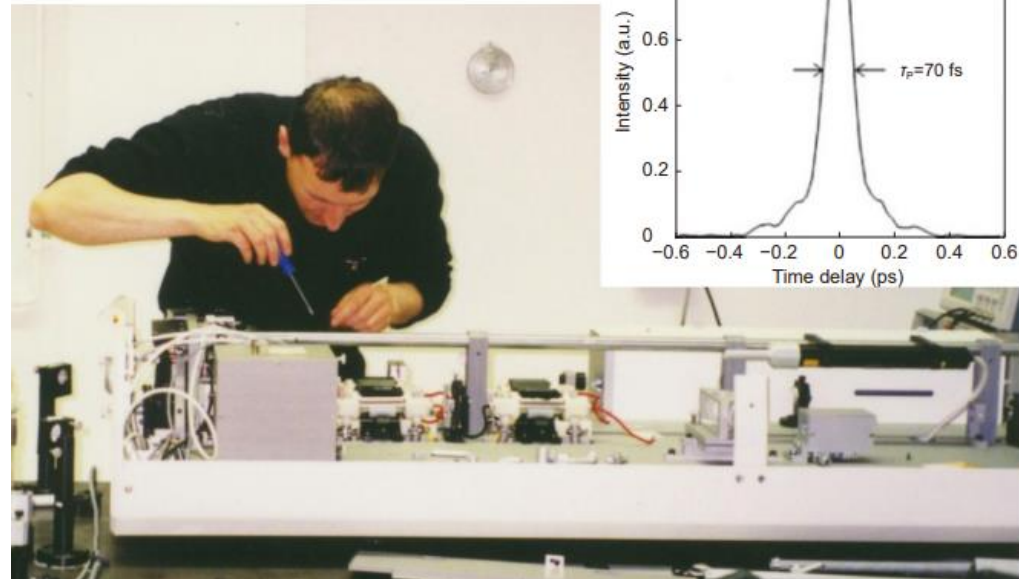
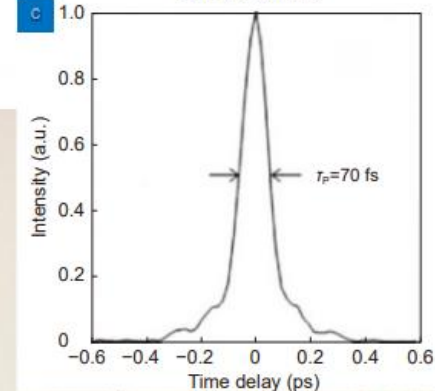
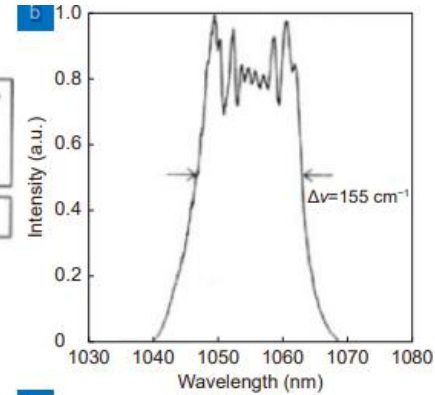
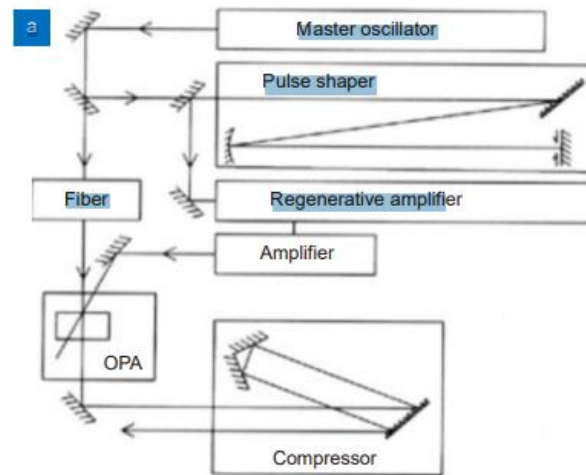


Fig. 2. Spectra of femtosecond pulses generated by non-collinear parametric amplification in a 2-mm BBO crystal.

Optical Parametric Chirped Pulse Amplification (OPCPA)

First demonstration of OPCPA: A. Dubietis, G. Jonušauskas, A. Piskarskas. *Powerful femtosecond pulse generation by chirped and stretched pulse parametric amplification in BBO crystal*. Opt Commun **88**, 437–440 (1992)

- 1.7ps, 1055nm pulses from a Nd:glass oscillator, spectrally broadened (and stretched) in a fiber
- Amplified in BBO crystal
- Compressed with grating pair: 70fs, 65μJ



Optical Parametric Chirped Pulse Amplification (OPCPA)

1 December 1997



OPTICS
COMMUNICATIONS

Optics Communications 144 (1997) 125–133

Full length article

The prospects for ultrashort pulse duration and ultrahigh intensity using optical parametric chirped pulse amplifiers

I.N. Ross ^{*}, P. Matousek, M. Towrie, A.J. Langley, J.L. Collier

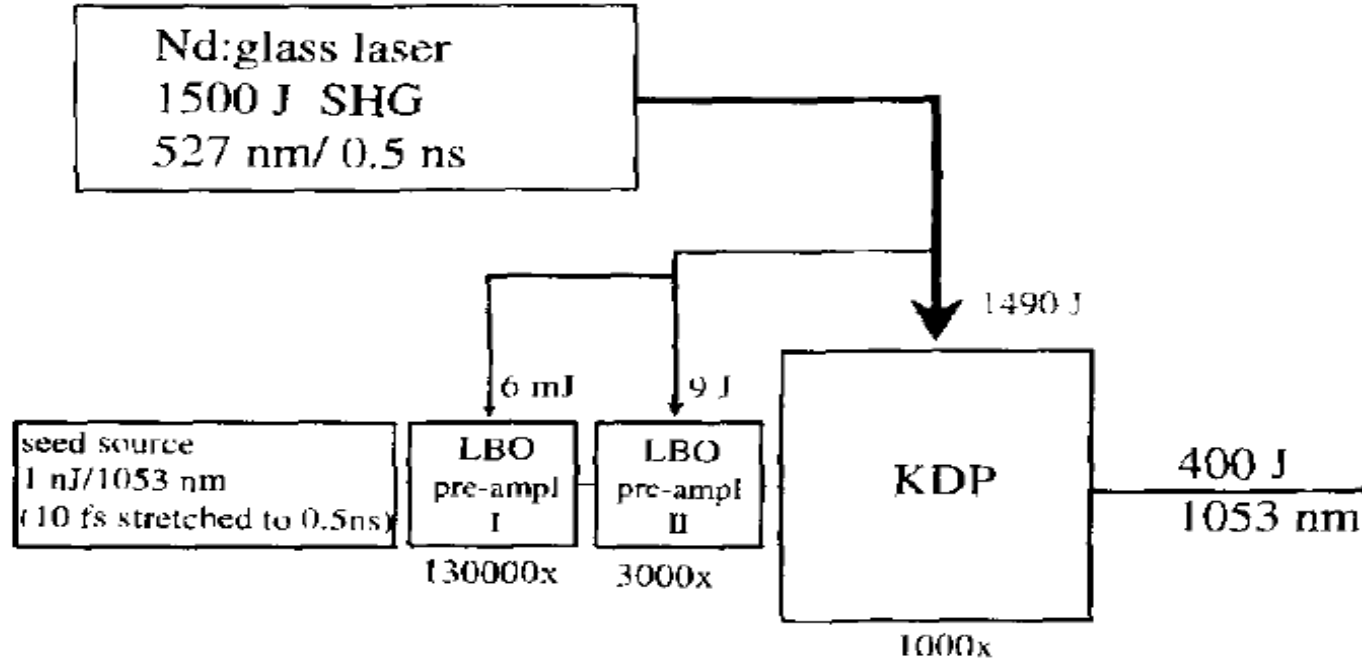
Central Laser Facility, Rutherford Appleton Laboratory, Chilton, Didcot, Oxon OX11 0QX, UK

Received 27 May 1997; accepted 8 July 1997

Abstract

The properties of optical parametric amplifiers are presented to show that, in addition to their use in providing tuneable pulses, they can form the critical component in systems generating shorter pulse duration with higher power and intensity than is possible with existing systems. Practical designs are evaluated leading to anticipated powers exceeding 10 PW and focused intensities greater than 10^{23} W cm⁻². © 1997 Elsevier Science B.V.

Optical Parametric Chirped Pulse Amplification (OPCPA)

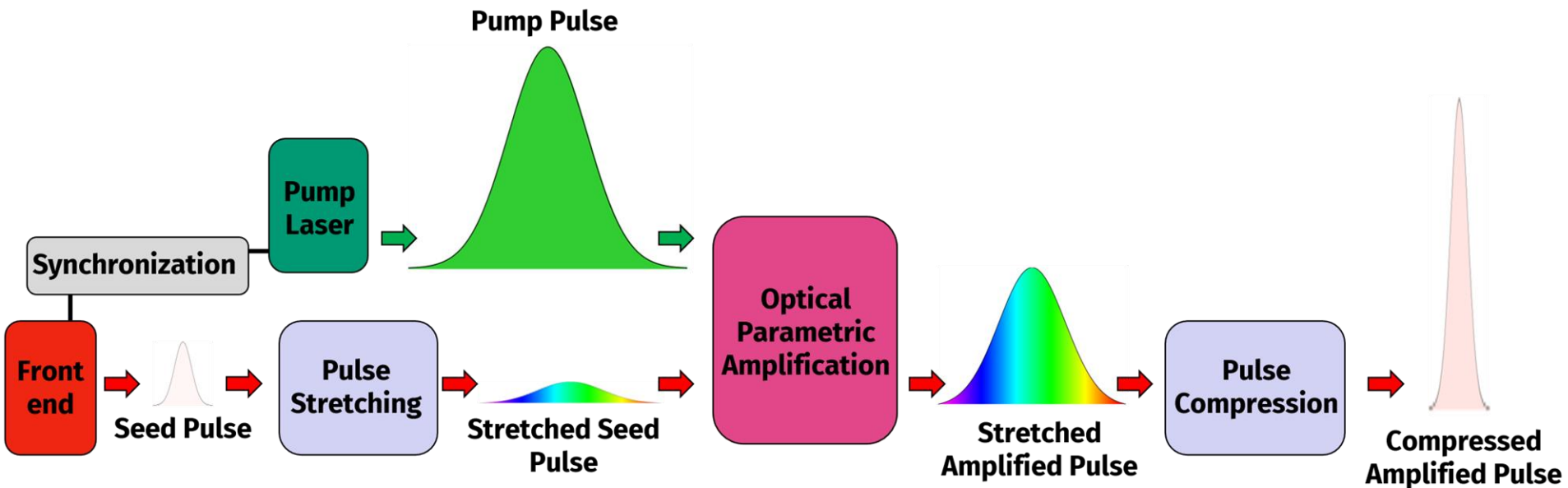


	beam diameter (mm)	crystal length (mm)	pump fluence (J/cm ²)
LBO / pre-amplifier I	0.6	11.5	2
LBO / pre-amplifier II	20	6.8	2.9
KDP	305	28.6	2

Fig. 6. PW OPCPA design for a high power Nd:glass laser.

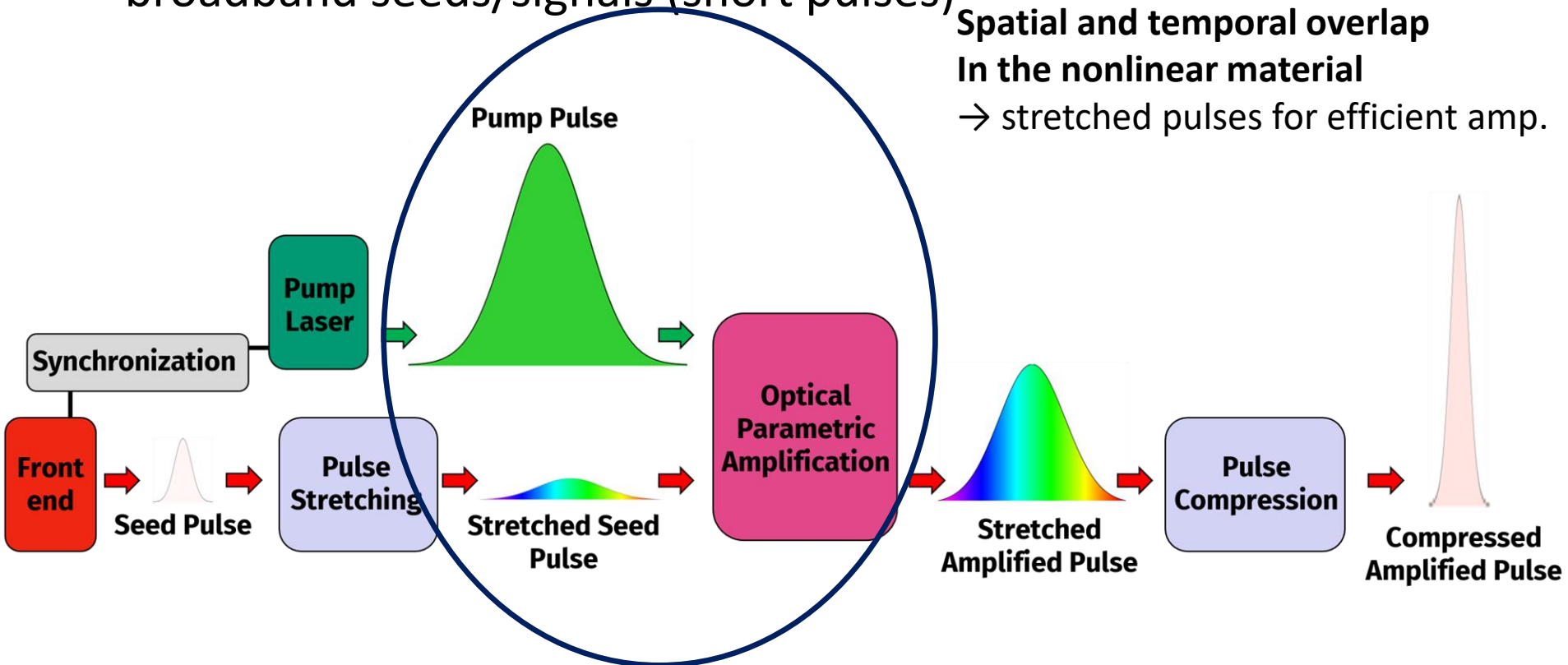
Optical Parametric Chirped Pulse Amplification (OPCPA)

- Amplification of short pulses avoiding unwanted nonlinear effects

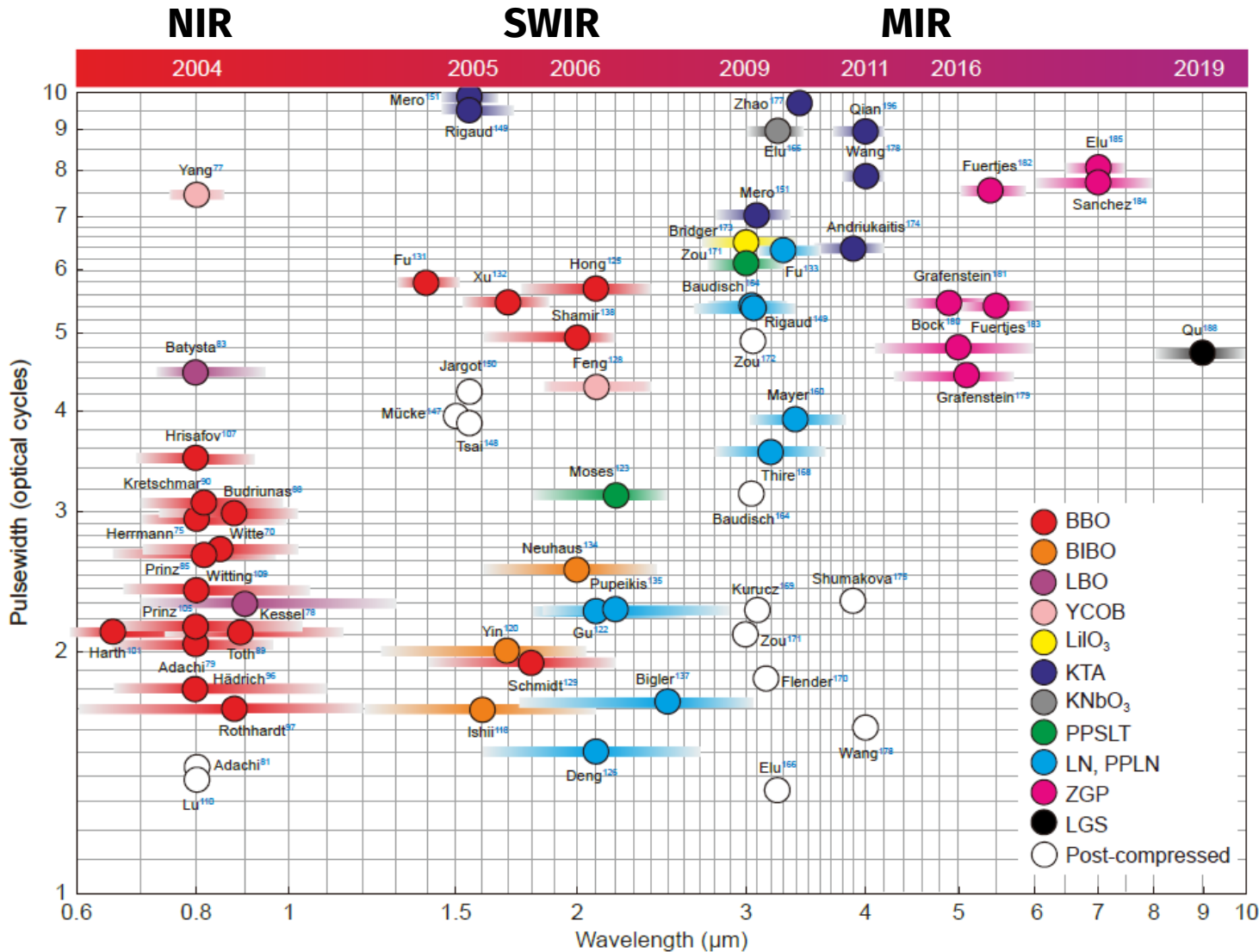


Optical Parametric Chirped Pulse Amplification (OPCPA)

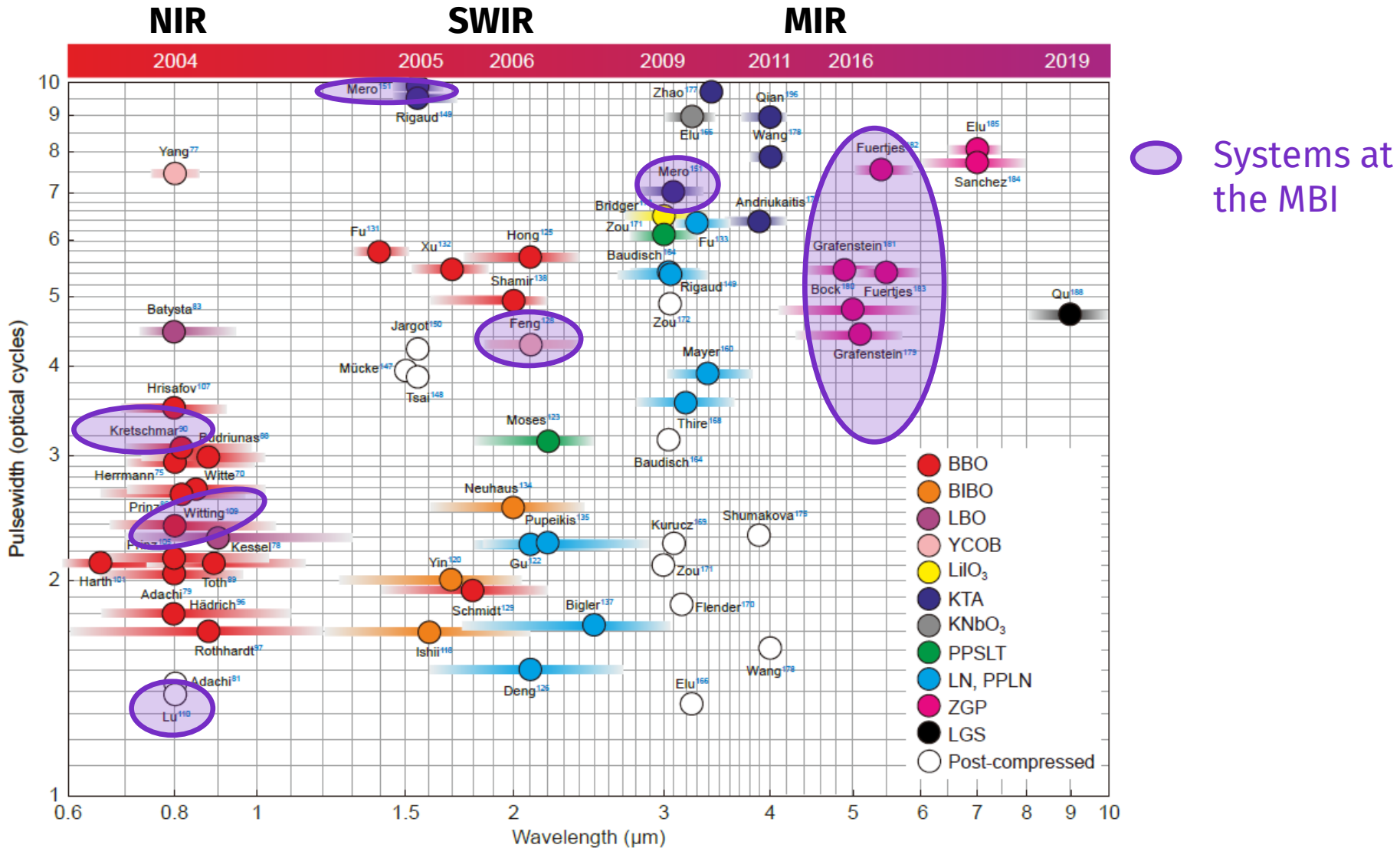
- Amplification of short pulses avoiding unwanted nonlinear effects
- Using narrowband pumps (long pulses) to amplify broadband seeds/signals (short pulses)



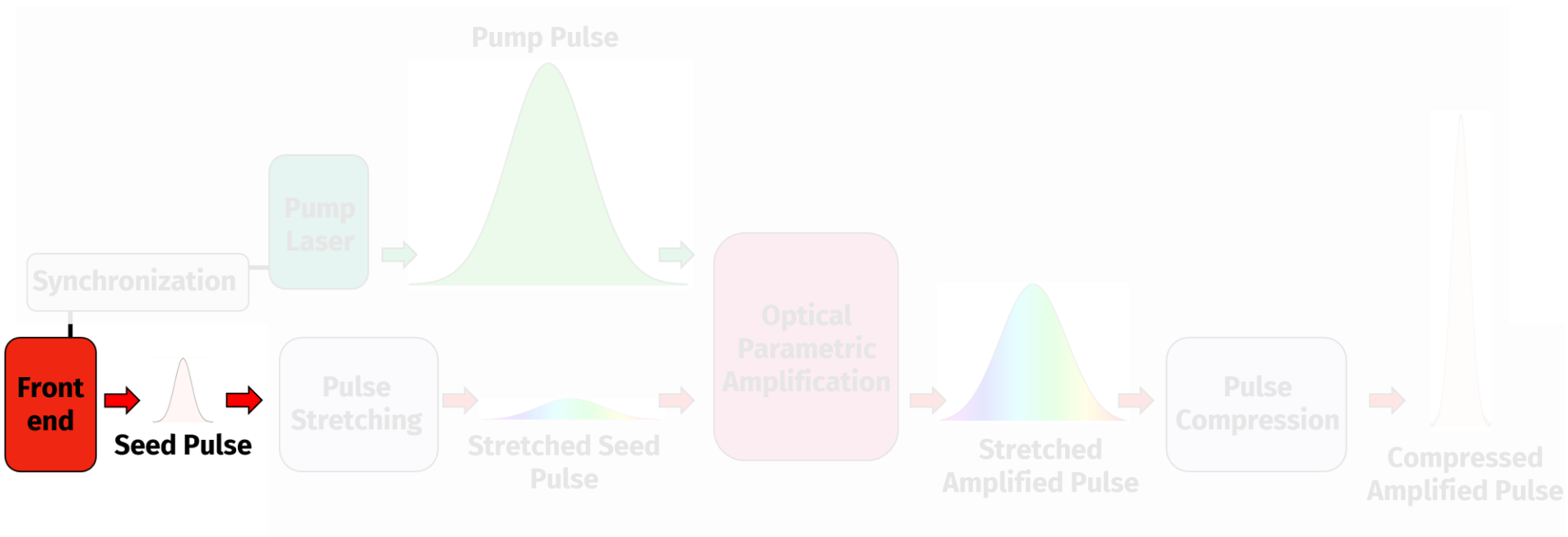
Few-cycle OPCPAs from VIS to MIR



Few-cycle OPCPAs from VIS to MIR



Building blocks



Front end

March 15, 2001 / Vol. 26, No. 6 / OPTICS LETTERS

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Generation of 5-fs pulses and octave-spanning spectra directly from a Ti:sapphire laser

R. Ell, U. Morgner, and F. X. Kärtner

High Frequency and Quantum Electronics Laboratory, University of Karlsruhe, D-76128 Karlsruhe, Germany

J. G. Fujimoto and E. P. Ippen

Department of Electrical Engineering and Computer Science and Research Laboratory of Electronics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139

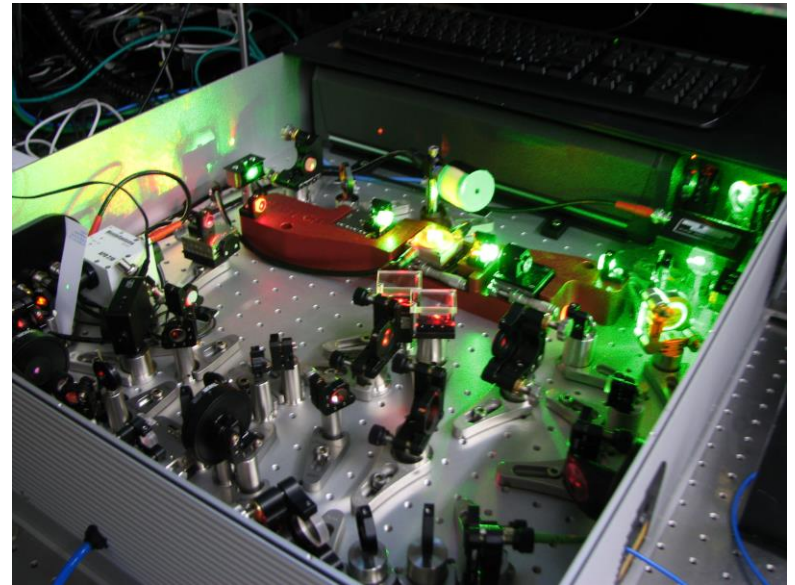
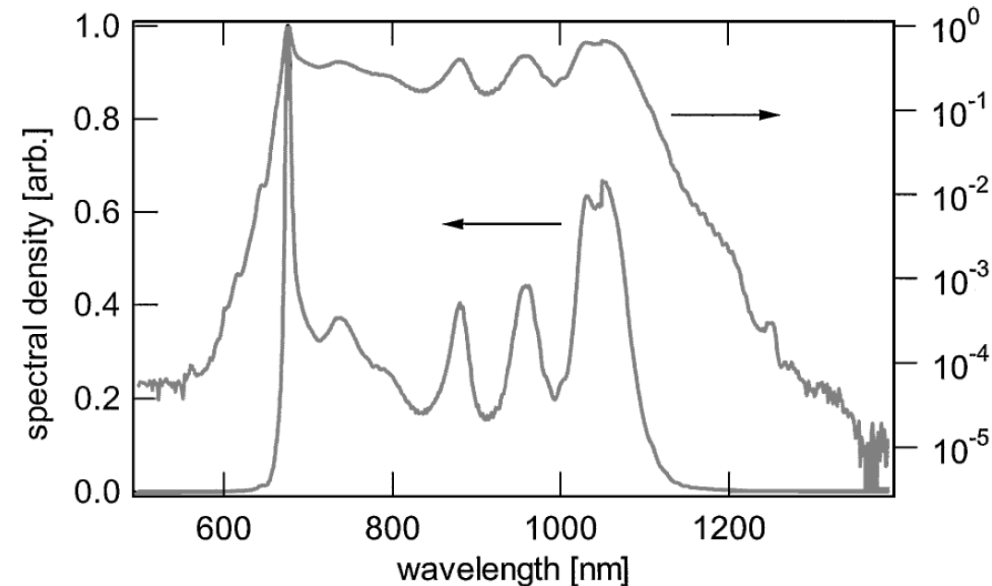
V. Scheuer, G. Angelow, and T. Tschudi

Institute for Applied Physics, TU Darmstadt, Germany

M. J. Lederer, A. Boiko, and B. Luther-Davies

Ultra-broadband Ti:Sapphire oscillators:

- Support ultrashort, down to few-cycle pulses
- Optical Synchronization
- CEP stability

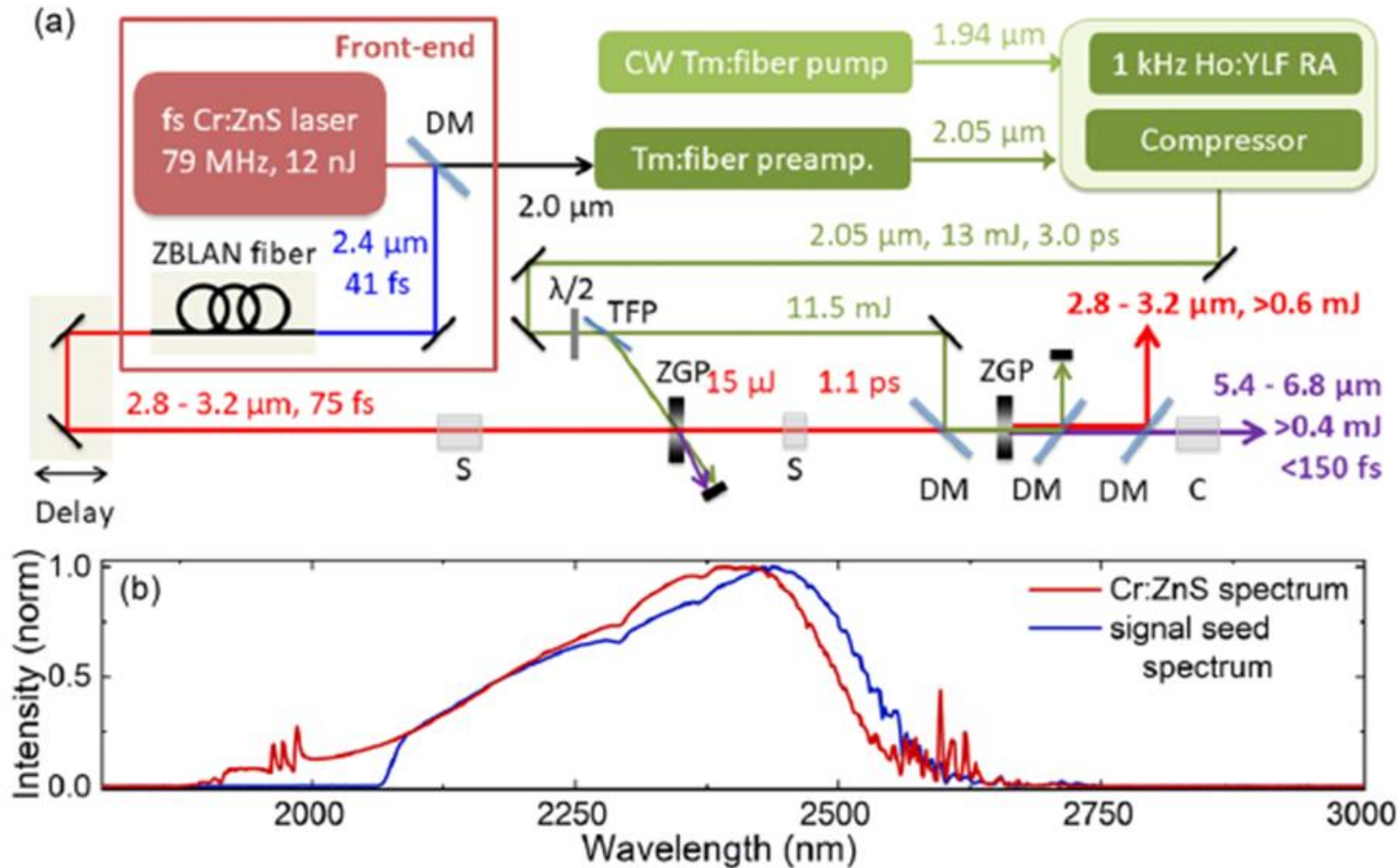


Photograph of the Vteon Ti:Sapphire oscillator at MBI

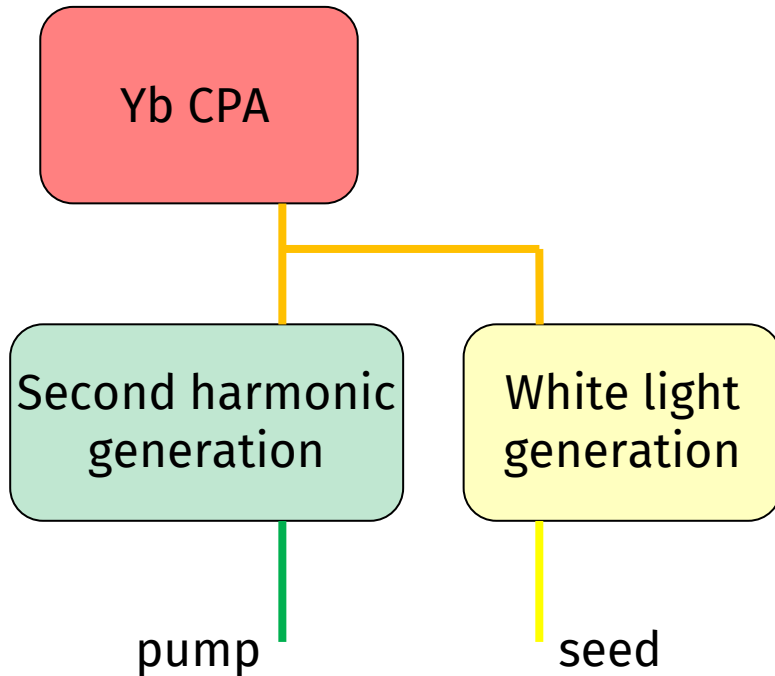
Front end

Broadband Cr:ZnS or Cr:ZnSe oscillators:

- Ultrashort pulses in the MWIR
- Optical Synchronization



Front end



Yb-based system+different nonlinear processes

- Access to different parts of the spectrum
- Seed for OPCPA pump (optical synchronization)
- Passive CEP stabilization

Front end

Yb-based system+different nonlinear processes

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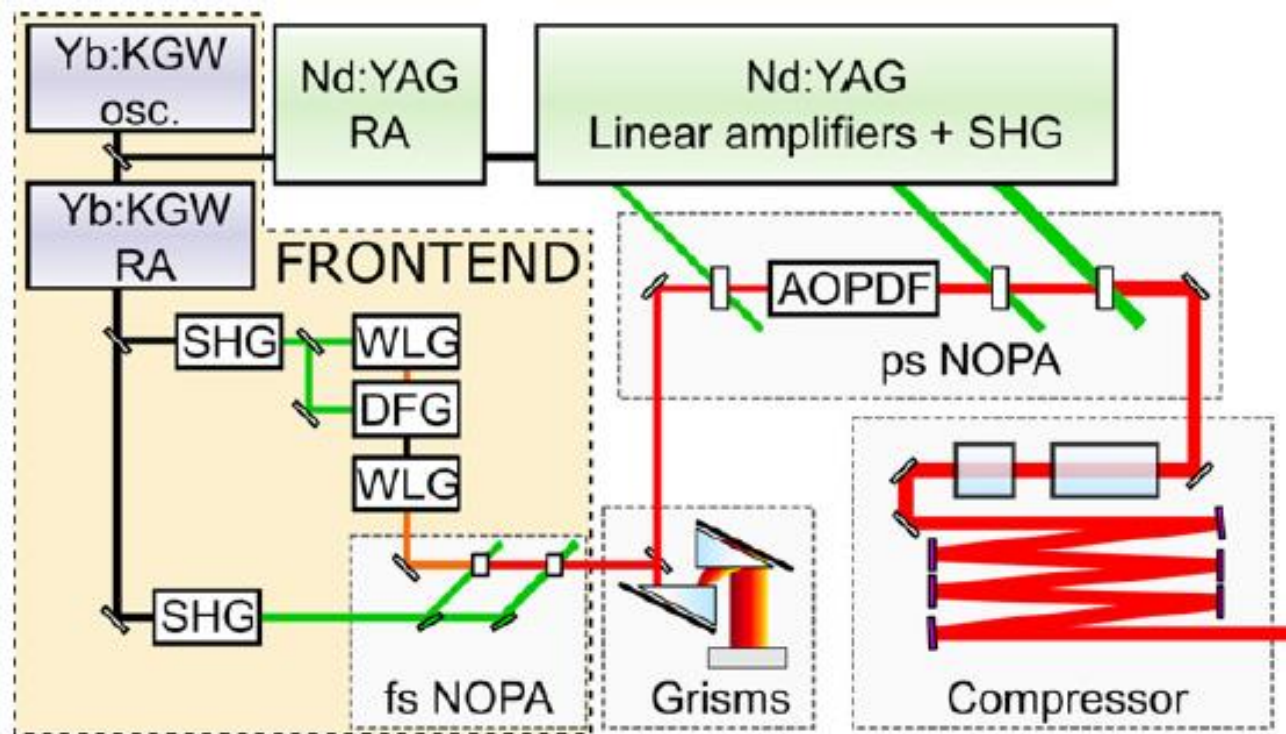
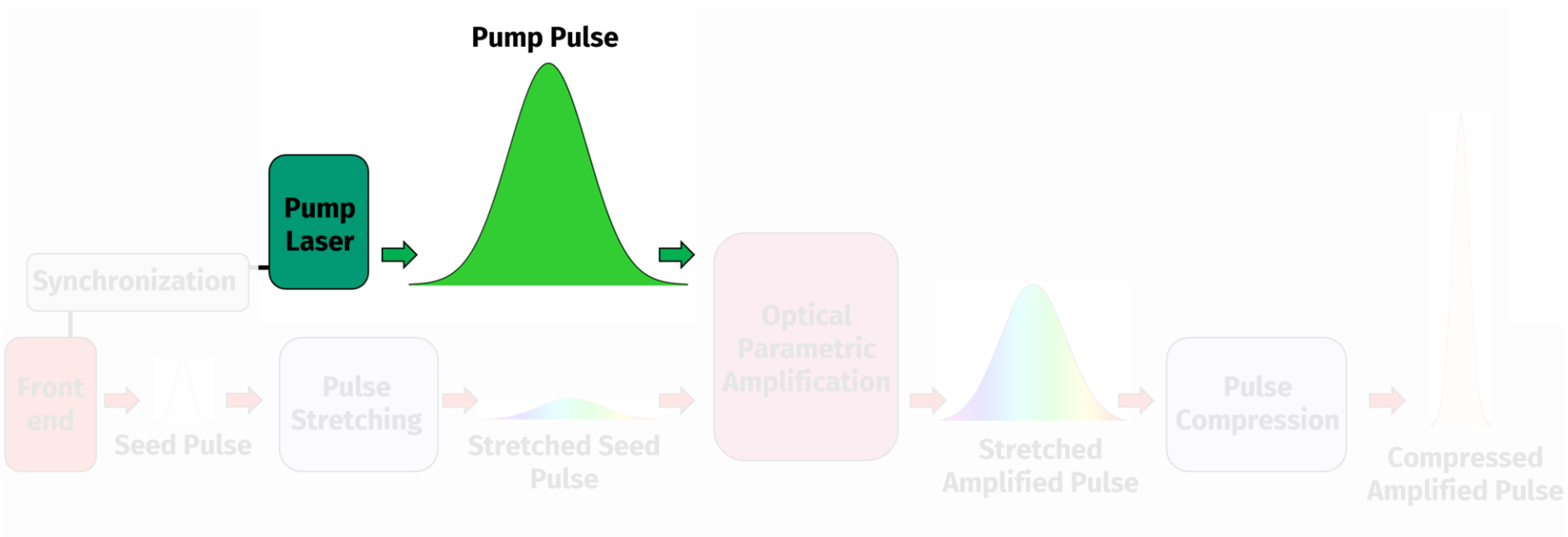


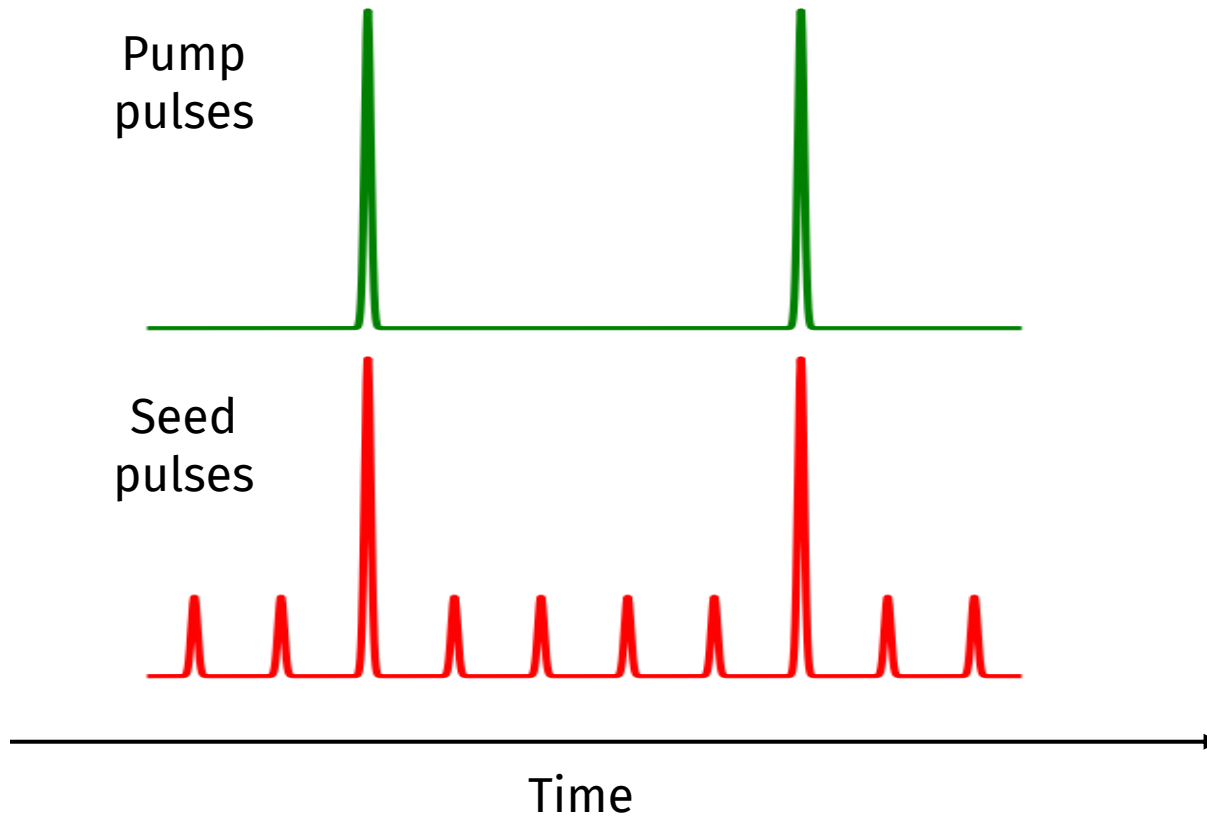
Figure 1. Principal layout of the OPCPA setup.

Building blocks



Pump

- Determines repetition rate and energy of the system

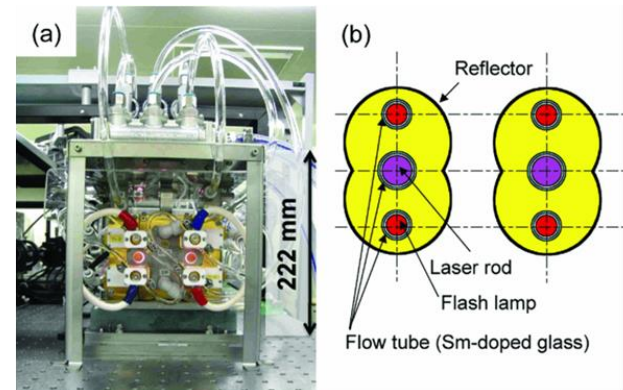


Pump

- Determines repetition rate and energy of the system
 - Flash lamp pumped systems: high energy at low repetition rates
 - DPSSL and fibers: high repetition rate, high average powers



Multipass thin disk Yb:YAG amplifier part of a DPSSL developed at the MBI



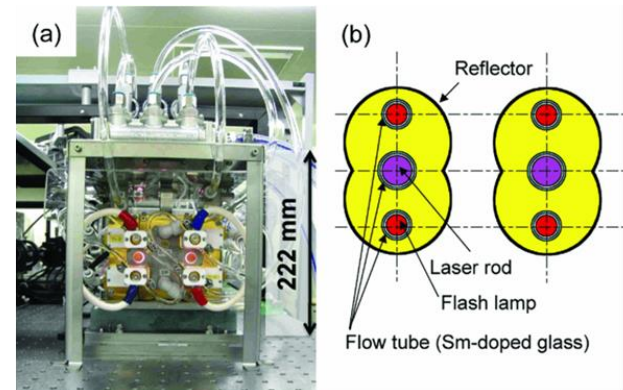
Flashed lamp pumped amplifier. Image from Hatae et al., Rev Sci Instrum **83**, 10E344 (2012)

Pump

- Determines repetition rate and energy of the system
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- Determines the chirp management
 - ns, ~ 100 ps pulses \Rightarrow grating stretcher / compressor
 - Sub-10ps pulses \Rightarrow material dispersion / dispersive mirrors / prism pairs



Multipass thin disk Yb:YAG amplifier part of a DPSSL developed at the MBI



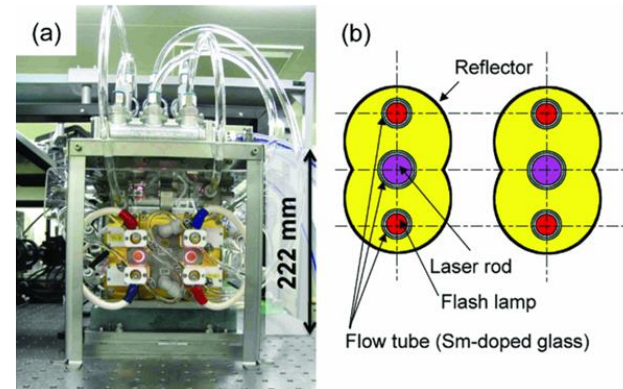
Flashed lamp pumped amplifier. Image from Hatae et al., Rev Sci Instrum **83**, 10E344 (2012)

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 - ns, ~ 100 ps pulses \Rightarrow grating stretcher / compressor
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- Different gain materials for pumping OPCPA in different spectral regions
 - $1\mu\text{m}$ systems for pumping NIR/SWIR: Nd:YAG, Nd:YLF, Nd:YVO₄, Yb:YAG, Yb:KGW, Yb-doped fibers
 - $2\mu\text{m}$ systems for pumping MIR: Ho:YAG, Ho:YLF, Tm:YLF

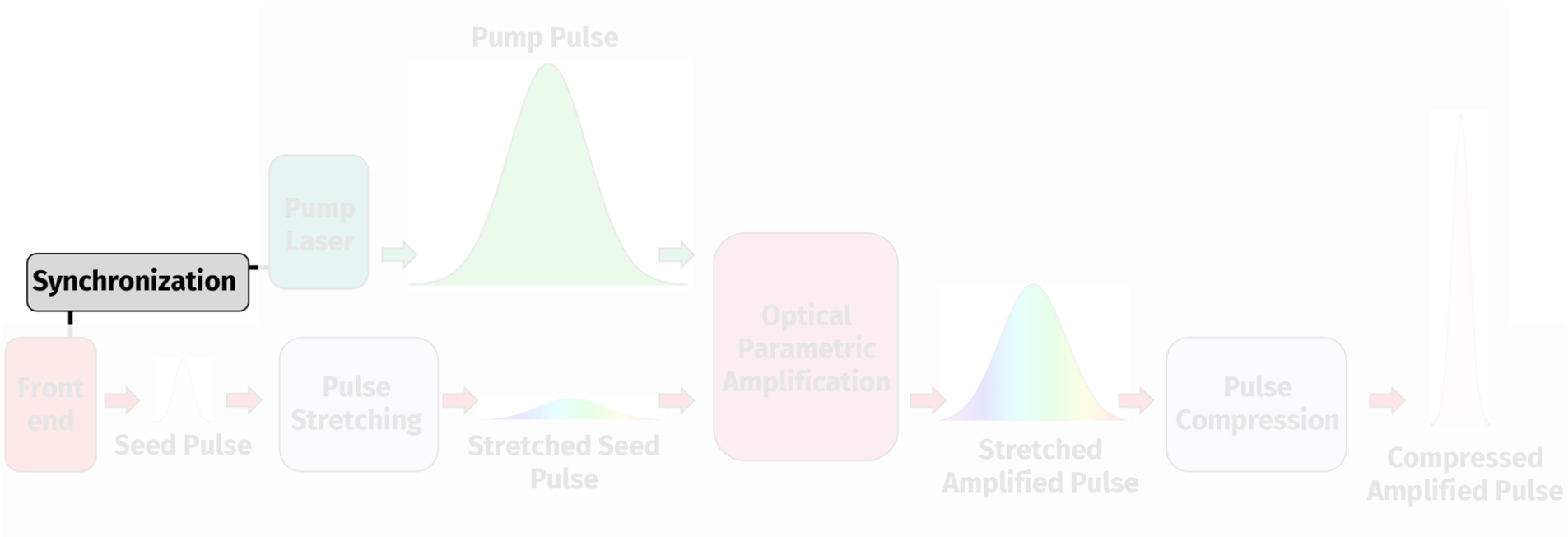


Multipass thin disk Yb:YAG amplifier part of a DPSSL developed at the MBI



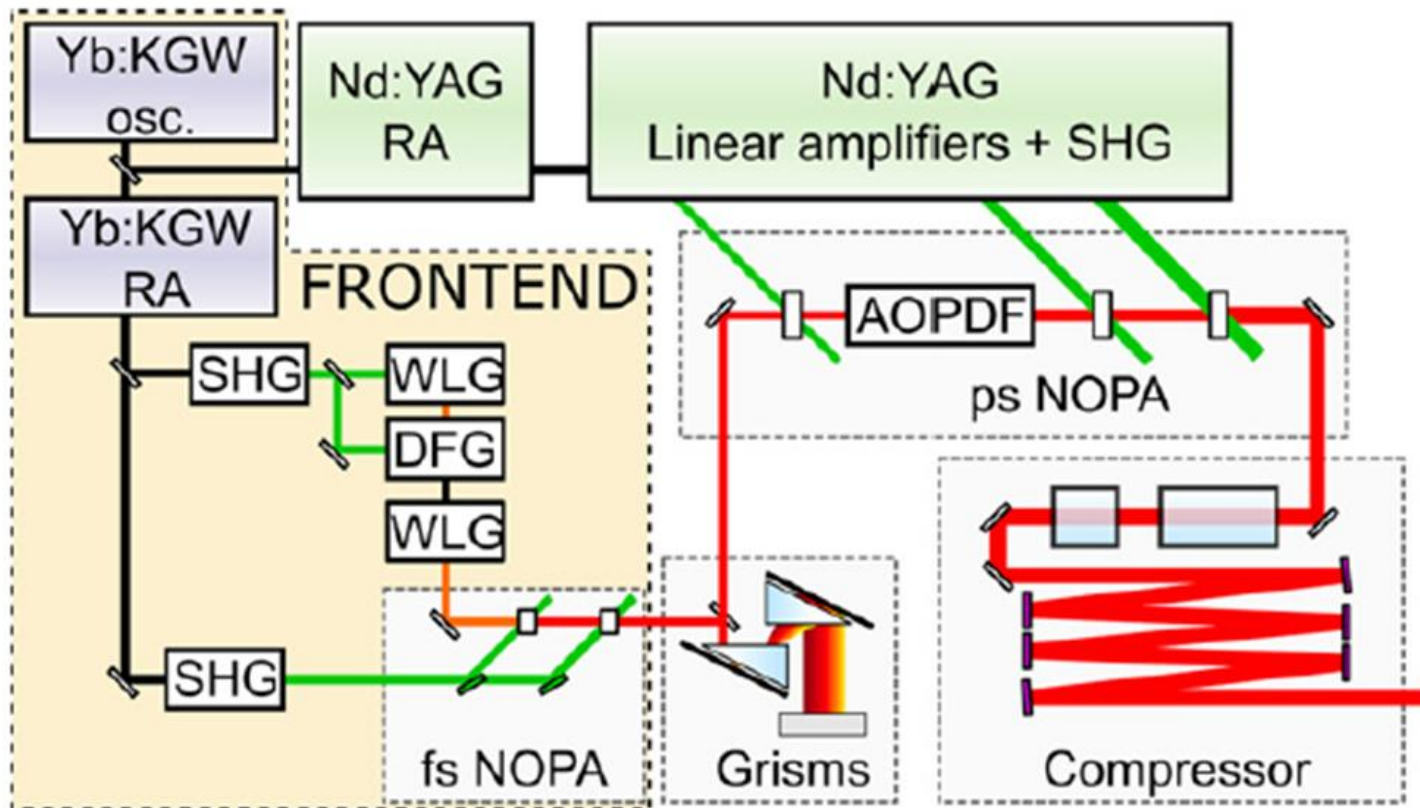
Flashed lamp pumped amplifier. Image from Hatae et al., Rev Sci Instrum **83**, 10E344 (2012)

Building blocks



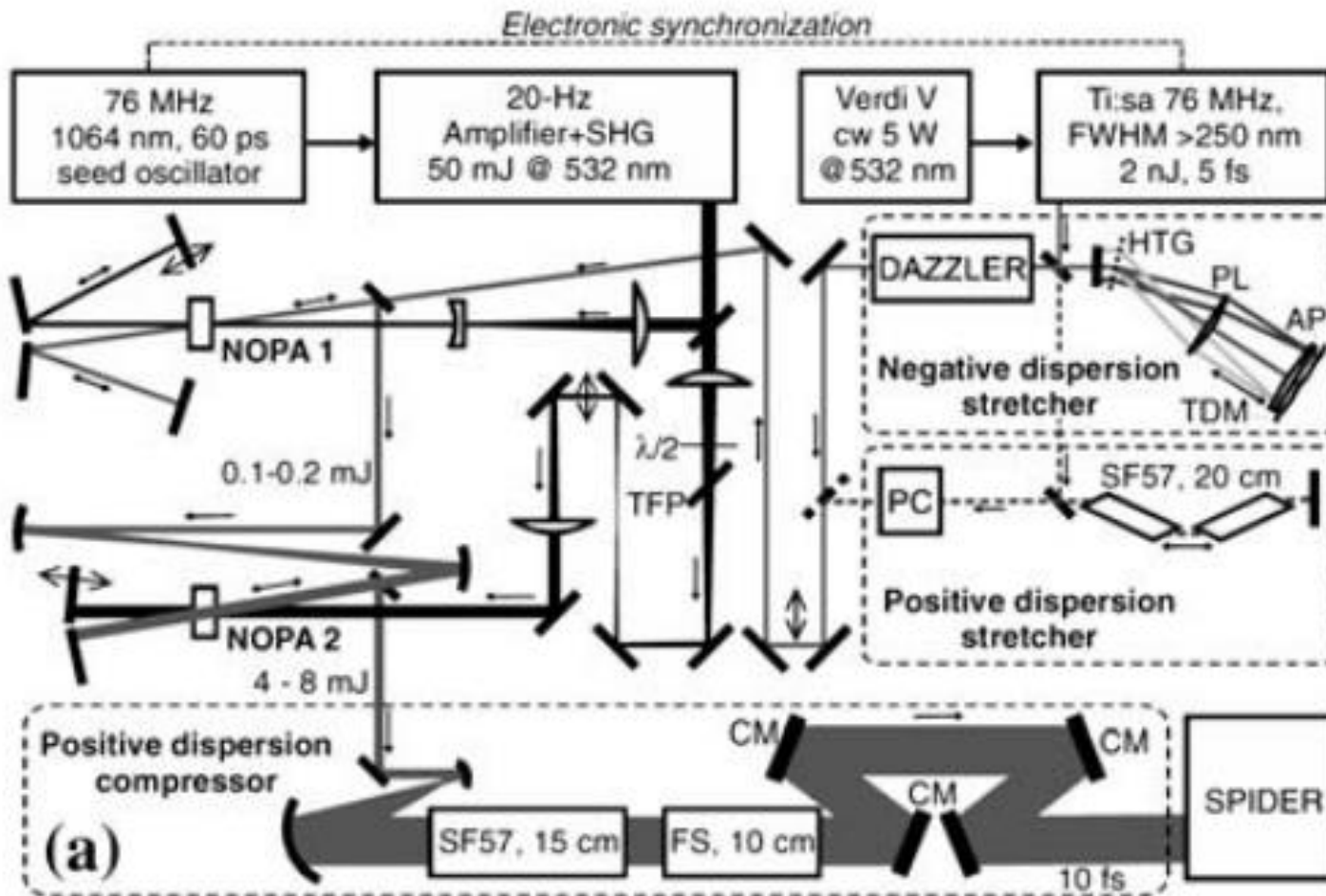
Synchronization

- Optical synchronization
 - Pulses derived from the same front end
 - Residual jitter: effects in amplifier chain

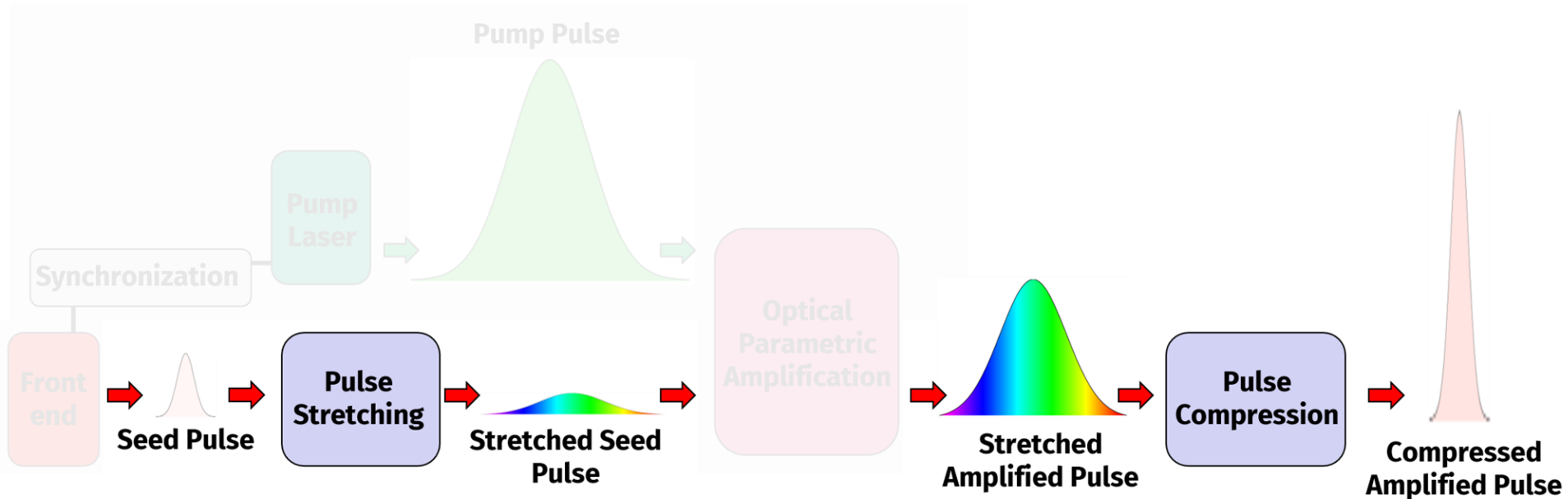


Synchronization

- Electronic synchronization
 - Oscillators locked to an external RF source
 - Residual jitter < 1ps



Building blocks



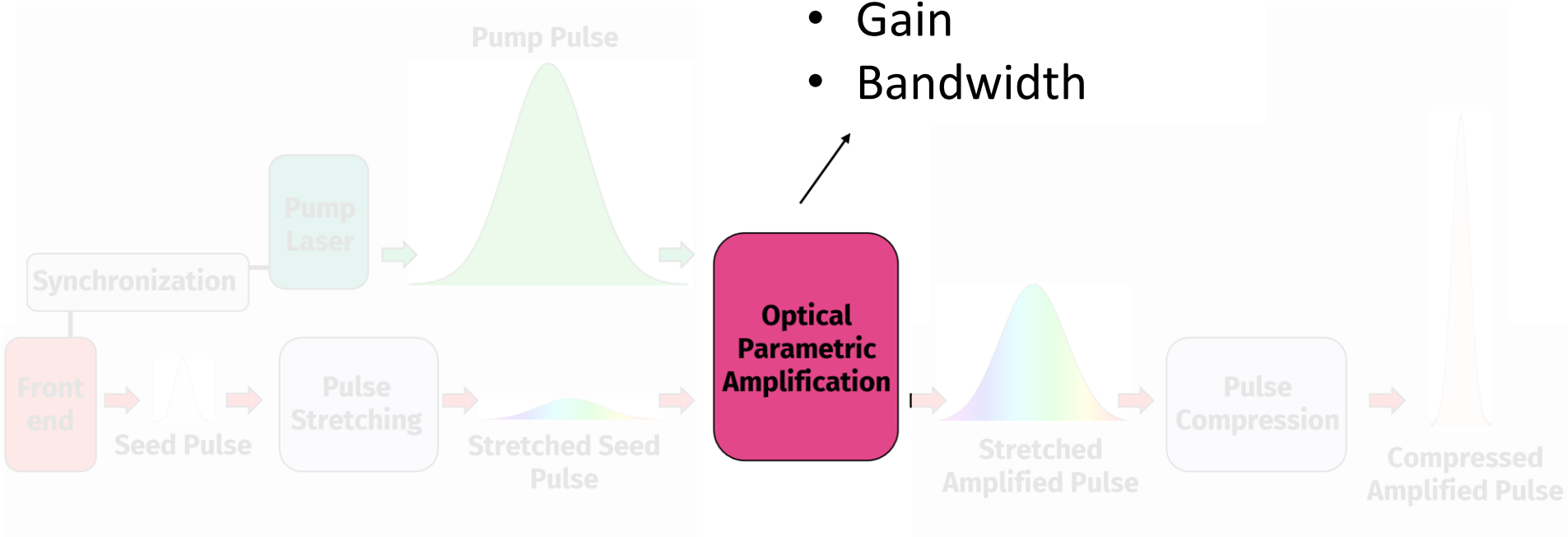
- Determines the chirp management

➤ ns, ~100ps pulses → grating stretcher / compressor

➤ Sub-10ps pulses → material dispersion / dispersive mirrors / prism pairs

Building blocks

- Spectral range
- Gain
- Bandwidth



OPCPA process

Coupled nonlinear differential equations for field envelopes: signal (s), idler (i) and pump (p)

$$\frac{\partial A_s}{\partial z} + \sum_{n=1}^{\infty} \frac{(-i)^{n-1}}{n!} k^{(n)} \frac{\partial^n A_s}{\partial t^n} = -i \frac{\chi^{(2)} \omega_s}{2n_s c} A_p A_i^* e^{-i\Delta k \cdot z}$$

$$\frac{\partial A_i}{\partial z} + \sum_{n=1}^{\infty} \frac{(-i)^{n-1}}{n!} k^{(n)} \frac{\partial^n A_i}{\partial t^n} = -i \frac{\chi^{(2)} \omega_i}{2n_i c} A_p A_s^* e^{-i\Delta k \cdot z}$$

$$\frac{\partial A_p}{\partial z} + \sum_{n=1}^{\infty} \frac{(-i)^{n-1}}{n!} k^{(n)} \frac{\partial^n A_p}{\partial t^n} = -i \frac{\chi^{(2)} \omega_p}{2n_p c} A_s A_i e^{i\Delta k \cdot z}$$

OPCPA process

PHYSICAL REVIEW

VOLUME 127, NUMBER 6

SEPTEMBER 15, 1962

Interactions between Light Waves in a Nonlinear Dielectric*

J. A. ARMSTRONG, N. BLOEMBERGEN, J. DUCUING,[†] AND P. S. PERSHAN

Division of Engineering and Applied Physics, Harvard University, Cambridge, Massachusetts

(Received April 16, 1962)

Solutions for monochromatic plane waves. No pump depletion, large gain:

$$I_s \approx I_s(0) \frac{1}{4} e^{2gz}$$

$$g = \sqrt{(\chi^{(2)})^2 \frac{\omega_s \omega_i I_p(0)}{2\epsilon_0 n_s n_i n_p c^3} - \left(\frac{\Delta k}{2}\right)^2}$$

Gain $\sim 10^3$ - 10^4 achievable

OPCPA process

PHYSICAL REVIEW

VOLUME 127, NUMBER 6

SEPTEMBER 15, 1962

Interactions between Light Waves in a Nonlinear Dielectric*

J. A. ARMSTRONG, N. BLOEMBERGEN, J. DUCUING,[†] AND P. S. PERSHAN

Division of Engineering and Applied Physics, Harvard University, Cambridge, Massachusetts

(Received April 16, 1962)

Phase of signal
independent of pump,
and dependent on
phase matching

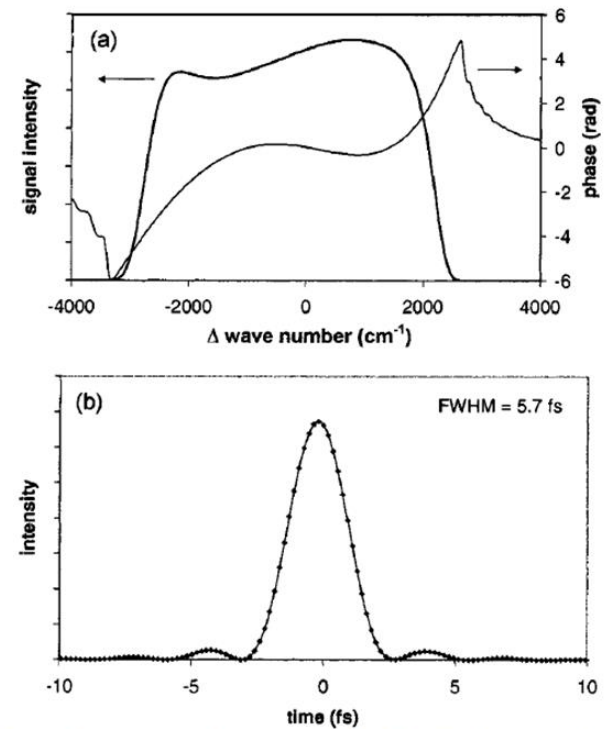


Fig. 9. Optimum performance of a BBO OPA of 1.5-mm length in noncollinear geometry for a 532-nm, 120-GW/cm² pump beam and 800-nm signal beam: (a) amplified spectral intensity and phase, and (b) pulse profile as determined by the Fourier transform of the spectral amplitude and phase after compensation for phase terms up to the quartic.

OPCPA process

PHYSICAL REVIEW

VOLUME 127, NUMBER 6

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62

VOLUME 88, NUMBER 13

PHYSICAL REVIEW LETTERS

1 APRIL 2002

Idler picks up phase difference between pump and signal

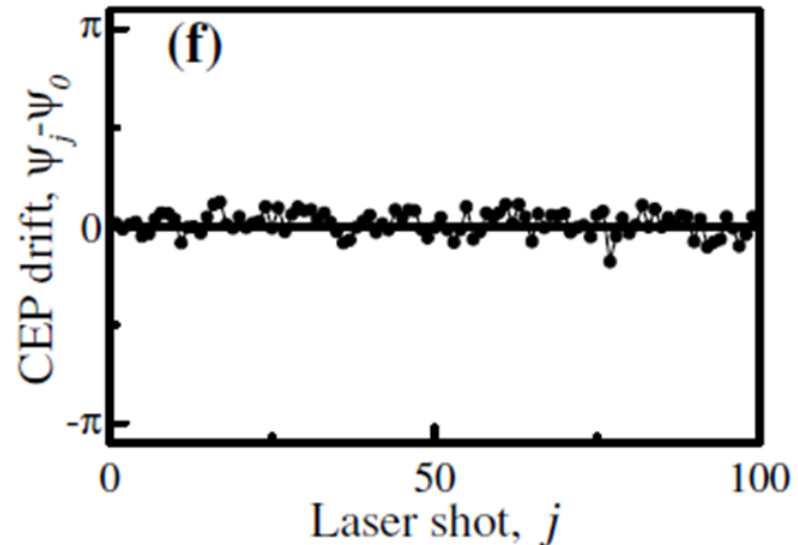
Controlling the Carrier-Envelope Phase of Ultrashort Light Pulses with Optical Parametric Amplifiers

Andrius Baltuška,* Takao Fuji, and Takayoshi Kobayashi

Department of Physics, Faculty of Science, University of Tokyo, 7-3-1 Hongo, Bunkyo, Tokyo 113-0033, Japan

(Received 1 November 2001; published 18 March 2002)

CEP-stable idler



OPCPA process

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Idler picks up phase
difference between
pump and signal



1672 Vol. 49, No. 7 / 1 April 2024 / *Optics Letters*

Letter

Optics Letters

Toward high-energy few-cycle optical vortices with minimized topological charge dispersion

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OPCPA process

Amplification rate is frequency dependent (through phase matching and time-dependent gain)

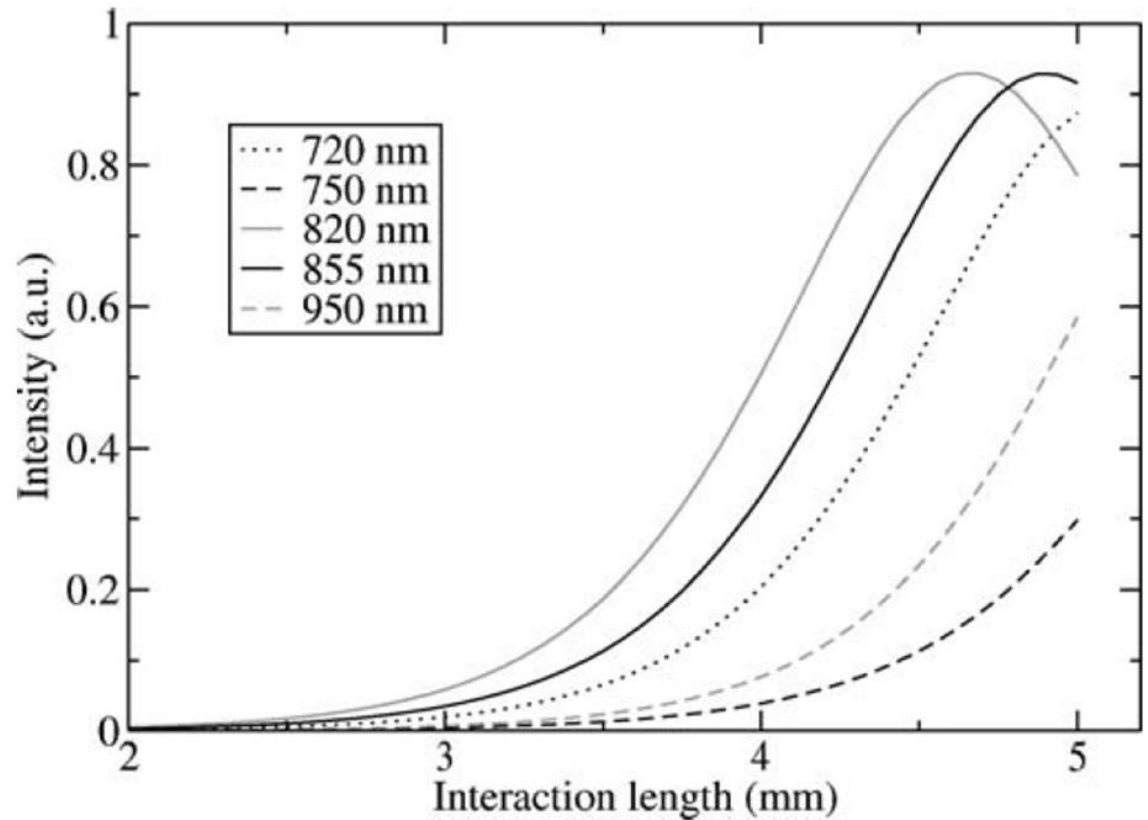


FIGURE 5 Wavelength-dependent gain saturation: Due to the differences in phase mismatch, the NOPCPA gain is frequency-dependent. Since all the frequencies are temporally separated in the chirped seed pulse, gain saturation and even backconversion independently occurs for every wavelength

OPCPA process

Amplification rate is frequency dependent (through phase matching and time-dependent gain)

Direction of energy flow can be reversed:
Back conversion

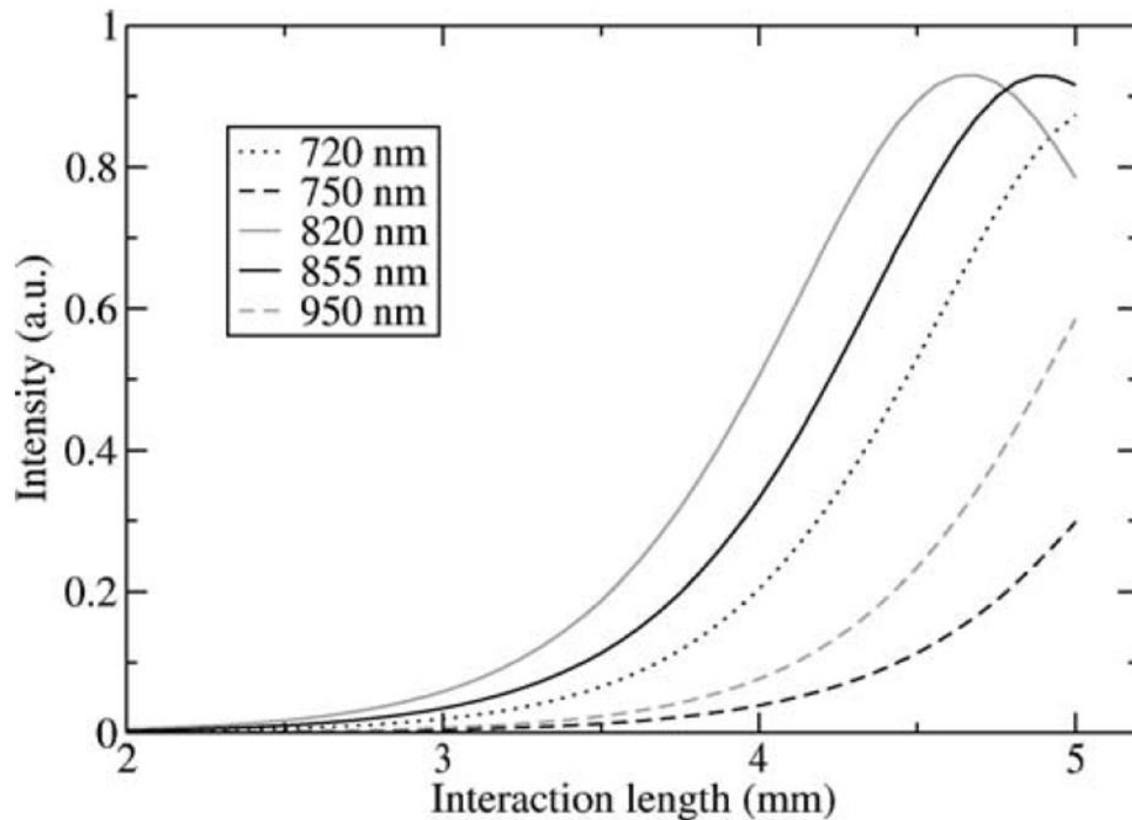


FIGURE 5 Wavelength-dependent gain saturation: Due to the differences in phase mismatch, the NOPCPA gain is frequency-dependent. Since all the frequencies are temporally separated in the chirped seed pulse, gain saturation and even backconversion independently occurs for every wavelength

$$\frac{\partial A_p}{\partial z} + \sum_{n=1}^{\infty} \frac{(-i)^{n-1}}{n!} k^{(n)} \frac{\partial^n A_p}{\partial t^n} = -i \frac{\chi^{(2)} \omega_p}{2n_p c} A_s A_i e^{i\Delta k \cdot z}$$

Witte et al., Appl. Phys. B **87**, 677 (2007)

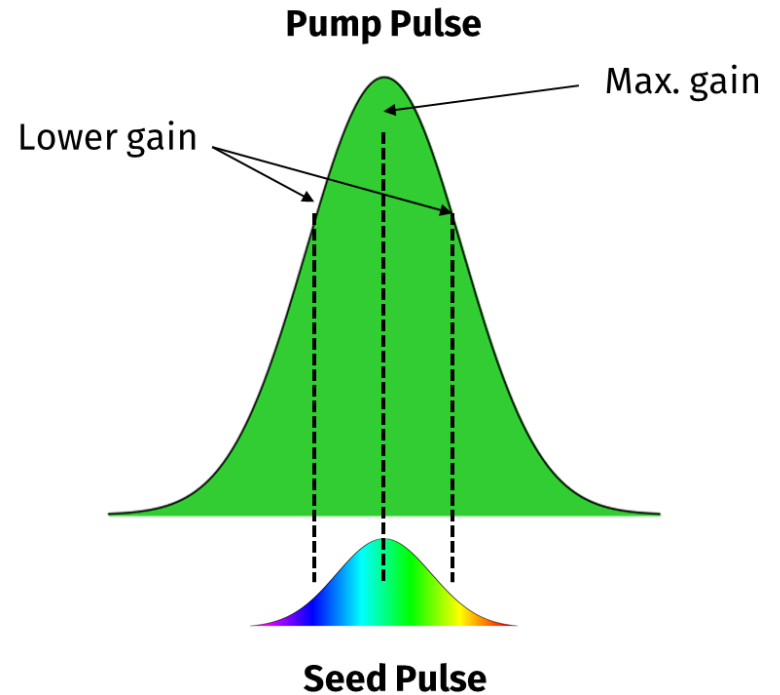
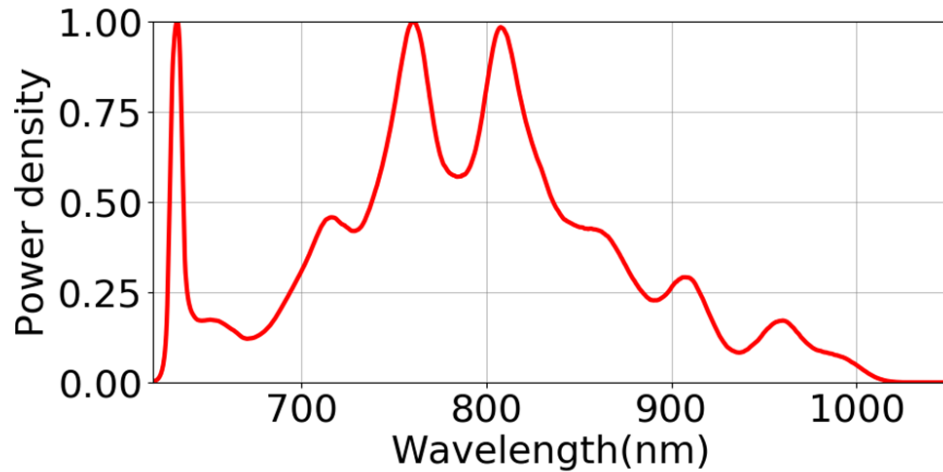
OPCPA process

Time dependent gain

Simulation with SISYFOS (G. Arisholm, Forsvarets forskningsinstitutt, Norway)

Seed: $<6\text{fs TL}, 800\text{fs}^2 (>500\text{fs}), 2.5\text{nJ}$

Pump: $1\text{ps}, 515\text{nm}, 100\text{GW}/\text{cm}^2, \text{BBO } 2.5\text{ mm}$



OPCPA process

Time dependent gain

Simulation with SISYFOS (G. Arisholm, Forsvarets forskningsinstitut, Norway)

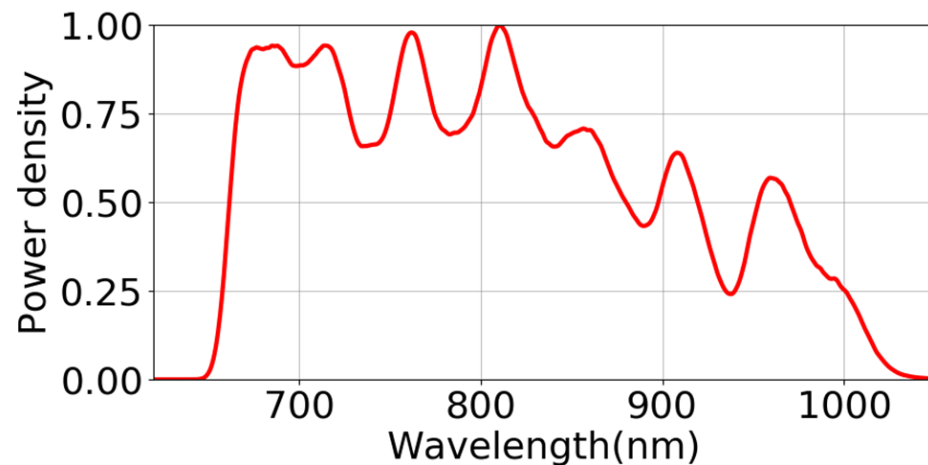
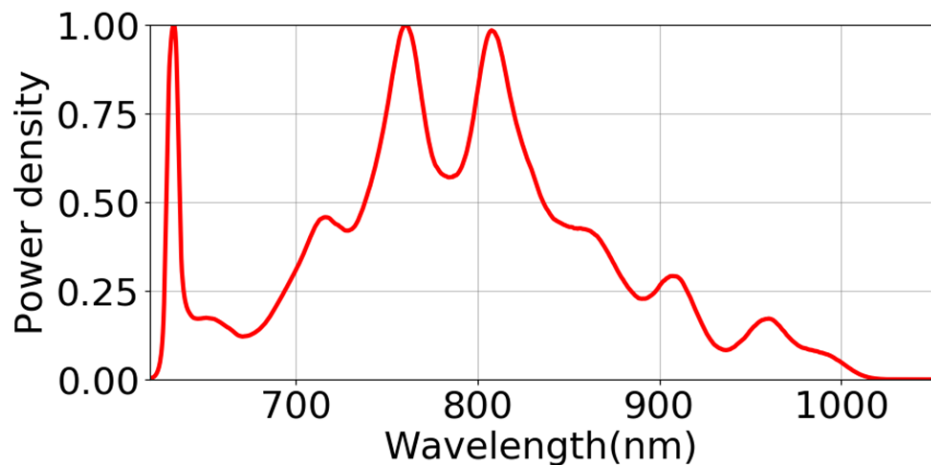
Seed: <6fs TL, 800fs²(>500fs), 2.5nJ

Pump: 1ps, 515nm, 100GW/cm², BBO 2.5 mm

Pump to
signal

$\eta \approx 17.3\%$

Saturation and back-conversion reshape spectrum



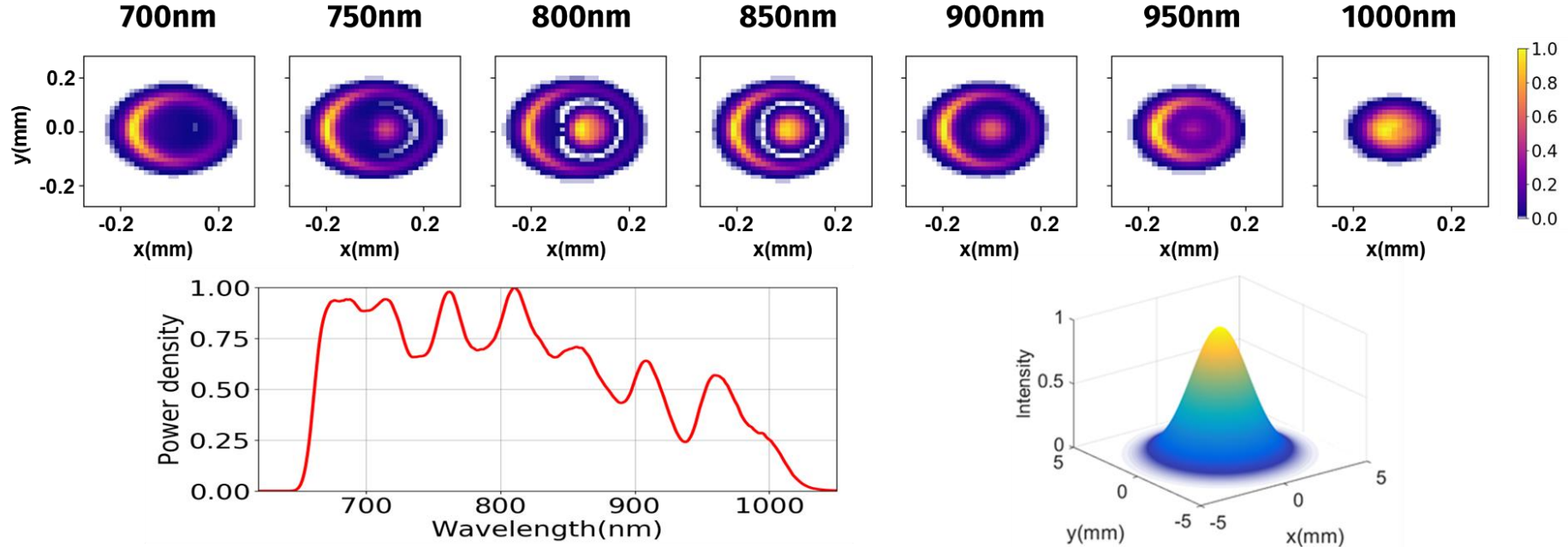
OPCPA process

Spatially dependent gain

Pump to
signal

$$\eta \approx 17.3\%$$

Saturation and back-conversion distort spatial distribution



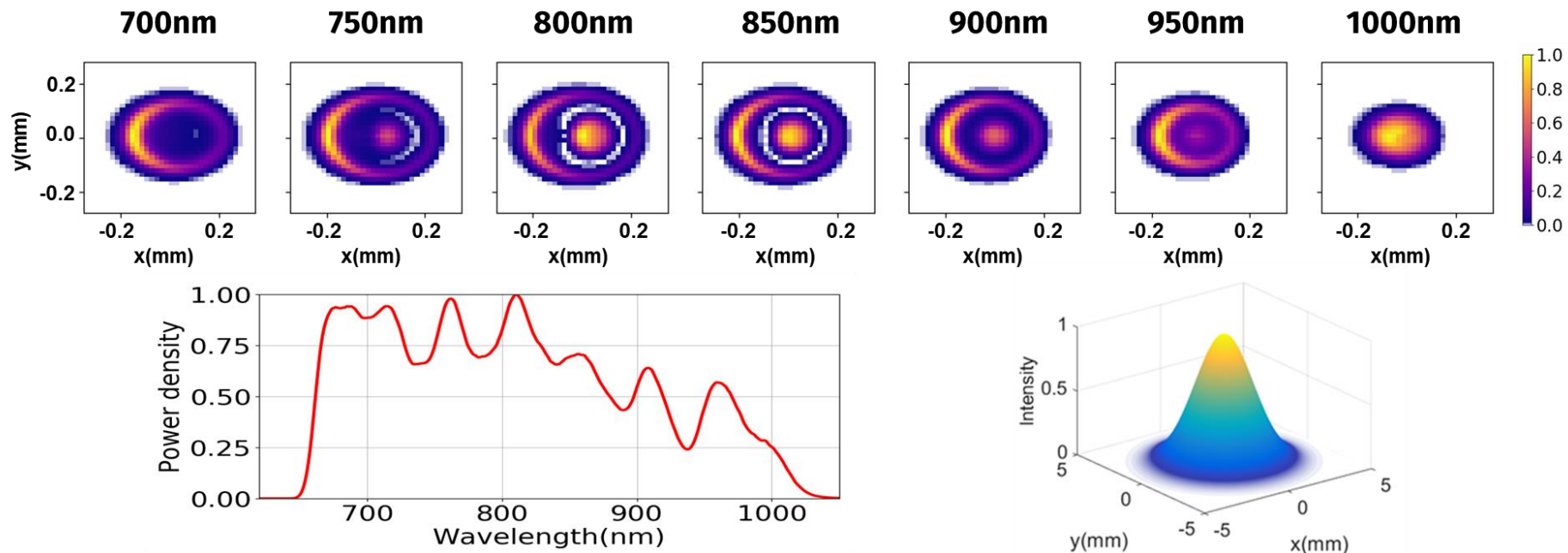
OPCPA process

Spatially dependent gain

Pump to
signal

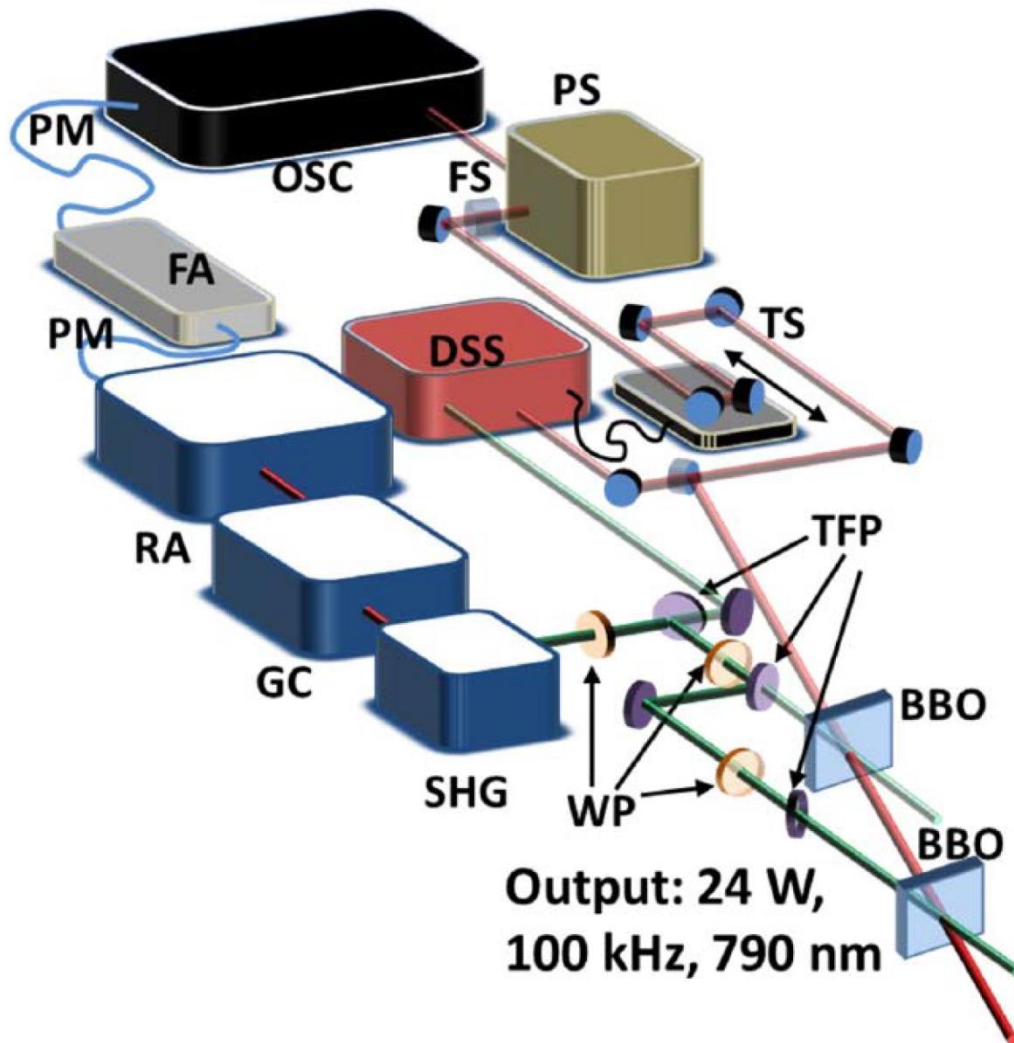
$\eta \approx 17.3\%$

Saturation and back-conversion distort spatial distribution



$E(x, y, z, t) \neq E_{sp}(x, y, z)E_{temp}(t) \rightarrow$ **Degradation of peak intensity**

Example: high rep. rate OPCPA at 800nm

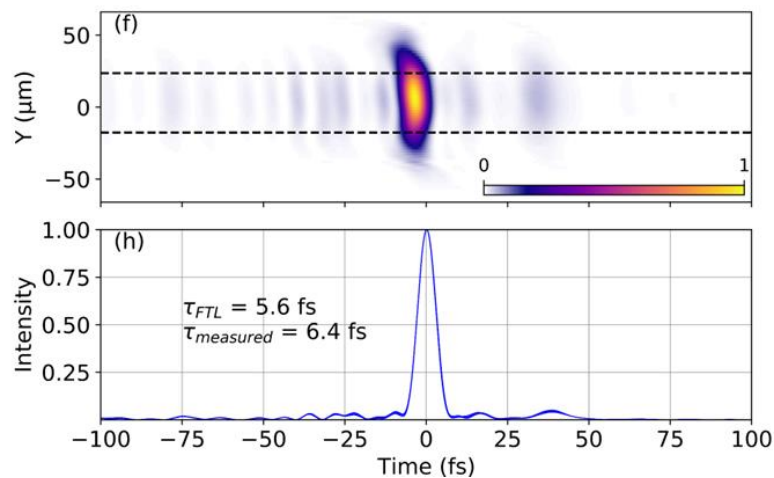
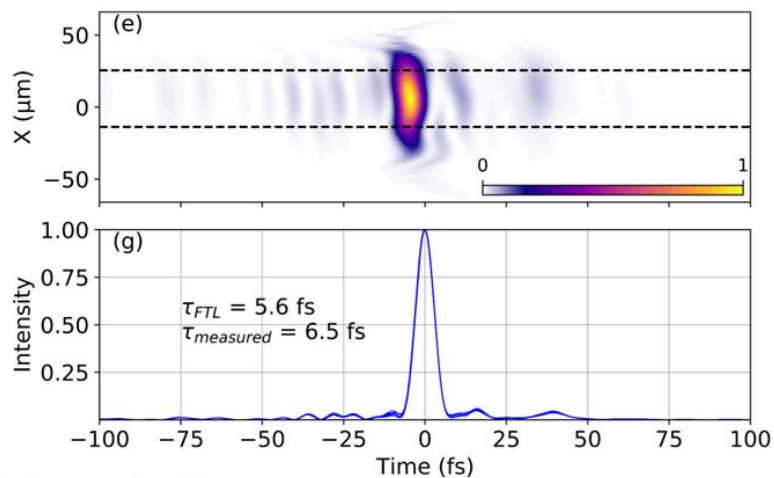
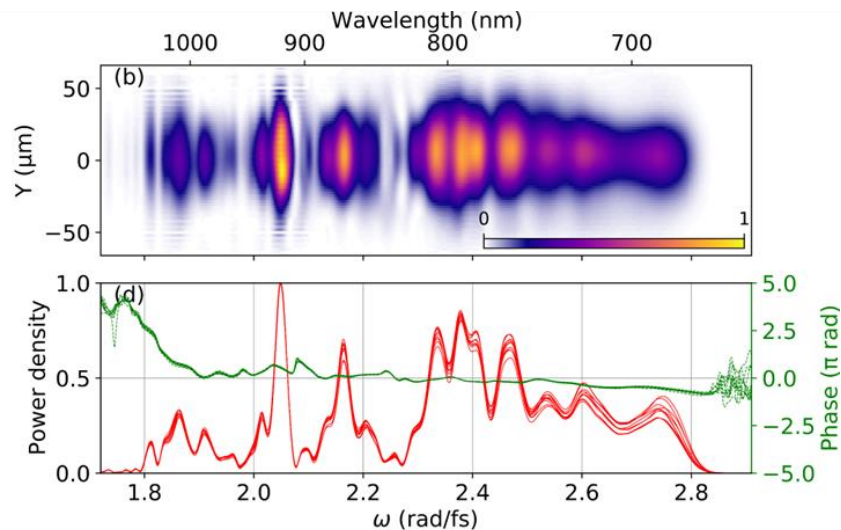
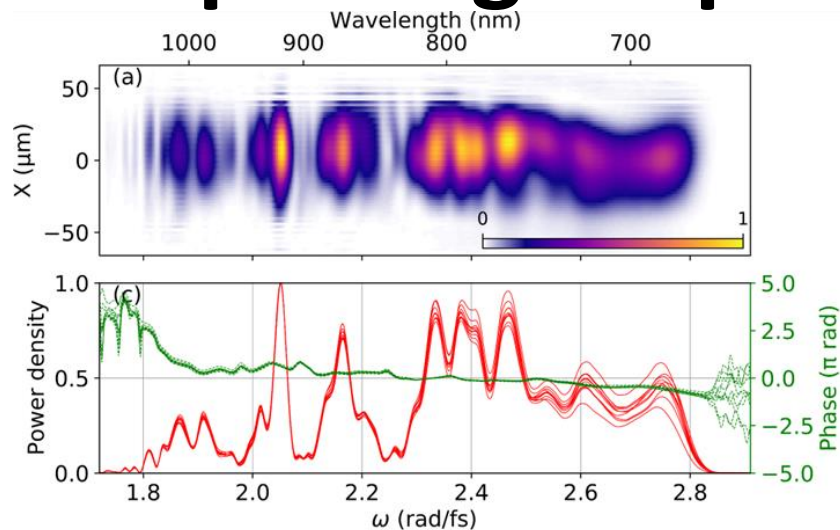


Seed: Ti:Sapphire oscillator
($<1\text{ nJ}$, $<6\text{ fs}$, 80 MHz)

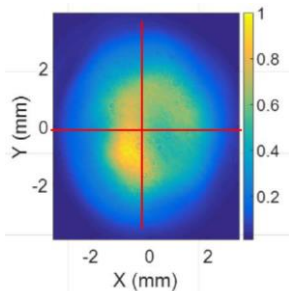
Pump: Yb:YAG thin-disk CPA
system + SHG
(1.2 mJ , 1 ps , 100 kHz , 515 nm)

OPCPA output: sub-7fs,
 $\approx 0.2\text{ mJ}$, 100 kHz , CEP-stable

Example: high rep. rate OPCPA at 800nm



SEA-F-SPIDER
characterization



F. Furch et al., J. Phys Photonics 4, 032001 (2022)

F. Furch et al., Optics Letters 42, 2495 (2017)

Example: high rep. rate OPCPA at 800nm

Generation and characterization of isolated attosecond pulses at 100 kHz repetition rate

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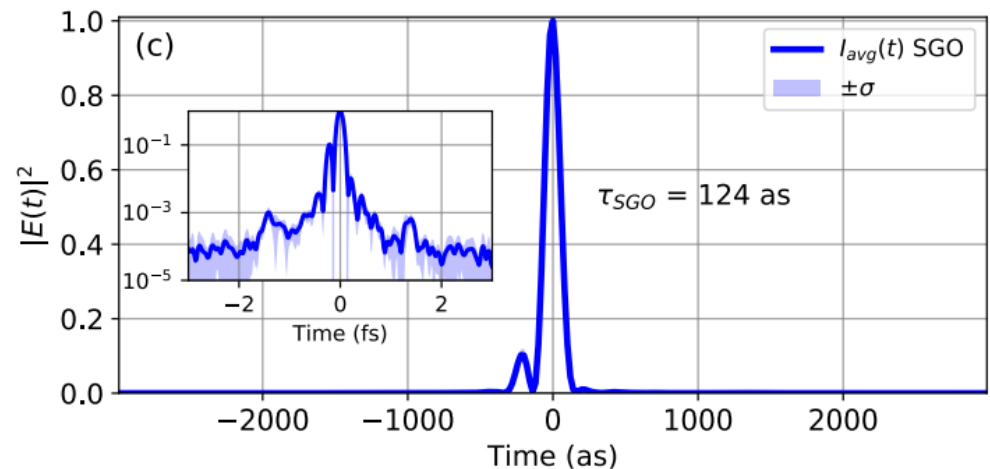
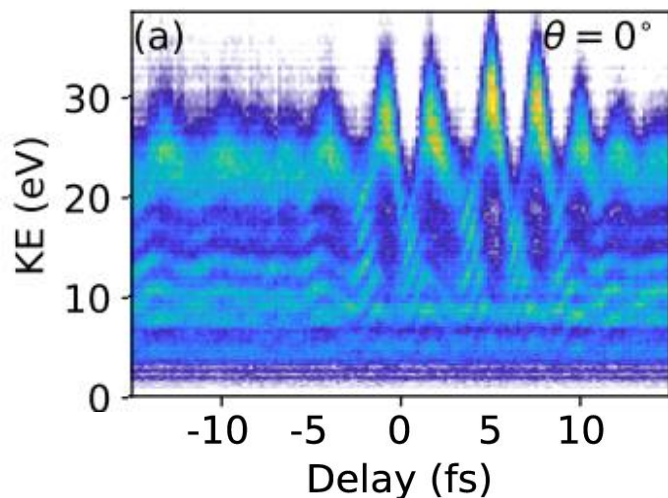
¹Max Born Institute for Nonlinear Optics and Short Pulse Spectroscopy, Max-Born-Strasse 2A, 12489 Berlin, Germany

²Department of Electrical and Computer Engineering, Colorado State University, Fort Collins, Colorado 80523, USA

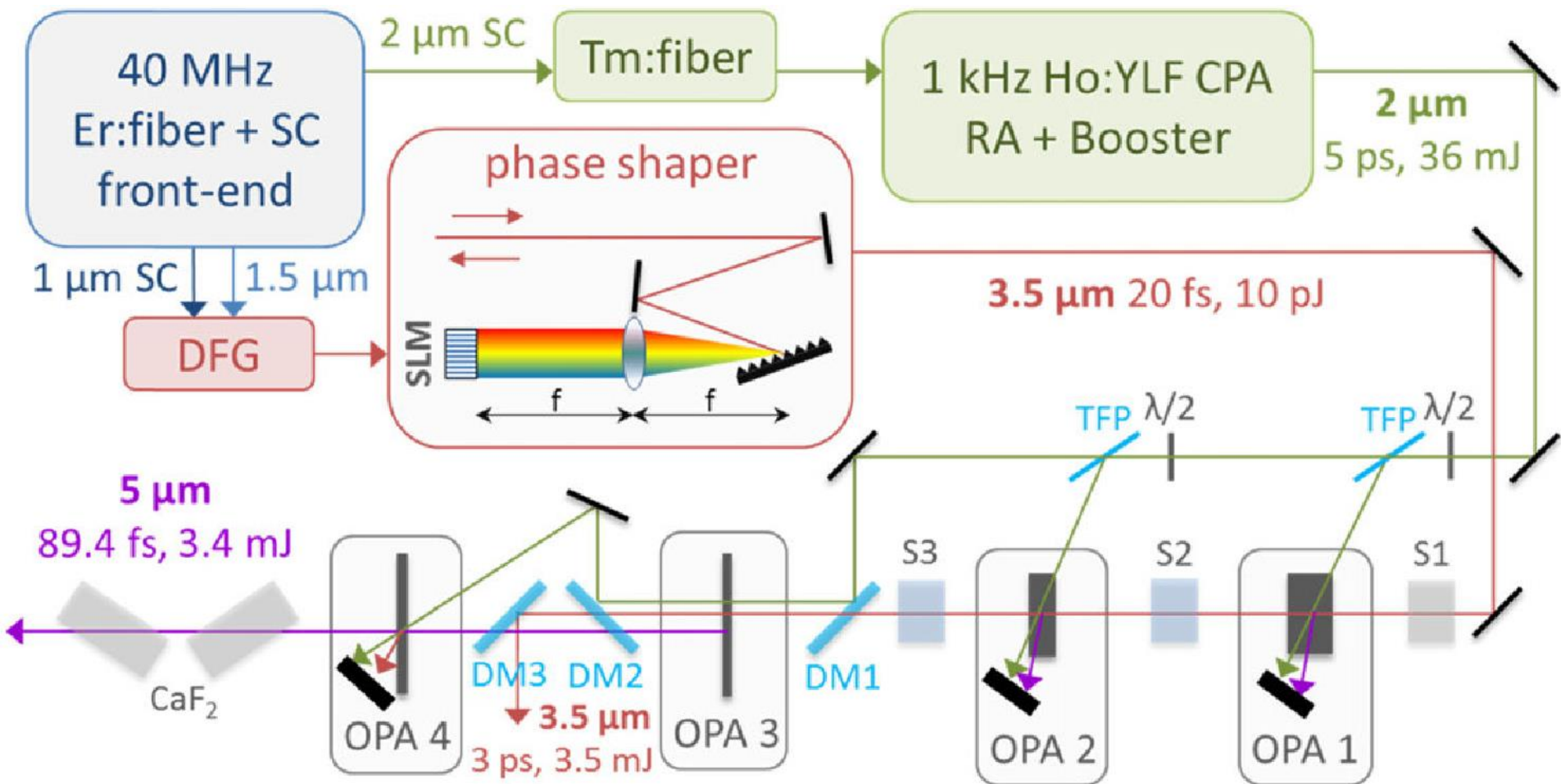
³e-mail: furch@mbi-berlin.de

*Corresponding author: tobias.witting@mbi-berlin.de

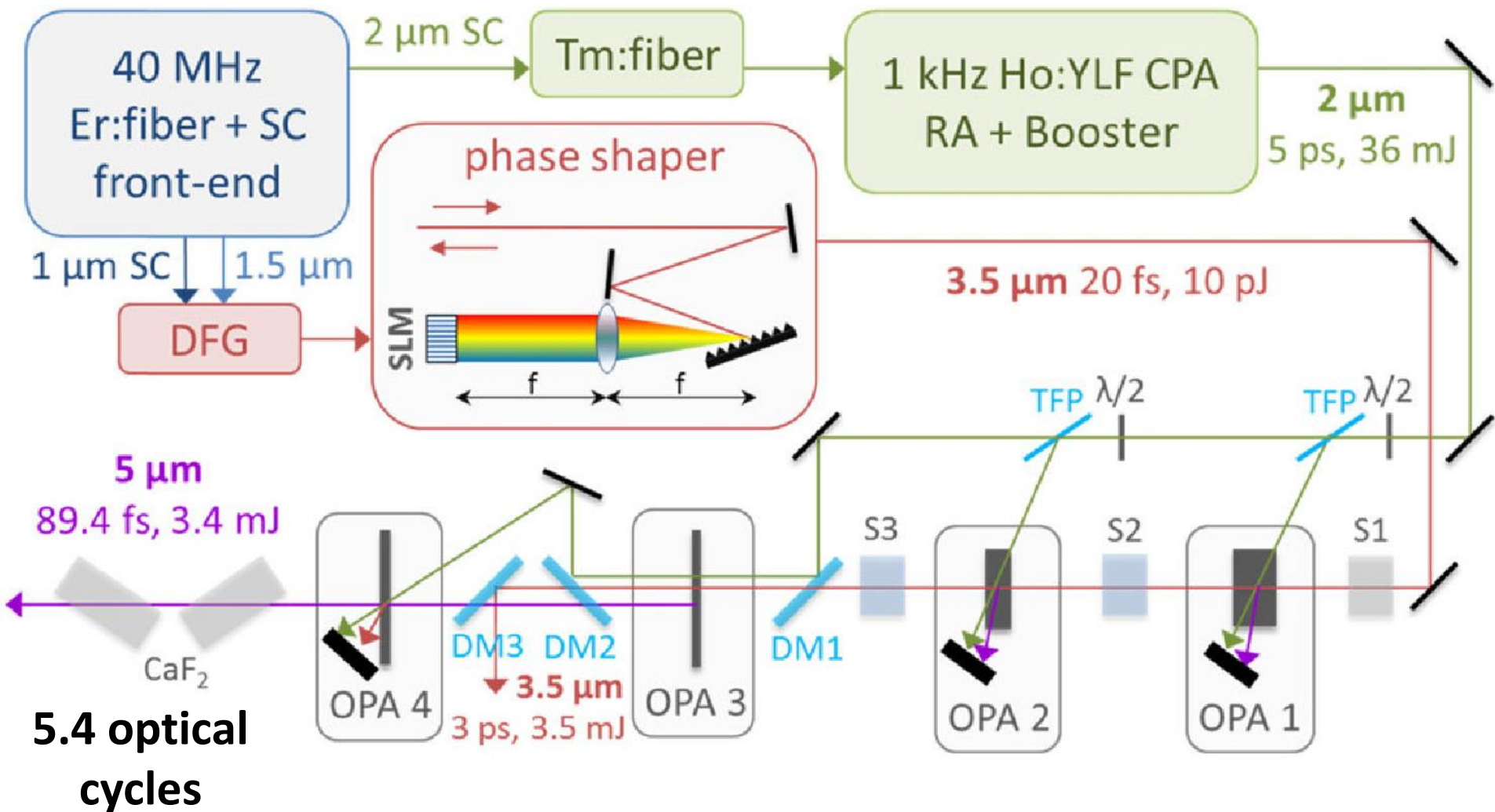
Received 15 September 2021; revised 30 November 2021; accepted 19 December 2021; published 28 January 2022



Example: high energy pulses at 5 microns



Example: high energy pulses at 5 microns



Useful materials for further reading:

C Manzoni and G Cerullo, Tutorial: Design criteria for ultrafast optical parametric amplifiers, *J. Opt.* 18 103501 (2016)

Hanieh Fattahi, et al., "Third-generation femtosecond technology," *Optica* 1, 45-63 (2014)

F. Furch et al., *J. Phys Photonics* 4, 032001 (2022)

Dubietis and Matijošius, *Opto-Electronic Advances* 6, 220046 (2023)