

# Estructura de la Materia 2

Clase 14 - Teoría

## Docentes

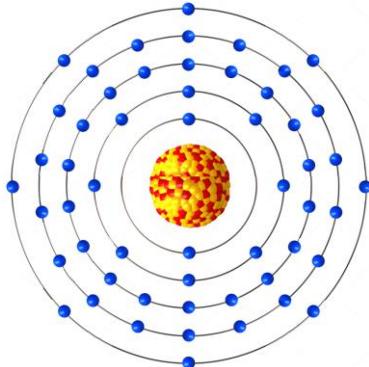
Gustavo Grinblat, Andrea Barral, Tomás Bortolin, Agustina Casafuz

Departamento de Física, FCEN, UBA – 2do Cuatrimestre, 2020

Web: <http://materias.df.uba.ar/edlm2a2020c2>

# Repaso

## Magnetismo en átomos aislados



$$\Delta\mathcal{H}_T = \mu_B \bar{H} \cdot (\bar{L} + g\bar{S}) + \frac{e^2}{8m_e c^2} \sum_i |\bar{H} \times \bar{r}_i|^2$$

Reglas de Hund

$$\Delta E_n = \mu_B \bar{H} \cdot \langle n | \bar{L} + g\bar{S} | n \rangle + \sum_{n' \neq n} \frac{|\langle n | \mu_B \bar{H} \cdot (\bar{L} + g\bar{S}) | n' \rangle|^2}{E_n - E_{n'}} + \frac{e^2}{8m_e c^2} \langle n | \sum_i |\bar{H} \times \bar{r}_i|^2 | n \rangle$$

Paramagnetismo de Curie

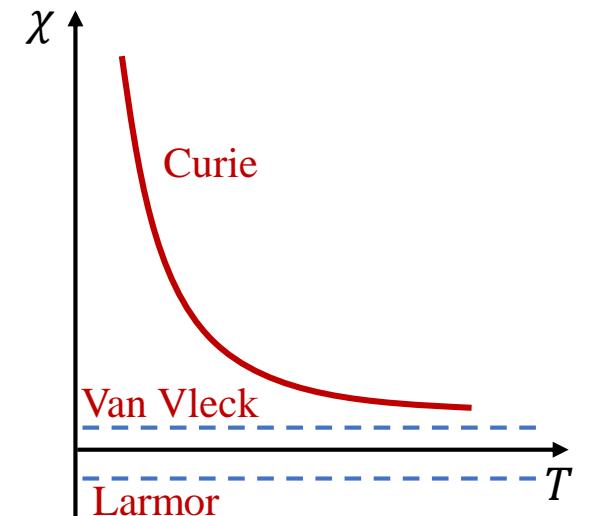
Paramagnetismo de Van Vleck

Diamagnetismo de Larmor

$$J \neq 0 \text{ (Paramagnetismo de Curie)} \longrightarrow \Delta E_n = g_L \mu_B \bar{H} \cdot \langle n | \bar{J} | n \rangle$$

$$\longrightarrow M = \frac{N}{V} \frac{g_L \mu_B}{2} \left[ (2J+1) \coth \left( \frac{\beta g_L \mu_B H (2J+1)}{2} \right) - \coth \left( \frac{\beta g_L \mu_B H}{2} \right) \right]$$

$$g_L \mu_B H \ll k_B T \longrightarrow \chi = \frac{\partial M}{\partial H} = \frac{N}{V} \frac{(g_L \mu_B)^2}{3k_B T} J(J+1)$$



# Repasso

---

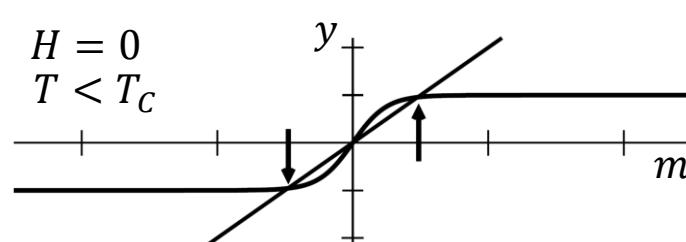
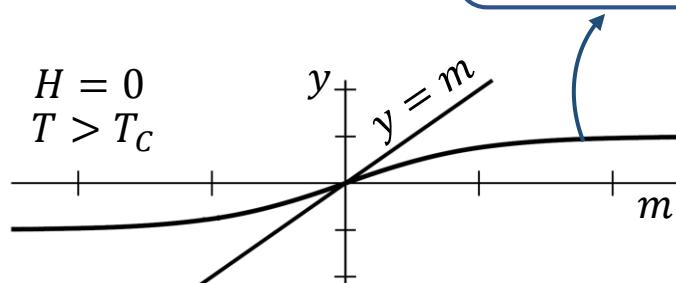
## Interacción entre espines: Orden ferromagnético

$$\mathcal{H} = -\frac{J_{int}}{2} \sum_{\langle i,j \rangle} \bar{S}_i \cdot \bar{S}_j + g_L \mu_B \bar{H} \cdot \sum_i \bar{S}_i, \quad J_{int} > 0$$

→ Para el sitio  $i$ :  $\mathcal{H}_i = \bar{S}_i \cdot \left[ -J_{int} \sum_j \langle \bar{S}_j \rangle + g_L \mu_B \bar{H} \right] = \bar{S}_i \cdot g_L \mu_B H_{ef}$

→  $g_L \mu_B H_{ef} \ll k_B T$  →  $\chi = \frac{N}{V} \frac{(g_L \mu_B)^2}{3k_B(T - T_C)} S(S+1)$ , con  $T_C = \frac{zJ_{int}S(S+1)}{3k_B}$  (Ley de Curie-Weiss)

A campo nulo:  $m = \frac{g_L \mu_B}{2} \left[ (2S+1) \coth \left( \frac{\beta(2S+1)zJ_{int}}{2g_L \mu_B} m \right) - \coth \left( \frac{\beta zJ_{int}}{2g_L \mu_B} m \right) \right]$

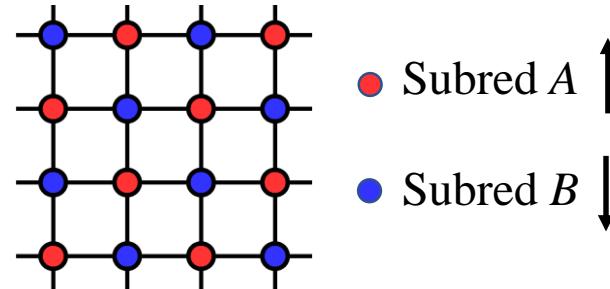


!Magnetización espontánea cuando  $T < T_c$ !

# Magnetismo: Antiferromagnetismo

Antiferromagneto: Interacción a primeros vecinos

$$\mathcal{H} = -\frac{J_{int}}{2} \sum_{\langle i,j \rangle} \bar{S}_i \cdot \bar{S}_j + g_L \mu_B \bar{H} \cdot \sum_i \bar{S}_i, \quad J_{int} < 0$$



Campo medio

Para el sitio  $i,A$ :  $\mathcal{H}_{i,A} = \bar{S}_{i,A} \cdot [-J_{int} z \langle \bar{S}_{j,B} \rangle + g_L \mu_B \bar{H}] = \bar{S}_{i,A} \cdot g_L \mu_B \bar{H}_{ef,B}$ ;  $\bar{H}_{ef,B} = \bar{H} - \frac{J_{int}}{g_L \mu_B} z \langle \bar{S}_{j,B} \rangle$

$\text{Nº de coordinación}$

Para el sitio  $i,B$ :  $\mathcal{H}_{i,B} = \bar{S}_{i,B} \cdot [-J_{int} z \langle \bar{S}_{j,A} \rangle + g_L \mu_B \bar{H}] = \bar{S}_{i,B} \cdot g_L \mu_B \bar{H}_{ef,A}$ ;  $\bar{H}_{ef,A} = \bar{H} - \frac{J_{int}}{g_L \mu_B} z \langle \bar{S}_{j,A} \rangle$

$$\bar{H} = H \hat{z} \rightarrow m_{A,B} = \frac{g_L \mu_B}{2} \left[ (2S+1) \coth \left( \frac{\beta g_L \mu_B H_{ef,B,A} (2S+1)}{2} \right) - \coth \left( \frac{\beta g_L \mu_B H_{ef,B,A}}{2} \right) \right]$$

$\cancel{g_L \mu_B H_{ef} \ll k_B T} \rightarrow m_{A,B} = \frac{(g_L \mu_B)^2}{3k_B T} S(S+1) H_{ef,B,A} = \frac{(g_L \mu_B)^2}{3k_B T} S(S+1) \left( H + \frac{z J_{int}}{(g_L \mu_B)^2} m_{B,A} \right)$

# Magnetismo: Antiferromagnetismo

## Antiferromagneto: Interacción a primeros vecinos

$$g_L \mu_B H_{ef} \ll k_B T: m_{A,B} = \frac{(g_L \mu_B)^2}{3k_B T} S(S+1) \left( H + \frac{zJ_{int}}{(g_L \mu_B)^2} m_{B,A} \right)$$

$$\rightarrow m_A = m_B = m \rightarrow m \left( 1 - \frac{zJ_{int}}{3k_B T} S(S+1) \right) = \frac{(g_L \mu_B)^2}{3k_B T} S(S+1) H$$

$$\rightarrow m = \frac{(g_L \mu_B)^2}{3k_B T - zJ_{int} S(S+1)} S(S+1) H$$

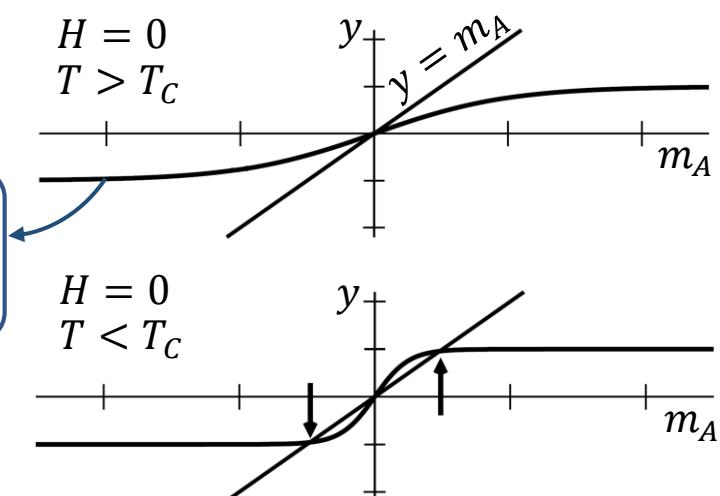
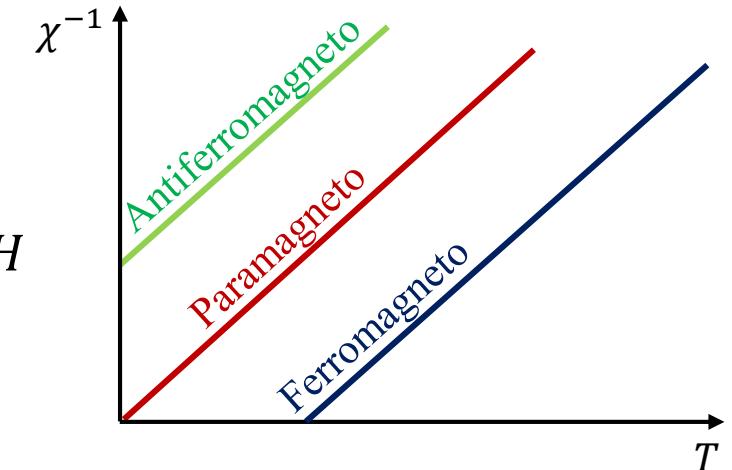
$$\rightarrow \chi = \frac{N}{V} \frac{(g_L \mu_B)^2}{3k_B(T + T_C)} S(S+1), \text{ con } T_C = -\frac{zJ_{int} S(S+1)}{3k_B}$$

Por otro lado, analizando a campo nulo:

$$m_{A,B} = \boxed{\frac{g_L \mu_B}{2} \left[ (2S+1) \coth \left( \frac{\beta(2S+1)zJ_{int}}{2g_L \mu_B} m_{B,A} \right) - \coth \left( \frac{\beta zJ_{int}}{2g_L \mu_B} m_{B,A} \right) \right]}$$

$$\rightarrow m_A = -m_B \text{ (coth es una función impar y } J_{int} < 0)$$

$\rightarrow$  Magnetización espontánea (para cada subred) cuando  $T < T_C$



# Magnetismo: Dominios e histéresis

## Anisotropía magnetocrystalina

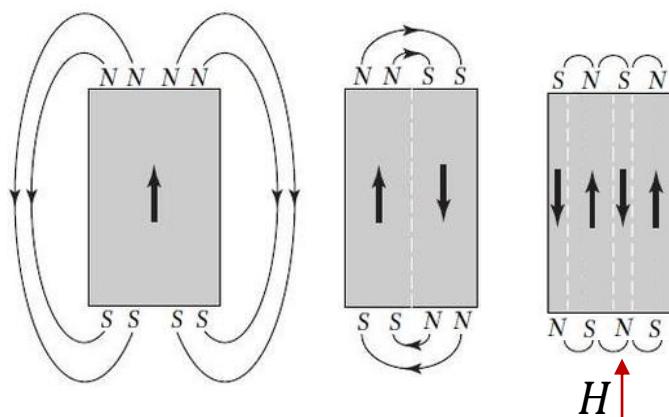
Debido al entorno anisotrópico de los átomos en una red cristalina, existen direcciones preferenciales para la orientación de los espines incluso en ausencia de campo magnético.

Una posibilidad sería:  $\mathcal{H} = -\frac{J_{int}}{2} \sum_{\langle i,j \rangle} \bar{S}_i \cdot \bar{S}_j - \kappa \sum_i (S_i^z)^2$  (Puede pensarse en un cristal tetragonal elongado en la dirección  $\hat{z}$ )

## Dominios magnéticos (ferromagneto)

Si bien la interacción dipolar magnética es despreciable frente a la interacción de intercambio, se torna apreciable cuando se tienen suficientes átomos, como ocurre en un material macroscópico ( $\sim 10^{23}$  átomos).

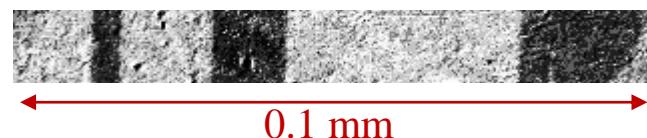
Esto sucede porque la interacción dipolar magnética es de largo alcance, mientras que la de intercambio es de corto alcance.



Energía magnetostática:  $\frac{1}{8\pi} \int H^2 dV$  Formar dominios magnéticos  
disminuye la energía total del sistema.

La región entre dominios recibe el nombre de “pared de dominio”.

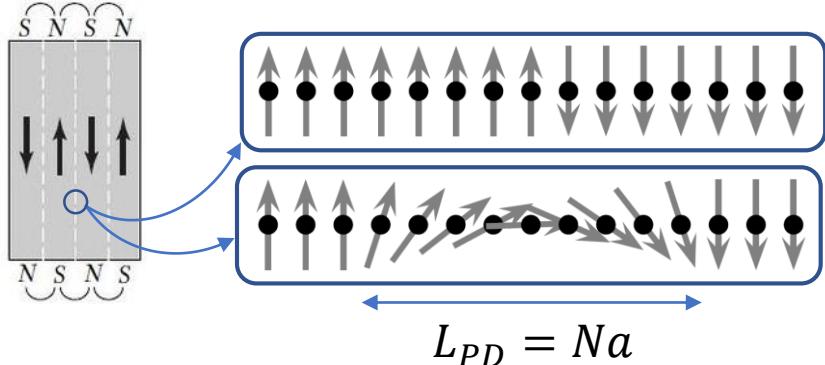
Aplicar un campo magnético externo desplaza paredes de dominio.



[https://en.wikipedia.org/wiki/Magnetic\\_domain](https://en.wikipedia.org/wiki/Magnetic_domain)

# Magnetismo: Dominios e histéresis

## Pared de dominio



$$E_{int} = -J_{int} \bar{S}_i \cdot \bar{S}_j = -J_{int} S^2 \cos(\theta_i - \theta_j) \approx -J_{int} S^2 \left( 1 - \frac{(\delta\theta)^2}{2} \right)$$

$$\rightarrow \delta E_{int} = J_{int} S^2 (\delta\theta)^2 / 2 = J_{int} S^2 (\pi/N)^2 / 2$$

$$\rightarrow E_{int}^{PD} = N J_{int} S^2 (\pi/N)^2 / 2$$

$$E_{ani} = -\kappa (S_i^z)^2 = -\kappa [S \cos \theta_i]^2 \rightarrow \delta E_{ani} = \kappa S^2 (1 - \cos^2 \theta_i) \rightarrow E_{ani}^{PD} = \sum_i \kappa S^2 \sin^2 \theta_i$$

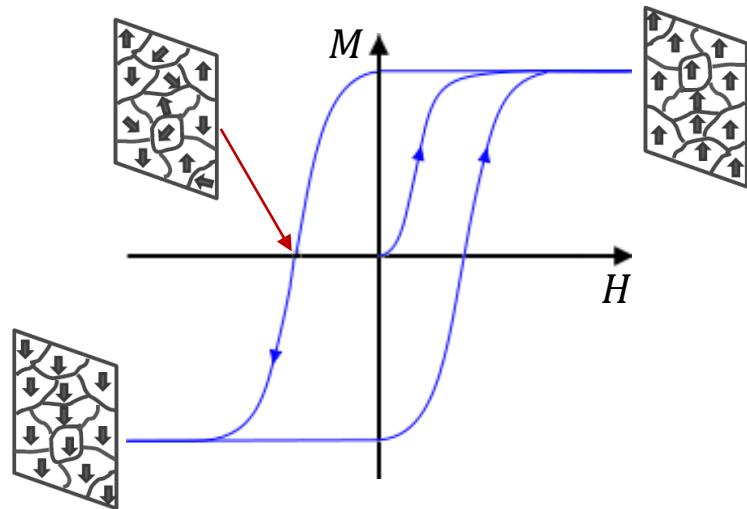
$$\rightarrow E_{ani}^{PD} = \kappa S^2 \int_0^\pi \sin^2 \theta d\theta \left( \frac{N}{\pi} \right) = \frac{N}{\pi} \kappa S^2 \int_0^\pi \frac{1 - \cos(2\theta)}{2} d\theta = \frac{N}{\pi} \kappa S^2 \left( \frac{\theta}{2} - \frac{\sin(2\theta)}{4} \right) \Big|_0^\pi = \frac{N \kappa S^2}{2}$$

$$\rightarrow E_{Total}^{PD} = \frac{J_{int} S^2 \pi^2}{2N} + \frac{N \kappa S^2}{2} \rightarrow \frac{\partial E_{Total}^{PD}}{\partial N} = 0 = -\frac{J_{int} S^2 \pi^2}{2N^2} + \frac{\kappa S^2}{2} \rightarrow \frac{J_{int} \pi^2}{N^2} = \kappa$$

$$\rightarrow N = \pi \sqrt{J_{int}/\kappa} \rightarrow E_{Total,min}^{PD} = \pi S^2 \sqrt{J_{int}\kappa} \quad (\text{Notar que si un cristal es extremadamente pequeño, puede resultar conveniente formar un monodominio})$$

# Magnetismo: Dominios e histéresis

## Histéresis en ferromagnéticos



Para un grano cristalino:  $\frac{\Delta E}{V} = g_L \mu_B \rho \bar{S} \cdot \bar{H} - \kappa \rho S_{eje-grano}^2 + int$

Densidad de espines  
cte

$$\rightarrow \frac{\Delta E}{V} = -\bar{M} \cdot \bar{H} - \kappa \frac{M_{eje-grano}^2}{(g_L \mu_B)^2 \rho} = -\bar{M} \cdot \bar{H} - \kappa' M_{eje-grano}^2$$

$$\bar{H} = H \hat{z} \rightarrow \frac{\Delta E}{V} = -|M||H|\cos\theta - \kappa' |M|^2 \cos^2(\theta - \phi)$$

$= 0$     $\hat{z}$   
 $\phi$     $Eje grano$

→  $\Delta E$  es mínimo cuando  $\bar{M}$  y  $\bar{H}$  son paralelos, y puede existir un mínimo relativo cuando son antiparalelos.

$$\frac{\partial(\Delta E)}{\partial \theta} = 0 \rightarrow -|M||H|\sin\theta - 2\kappa' |M|^2 \cos\theta \sin\theta = 0 \rightarrow \theta = 0, \pi, \arcsin\left(-\frac{|H|}{2\kappa' |M|}\right)$$

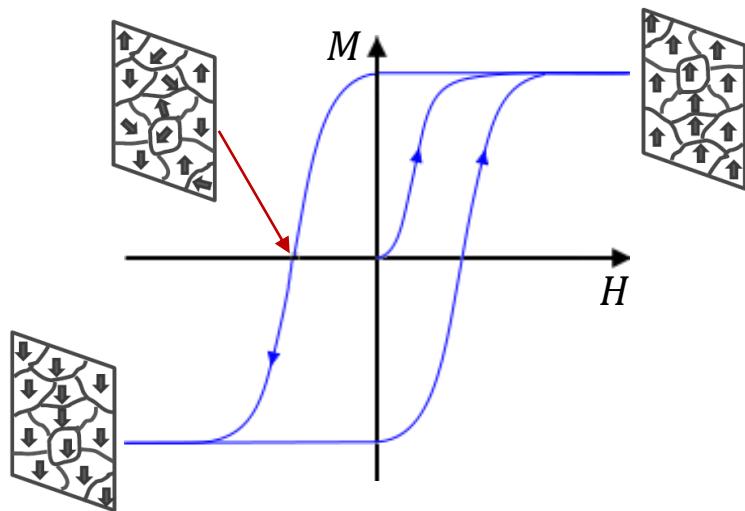
→  $H_{crítico} = 2\kappa' |M|$  → Si  $|H| < H_{crítico}$  existe un mínimo relativo, además del mínimo absoluto

Si  $\bar{M}$  apunta en la dirección  $-\hat{z}$  y se aplica un campo en la dirección  $\hat{z}$

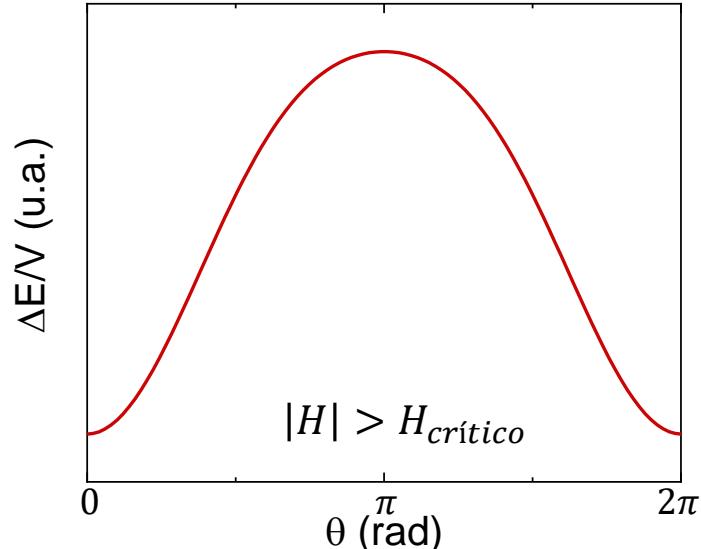
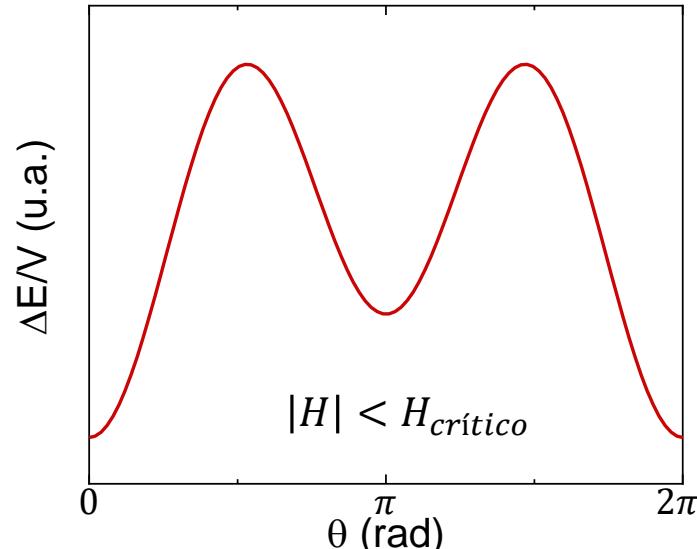
→ Debe ser suficientemente grande ( $> H_{crítico}$ ) para rotar a los espines.

# Magnetismo: Dominios e histéresis

## Histéresis en ferromagnéticos



$$\Delta E/V = -|M||H|\cos\theta - \kappa'|M|^2\cos^2\theta$$



→  $\Delta E$  es mínimo cuando  $\bar{M}$  y  $\bar{H}$  son paralelos, y puede existir un mínimo relativo cuando son antiparalelos.

$$\frac{\partial(\Delta E)}{\partial\theta} = 0 \rightarrow -|M||H|\sin\theta - 2\kappa'|M|^2\cos\theta\sin\theta = 0 \rightarrow \theta = 0, \pi, \arcsin\left(-\frac{|H|}{2\kappa'|M|}\right)$$

→  $H_{crítico} = 2\kappa'|M| \rightarrow$  Si  $|H| < H_{crítico}$  existe un mínimo relativo, además del mínimo absoluto

Si  $\bar{M}$  apunta en la dirección  $-\hat{z}$  y se aplica un campo en la dirección  $\hat{z}$

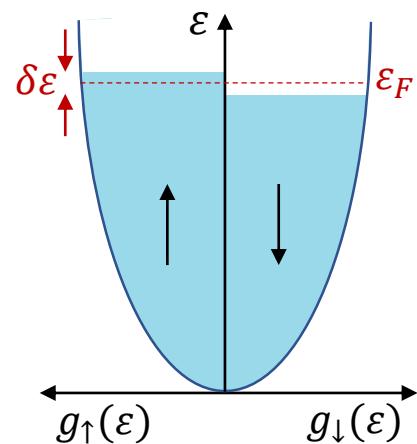
→ Debe ser suficientemente grande ( $> H_{crítico}$ ) para rotar a los espines.

# Magnetismo: Ferromagnetismo itinerante

## Magnetismo itinerante: Orden magnético a campo nulo

$$\mathcal{H} = \mathcal{H}_{\text{Tight-Binding}} + \mathcal{H}_{\text{Hubbard}} \quad \mathcal{H}_{\text{Hubbard}} = U \sum_i n_{i\uparrow} n_{i\downarrow}$$

¿Conviene tener  $\rho_\uparrow \neq \rho_\downarrow$ ?



$$\begin{cases} \varepsilon_{F,\uparrow} = \varepsilon_F + \delta\varepsilon/2 \\ \varepsilon_{F,\downarrow} = \varepsilon_F - \delta\varepsilon/2 \end{cases}$$

$$\rightarrow \rho_\uparrow - \rho_\downarrow =$$

$i$ : Corre sobre sitios de la red

$U$ : Costo energético por tener  $2e^-$  en el mismo sitio

$n_{i\uparrow}(n_{i\downarrow})$ : N° de  $e^-$  con espín  $\uparrow$  ( $\downarrow$ ) en el sitio  $i$

$$\int_0^{\varepsilon_F + \frac{\delta\varepsilon}{2}} \frac{g(\varepsilon)}{2} d\varepsilon - \int_0^{\varepsilon_F - \frac{\delta\varepsilon}{2}} \frac{g(\varepsilon)}{2} d\varepsilon \approx \delta\varepsilon \frac{g(\varepsilon_F)}{2}$$

Nº medio de  $e^-$  por sitio:  $x = (\rho_\uparrow - \rho_\downarrow)v$

$$\text{Aproximo: } Un_{i\uparrow}n_{i\downarrow} = \frac{U}{4}(n_{i\uparrow} + n_{i\downarrow})^2 - \frac{U}{4}(n_{i\uparrow} - n_{i\downarrow})^2 \approx \frac{U}{4}\overbrace{\langle n_{i\uparrow} + n_{i\downarrow} \rangle^2} - \frac{U}{4}\overbrace{\langle n_{i\uparrow} - n_{i\downarrow} \rangle^2}$$

$$\text{Magnetización: } M = -\frac{g}{2}\mu_B(\rho_\uparrow - \rho_\downarrow) = -\frac{g\mu_B\delta\varepsilon}{4}g(\varepsilon_F) \rightarrow \frac{E_{\text{Hubbard}}}{V} = \boxed{\frac{Ux^2}{4v} - \frac{U}{v}\left(\frac{Mv}{g\mu_B}\right)^2}$$

$$E \text{ cinética: } \frac{K}{V} = \int_0^{\varepsilon_F + \frac{\delta\varepsilon}{2}} \varepsilon \frac{g(\varepsilon)}{2} d\varepsilon + \int_0^{\varepsilon_F - \frac{\delta\varepsilon}{2}} \varepsilon \frac{g(\varepsilon)}{2} d\varepsilon = 2 \int_0^{\varepsilon_F} \varepsilon \frac{g(\varepsilon)}{2} d\varepsilon + \int_{\varepsilon_F}^{\varepsilon_F + \frac{\delta\varepsilon}{2}} \varepsilon \frac{g(\varepsilon)}{2} d\varepsilon - \int_{\varepsilon_F - \frac{\delta\varepsilon}{2}}^{\varepsilon_F} \varepsilon \frac{g(\varepsilon)}{2} d\varepsilon$$

$\stackrel{=} K(M=0)$

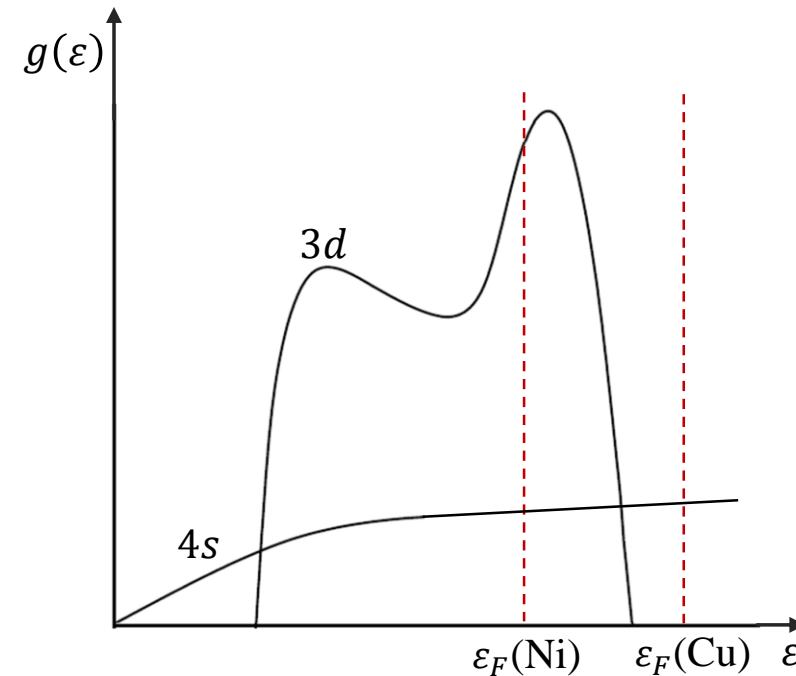
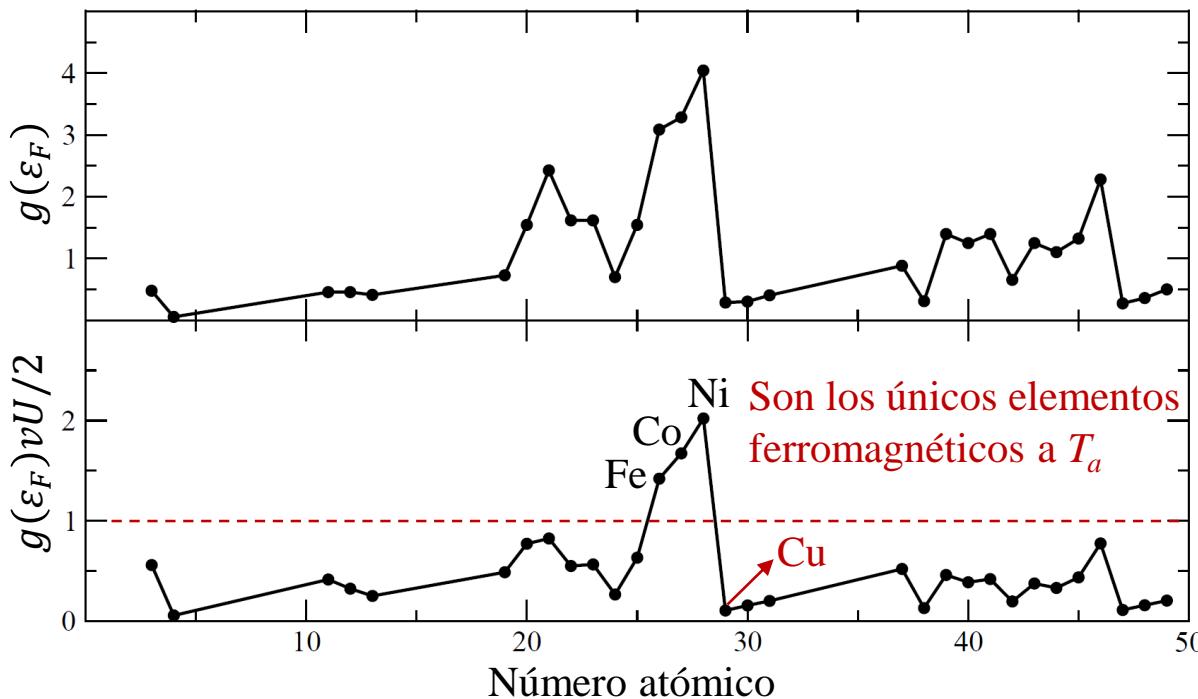
$$\approx \frac{K_0}{V} + \frac{g(\varepsilon_F)}{2} \left\{ \left[ \frac{(\varepsilon_F + \delta\varepsilon/2)^2}{2} - \frac{\varepsilon_F^2}{2} \right] - \left[ \frac{\varepsilon_F^2}{2} - \frac{(\varepsilon_F - \delta\varepsilon/2)^2}{2} \right] \right\} = \frac{K_0}{V} + \frac{g(\varepsilon_F)}{2} \left( \frac{\delta\varepsilon}{2} \right)^2 = \boxed{\frac{K_0}{V} + \frac{g(\varepsilon_F)}{2} \left( \frac{2M}{g\mu_B g(\varepsilon_F)} \right)^2}$$

# Magnetismo: Ferromagnetismo itinerante

## Magnetismo itinerante: Orden magnético a campo nulo

$$\frac{E}{V} = \frac{K + E_{Hubbard}}{V} = \frac{K_0 + \cancel{U_0}}{V} + \frac{1}{2g(\varepsilon_F)} \left( \frac{2M}{g\mu_B} \right)^2 - \frac{U}{v} \left( \frac{Mv}{g\mu_B} \right)^2 = \frac{\cancel{E_0}}{V} + \left( \frac{M}{g\mu_B} \right)^2 \left[ \frac{2}{g(\varepsilon_F)} - vU \right]$$

¿Cuándo conviene tener  $M \neq 0$ ?  $\rightarrow \frac{2}{g(\varepsilon_F)} - vU < 0 \rightarrow g(\varepsilon_F) > \frac{2}{vU} \rightarrow \frac{g(\varepsilon_F)vU}{2} > 1$  (Criterio de Stoner)

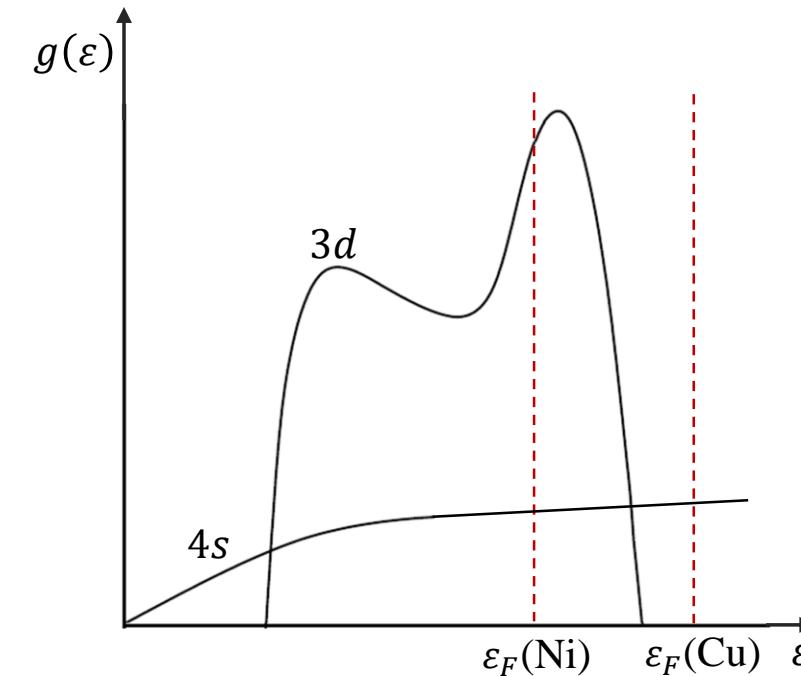
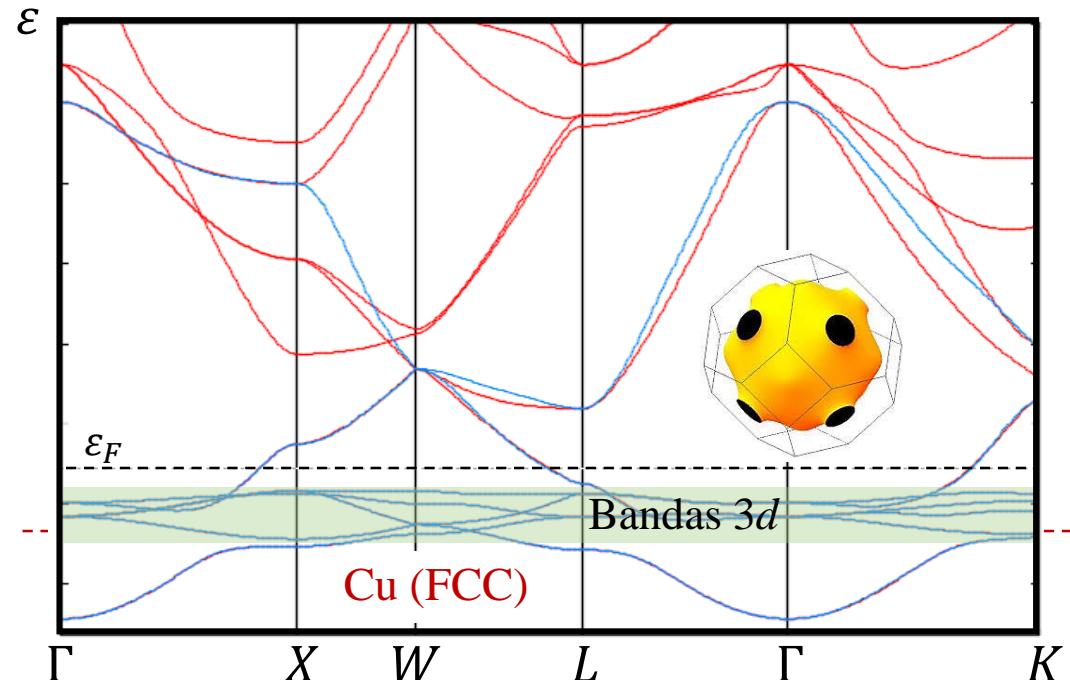


# Magnetismo: Ferromagnetismo itinerante

## Magnetismo itinerante: Orden magnético a campo nulo

$$\frac{E}{V} = \frac{K + E_{Hubbard}}{V} = \frac{K_0 + \cancel{U_0}}{V} + \frac{1}{2g(\varepsilon_F)} \left( \frac{2M}{g\mu_B} \right)^2 - \frac{U}{v} \left( \frac{Mv}{g\mu_B} \right)^2 = \cancel{\frac{E_0}{V}} + \left( \frac{M}{g\mu_B} \right)^2 \left[ \frac{2}{g(\varepsilon_F)} - vU \right]$$

¿Cuándo conviene tener  $M \neq 0$ ?  $\rightarrow \frac{2}{g(\varepsilon_F)} - vU < 0 \rightarrow g(\varepsilon_F) > \frac{2}{vU} \rightarrow \frac{g(\varepsilon_F)vU}{2} > 1$  (Criterio de Stoner)



# Magnetismo: La tabla periódica

# Magnetismo en la tabla periódica

Considerando todas las contribuciones, se obtiene un comportamiento dominante.

(e- de conducción: paramagnetismo de Pauli:  $\chi_{Pauli} = \mu_B^2 g(\varepsilon_F)$ ; diamagnetismo de Landau:  $\chi_{Landau} = -\frac{1}{3} \chi_{Pauli}$ )

<b>1</b>	<b>1.0080</b>	<b>H</b>	<b>18</b>
20.28			4.002602
17.81			2.416 (8.00 bar)
0.08988			0.1787
1s <sup>1</sup>			He
<b>2</b>			
<b>3</b> 6.97	<b>4</b> 9.012192	<b>Ferromagneto</b>	
1615 453.69	2745 1560 1.848	<b>Be</b> [He]2s <sup>1</sup>	
0.53			
<b>11</b> 1 22.89769	<b>12</b> 24.3050	<b>Diamagneto</b>	
1156.1 370.96	1363 922 1.738	<b>Mg</b> [Ne]3s <sup>2</sup>	
0.971			
<b>19</b> 1 39.0983	<b>20</b> 40.078	<b>Anti-ferromagneto</b>	
1033.1 336.4 0.862	21 44.955912 3 2.3,(4) 3103 1814 1.55	<b>Sc</b> 47.867 50.9415 [Ar]4s <sup>2</sup> 3d <sup>1</sup>	
	22 47.867	<b>Ti</b> 50.9416 2.3,(4),(5) 3560 2163+10 4.50 [Ar]4s <sup>2</sup> 3d <sup>2</sup>	
	23 50.9416	<b>V</b> 54.938045 2.3,(4,6,7) 3680 2163+10 5.80 [Ar]4s <sup>2</sup> 3d <sup>3</sup>	
	24 51.9961	<b>Cr</b> 55.845 2.3,(3) 2945 1517 7.43 [Ar]4s <sup>2</sup> 3d <sup>4</sup>	
	25 54.938045	<b>Mn</b> 58.933195 2.3,(2,3) 3134 1808 7.86 [Ar]4s <sup>2</sup> 3d <sup>5</sup>	
	26 55.845	<b>Fe</b> 58.6934 2.3,(2,3) 3020 1768 8.90 [Ar]4s <sup>2</sup> 3d <sup>6</sup>	
	27 58.933195	<b>Co</b> 63.546 1.2,(2) 3186 1726 8.90 [Ar]4s <sup>2</sup> 3d <sup>7</sup>	
	28 58.6934	<b>Ni</b> 65.538 2 1.2,(2) 2860 1825 8.96 [Ar]4s <sup>2</sup> 3d <sup>8</sup>	
	29 63.546	<b>Cu</b> 69.723 2 1.2,(2) 1180 692.73 7.14 [Ar]4s <sup>2</sup> 3d <sup>9</sup>	
	30 65.538	<b>Zn</b> 72.63 2 1.2,(4) 2477 302.93 5.907 [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	31 69.723	<b>Ga</b> 74.92160 1.2,(3,5) 3103 1210.6 5.72 [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	32 72.63	<b>Ge</b> 78.96 1.2,(4,6) 266 1.2,(3,5) [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	33 74.92160	<b>As</b> 79.904 1.2,(4,6) 886 (sub) 1.2,(3,5) [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	34 78.96	<b>Se</b> 83.798 1.2,(4,6) 490 1.2,(3,5) [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	35 79.904	<b>Kr</b> 83.798 0 1.2,(4,6) 331.93 1.2,(3,5) [Ar]4s <sup>2</sup> 3d <sup>10</sup>	
	36 83.798		
	37 85.4678	<b>Rb</b> 87.62 2 1.2,(4,6) 1657 1042 2.60 [Kr]5s <sup>1</sup>	
	38 87.62	<b>Sr</b> 88.90585 3 1.2,(4,6) 3611 1795+8 4.47 [Kr]5s <sup>2</sup>	
	39 88.90585	<b>Y</b> 91.224 4 1.2,(4,6) 4650 2125+2 6.40 [Kr]5s <sup>2</sup>	
	40 91.224	<b>Zr</b> 92.906 4 1.2,(4,6) 5015 2471+10 8.57 [Kr]5s <sup>2</sup>	
	41 92.906	<b>Nb</b> 95.96 3 1.2,(4,6) 4885 2890 10.20 [Kr]5s <sup>2</sup>	
	42 95.96	<b>Tc</b> 97.9072 4 1.2,(4,6,7) 5150 2445 11.5 [Kr]5s <sup>2</sup>	
	43 97.9072	<b>Mo</b> 101.07 2 1.2,(3,4,6,8) 4423 2607 12.20 [Kr]5s <sup>2</sup>	
	44 101.07	<b>Ru</b> 102.9050 2 1.2,(3,4) 3968 2239+3 12.40 [Kr]4d <sup>10</sup>	
	45 102.9050	<b>Rh</b> 106.42 2 1.2,(3,4) 3213 1825 12.02 [Kr]5s <sup>2</sup>	
	46 106.42	<b>Pd</b> 107.8682 1 1.2,(4) 2435 1235.08 10.50 [Kr]5s <sup>2</sup>	
	47 107.8682	<b>Ag</b> 112.411 1 1.2,(4) 1038 594.1 8.65 [Kr]5s <sup>2</sup>	
	48 112.411	<b>Cd</b> 118.710 2 1.2,(4) 2353 429.76 7.31 [Kr]5s <sup>2</sup>	
	49 118.710	<b>In</b> 121.760 2 1.2,(4) 2875 505.12 7.30 [Kr]5s <sup>2</sup>	
	50 121.760	<b>Sb</b> 127.60 2 1.2,(4,6) 1860 903.89 6.84 [Kr]5s <sup>2</sup>	
	51 127.60	<b>Te</b> 131.293 2 1.2,(4,6) 1263.1 722.7 6.24 [Kr]5s <sup>2</sup>	
	52 131.293		
	53 132.905452	<b>Cs</b> 137.327 3 1.2,(4,6) 2170 1936 3.51 [Xe]6s <sup>1</sup>	
	54 137.327	<b>Ba</b> 174.96668 3 1.2,(4,6) 3675 2875 9.85 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>5</sup>	
	55 174.96668	<b>Lu</b> 178.49 4 1.2,(4,6) 570+100 3269 13.2 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>5</sup>	
	56 178.49	<b>Ta</b> 180.94788 5 1.2,(4,6) 5933 3863+20 19.3 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>5</sup>	
	57 180.94788	<b>W</b> 183.84 2 1.2,(4,6,7) 5900 (est.) 3453 21.0 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>5</sup>	
	58 183.84	<b>Os</b> 186.207 2 1.2,(4,6,8) 5300 3237 22.4 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>5</sup>	
	59 186.207	<b>Ir</b> 192.217 2 1.2,(4,6,8) 4100 2045 21.45 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>7</sup>	
	60 192.217	<b>Pt</b> 195.084 2 1.2,(4,6,8) 4403 1337.58 19.32 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>9</sup>	
	61 195.084	<b>Au</b> 196.966569 1 1.2,(4,6,8) 628.73 232.48 13.546 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	62 196.966569	<b>Hg</b> 200.59 1 1.2,(4,6,8) 576.7 11.85 11.34 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	63 200.59	<b>Tl</b> 204.384 2 1.2,(4,6) 81 1.2,(4,6) 1730+10 2013 600.652 9.80 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	64 204.384	<b>Pb</b> 207.2 2 1.2,(4,6) 1833.5 544.5 9.4 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	65 207.2	<b>Bi</b> 208.98040 3 1.2,(4,6) 1235 527 9.4 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	66 208.98040	<b>Po</b> 208.3824 2 1.2,(4,6) 810 575 9.4 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	67 208.3824	<b>At</b> 209.9871 0 1.2,(4,6) 211 202 9.73 [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	68 209.9871		
	69 168.93421	<b>Rn</b> 213.054 2 1.2,(3) 1092 169 Yb [Xe]6s <sup>2</sup> 4f <sup>1</sup> 5d <sup>10</sup>	
	70 213.054		
	71 223.0197	<b>Fr</b> 226.0254 1 1.2,(4,6) 950 300 — [Rn]7s <sup>1</sup>	
	72 226.0254	<b>Ra</b> 227.02781 2 1.2,(4,6) 1413 973 6.0 [Rn]7s <sup>1</sup>	
	73 227.02781	<b>La</b> 138.90547 3 1.2,(4,6) 3706 1071+3 6.78 [Xe]6s <sup>2</sup> 5d <sup>1</sup>	
	74 138.90547	<b>Ce</b> 140.116 3 1.2,(4,6) 3793 1204 6.77 [Xe]6s <sup>2</sup> 4f <sup>1</sup>	
	75 140.116	<b>Pr</b> 144.242 3 1.2,(4,6) 3347 1294 7.00 [Xe]6s <sup>2</sup> 4f <sup>2</sup>	
	76 144.242	<b>Nd</b> 144.9127 3 1.2,(4,6) 3300 (est.) 1315 6.475 [Xe]6s <sup>2</sup> 4f <sup>3</sup>	
	77 144.9127	<b>Pm</b> 150.36 2 1.2,(4,6) 2067 1347 7.54 [Xe]6s <sup>2</sup> 4f <sup>4</sup>	
	78 150.36	<b>Sm</b> 151.964 2 1.2,(4,6) 1800 1095 7.895 [Xe]6s <sup>2</sup> 4f <sup>5</sup>	
	79 151.964	<b>Eu</b> 157.25 3 1.2,(4,6) 3546 1629 8.27 [Xe]6s <sup>2</sup> 4f <sup>6</sup>	
	80 157.25	<b>Gd</b> 158.925 3 1.2,(4,6) 3503 1747 8.80 [Xe]6s <sup>2</sup> 4f <sup>7</sup>	
	81 158.925	<b>Tb</b> 162.500 3 1.2,(4,6) 2840 1865 8.536 [Xe]6s <sup>2</sup> 4f <sup>8</sup>	
	82 162.500	<b>Dy</b> 164.93032 3 1.2,(4,6) 2973 1747 9.05 [Xe]6s <sup>2</sup> 4f <sup>9</sup>	
	83 164.93032	<b>Ho</b> 167.259 3 1.2,(4,6) 3141 1802 9.33 [Xe]6s <sup>2</sup> 4f <sup>10</sup>	
	84 167.259	<b>Er</b> 168.93421 2 1.2,(4,6) 2223 1818 9.98 [Xe]6s <sup>2</sup> 4f <sup>11</sup>	
	85 168.93421	<b>Tm</b> 173.054 2 1.2,(4,6) 3149 1092 Yb [Xe]6s <sup>2</sup> 4f <sup>12</sup>	
	86 173.054		
	87 223.0197	<b>Fr</b> 226.0254 1 1.2,(4,6) 950 300 — [Rn]7s <sup>1</sup>	
	88 226.0254	<b>Ra</b> 227.02781 2 1.2,(4,6) 1413 973 6.0 [Rn]7s <sup>1</sup>	
	89 227.02781	<b>Ac</b> 232.0381 3 1.2,(4,6) 4273 2023 11.70 [Rn]7s <sup>2</sup> 6d <sup>2</sup>	
	90 232.0381	<b>Th</b> 238.02891 4 1.2,(4,6) 4091 1405 15.40 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>1</sup>	
	91 238.02891	<b>Pa</b> 237.0462 5 1.2,(4,6) 4175 913 18.90 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>2</sup>	
	92 237.0462	<b>U</b> 244.0642 5 1.2,(4,6) 3505 914 20.45 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>3</sup>	
	93 244.0642	<b>Np</b> 246.0642 5 1.2,(4,6) 2880 1267 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>4</sup>	
	94 246.0642	<b>Pu</b> 249.0642 5 1.2,(4,6) 1613 — 13.5 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>5</sup>	
	95 249.0642	<b>Am</b> 257.0951 5 1.2,(4,6) 292 K 221 K 179 K 85 K 20 K 55 K [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>6</sup>	
	96 257.0951	<b>Cm</b> 257.0951 5 1.2,(4,6) 293 K 221 K 179 K 85 K 20 K 55 K [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>7</sup>	
	97 257.0951	<b>Bk</b> 257.0951 5 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>8</sup>	
	98 257.0951	<b>Cf</b> 257.0951 5 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>9</sup>	
	99 257.0951	<b>Es</b> 258.0984 3 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>10</sup>	
	100 258.0984	<b>Fm</b> 259.1010 3 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>11</sup>	
	101 259.1010	<b>Md</b> 259.1010 3 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>12</sup>	
	102 259.1010	<b>No</b> 259.1010 3 1.2,(4,6) 251.0796 251.0796 13.6 [Rn]7s <sup>2</sup> 5f <sup>6</sup> d <sup>13</sup>	

# Resumen

---

- Modelo de campo medio en un antiferromagneto
- Anisotropía magnetocristalina y formación de dominios
- Ciclo de histéresis de un ferromagneto
- Magnetismo itinerante
- Hamiltoniano de Hubbard y Criterio de Stoner

