



Campo y potencial electrostático de un distribución de cargas (discreta y continua)

$$\boxed{\vec{E}(\vec{r}) = -\vec{\nabla} V(\vec{r})}$$

$$\vec{E}(\vec{r}) = \sum_{n=1}^N \frac{1}{4\pi\epsilon_0} q_n \frac{\vec{r} - \vec{r}_n}{|\vec{r} - \vec{r}_n|^3}$$

$$V(\vec{r}) = \sum_{n=1}^N \frac{1}{4\pi\epsilon_0} q_n \frac{1}{|\vec{r} - \vec{r}_n|}$$



$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} dV'$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

Conectando distribuciones continuas y discretas



Delta de Dirac (1D)

$$\delta(x) = \begin{cases} +\infty, & x = 0 \\ 0, & x \neq 0 \end{cases}$$

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

Delta de Dirac (3D)

$$\delta(\vec{r}) = \delta(x) \delta(y) \delta(z)$$

$$\iiint_{-\infty}^{\infty} \delta(\vec{r}) dV = 1$$

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$$\delta(\vec{r}) = \delta(x) \delta(y) \delta(z)$$

$$\iiint_{-\infty}^{\infty} \delta(\vec{r}) dV = 1$$

Una carga en el origen

$$\rho(\vec{r}') = q \delta(\vec{r}')$$

$$\begin{aligned} V(\vec{r}) &= \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV' \\ &= \frac{1}{4\pi\epsilon_0} \iiint_V q \delta(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV' \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{|\vec{r}|} \end{aligned}$$

Conectando distribuciones continuas y discretas



Delta de Dirac (1D)

$$\delta(x) = \begin{cases} +\infty, & x = 0 \\ 0, & x \neq 0 \end{cases}$$

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

Delta de Dirac (3D)

$$\delta(\vec{r}) = \delta(x) \delta(y) \delta(z)$$

$$\iiint_{-\infty}^{\infty} \delta(\vec{r}) dV = 1$$

Una carga en \vec{r}_0

$$\rho(\vec{r}') = q \delta(\vec{r}' - \vec{r}_0)$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

$$= \frac{1}{4\pi\epsilon_0} \iiint_V q \delta(\vec{r}' - \vec{r}_0) \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{|\vec{r} - \vec{r}_0|}$$

Conectando distribuciones continuas y discretas



Delta de Dirac (1D)

$$\delta(x) = \begin{cases} +\infty, & x = 0 \\ 0, & x \neq 0 \end{cases}$$

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

Delta de Dirac (3D)

$$\delta(\vec{r}) = \delta(x) \delta(y) \delta(z)$$

$$\iiint_{-\infty}^{\infty} \delta(\vec{r}) dV = 1$$

Multiples carga q_n en \vec{r}_n

$$\rho(\vec{r}') = \sum_{n=1}^N q_n \delta(\vec{r}' - \vec{r}_n)$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

$$= \frac{1}{4\pi\epsilon_0} \iiint_V \sum_{n=1}^N q_n \delta(\vec{r}' - \vec{r}_n) \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

$$= \frac{1}{4\pi\epsilon_0} \sum_{n=1}^N q_n \frac{1}{|\vec{r} - \vec{r}_n|}$$

Conectando distribuciones en

volumétricas, **superficiales** y **lineales**

$$\rho(x', y', z') = \sigma(x', y') \delta(z') = \lambda(x') \delta(y') \delta(z')$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dx' dy' dz'$$

Expansión multipolar

campo y potencial
lejos de la distribución de cargas

El potencial lejos de la distribución de cargas

$$|\vec{r}| \gg |\vec{r}'|$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \iiint_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV'$$

$$\hat{r} = \frac{\vec{r}}{r}$$

$$|\vec{r} - \vec{r}'| = \sqrt{|r^2 - 2\vec{r}\vec{r}' + r'^2|} = r \sqrt{\left|1 - 2\frac{\hat{r}\vec{r}'}{r} + \left(\frac{r'}{r}\right)^2\right|}$$

El potencial lejos de la distribución de cargas

$$|\vec{r}| \gg |\vec{r}'|$$

Teorema del Binomio
(generalizado)

$$(a+x)^s = \sum_{n=0}^{\infty} \frac{s!}{n!(s-n)!} x^n a^{s-n}$$

$$(1+x)^s = \sum_{n=0}^{\infty} \frac{s!}{n!(s-n)!} x^n$$

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{r} \frac{1}{\sqrt{\left|1 - 2\frac{\hat{r}\vec{r}'}{r} + \left(\frac{r'}{r}\right)^2\right|}}$$

$\underbrace{\hspace{10em}}_{\epsilon}$

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{r} \left(1 - \frac{1}{2}\epsilon + \frac{3}{8}\epsilon^2 - \frac{5}{16}\epsilon^3 + \dots\right)$$

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{r} \left[1 - \frac{1}{2} \left(-2\frac{\hat{r}\vec{r}'}{r} + \left(\frac{r'}{r}\right)^2\right) + \frac{3}{8} \left(-2\frac{\hat{r}\vec{r}'}{r} + \left(\frac{r'}{r}\right)^2\right)^2 + \dots\right]$$

El potencial lejos de la distribución de cargas

$$|\vec{r}| \gg |\vec{r}'|$$

$$\begin{aligned} \frac{1}{|\vec{r} - \vec{r}'|} &= \frac{1}{r} \left[1 - \frac{1}{2} \left(-2 \frac{\hat{r} \cdot \vec{r}'}{r} + \left(\frac{r'}{r} \right)^2 \right) + \frac{3}{8} \left(-2 \frac{\hat{r} \cdot \vec{r}'}{r} + \left(\frac{r'}{r} \right)^2 \right)^2 + \dots \right] \\ &= \frac{1}{r} \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} - \frac{1}{2} \left(\frac{r'}{r} \right)^2 + \frac{3}{2} \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \frac{3}{2} \frac{\hat{r} \cdot \vec{r}'}{r} \left(\frac{r'}{r} \right)^2 + \frac{3}{8} \left(\frac{r'}{r} \right)^4 + \dots \right] \\ &= \frac{1}{r} \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) + \dots \right] \end{aligned}$$

El potencial lejos de la distribución de cargas

$$|\vec{r}| \gg |\vec{r}'|$$

$$\begin{aligned} V(\vec{r}) &= \frac{1}{4\pi\epsilon_0} \int_V \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV' \\ &= \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\vec{r}') \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) + \dots \right] dV' \end{aligned}$$

El potencial lejos de la distribución de cargas

Término monopolar

$$|\vec{r}| \gg |\vec{r}'|$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\vec{r}') \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) + \dots \right] dV'$$

$$V_{\text{mon}}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\vec{r}') dV' = \frac{Q}{4\pi\epsilon_0 r}$$

Carga total (escalar)

$$Q = \int_V \rho(\vec{r}') dV'$$

$$V_{\text{mon}}(\vec{r}) = \frac{Q}{4\pi\epsilon_0 r}$$

El potencial lejos de la distribución de cargas

Término dipolar

$$|\vec{r}| \gg |\vec{r}'|$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\vec{r}') \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) + \dots \right] dV'$$

$$V_{\text{dip}}(\vec{r}) = \frac{1}{4\pi\epsilon_0 r} \int_V \rho(\vec{r}') \frac{(\hat{r} \cdot \vec{r}')}{r} dV'$$

Momento dipolar (vector)

$$\vec{p} = \int_V \rho(\vec{r}') \vec{r}' dV'$$

$$= \frac{\hat{r}}{4\pi\epsilon_0 r^2} \int_V \rho(\vec{r}') \vec{r}' dV'$$

$$V_{\text{dip}}(\vec{r}) = \frac{\hat{r}}{4\pi\epsilon_0 r^2} \vec{p}$$

El potencial lejos de la distribución de cargas

Término cuadrupolar

$$|\vec{r}| \gg |\vec{r}'|$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\vec{r}') \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) + \dots \right] dV'$$

$$V_{\text{quad}}(\vec{r}) = \frac{1}{4\pi\epsilon_0 r} \int_V \rho(\vec{r}') \frac{1}{2} \left(3 \left(\frac{\hat{r} \cdot \vec{r}'}{r} \right)^2 - \left(\frac{r'}{r} \right)^2 \right) dV'$$

Momento cuadrupolar (tensor)

$$Q_{i,j} = \int_V \rho(\vec{r}') \frac{1}{2} (3 r_i r_j - r \delta_{i,j}) dV'$$

$$V_{\text{quad}}(\vec{r}) = \frac{\hat{r}}{4\pi\epsilon_0 r^3} \sum_{i,j} \frac{1}{2} Q_{i,j} \hat{r}_i \hat{r}_j$$

$$V(\vec{r}) = V_{\text{mon}}(\vec{r}) + V_{\text{dip}}(\vec{r}) + V_{\text{quad}}(\vec{r}) + \dots$$



