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Física Teórica 2 - guía 6.

Ejercicio Integrador guía 6.

(a) verifico que $S(r)$ es unitario, es decir, que se cumple que $S(r) S^\dagger(r) = \mathbb{1}$.

$$S(r) S^\dagger(r) = e^{\frac{i}{2}(a^2 - a'^2)} e^{\frac{i}{2}(a'^2 - a^2)} = e^{\frac{i}{2}(a^2 - a^2)} = e^0 = 1$$

o llamo $A = \frac{i}{2}(a^2 - a'^2)$ y $B = 1$

$[A, B] = 0$
 y se cumple $e^A B e^{-A} = e^{B^0} = 1$

$$S^{-1}(r) = \left(e^{\frac{i}{2}(a^2 - a'^2)} \right)^{\dagger} = e^{\frac{i}{2}(a'^2 - a^2)} = e^{-\frac{i}{2}(a^2 - a'^2)} = S^\dagger(r)$$

como $S(r) S^\dagger(r) = 1$

$S^\dagger(r) = S^{-1}(r)$

(b) uso la propiedad $e^A B e^{-A} = e^C B$ donde $C[A, B] = CB$.
 y llamo $A = \frac{i}{2}(a^2 - a'^2)$ y $B = x$

$$[A, B] = \frac{i}{2} [a^2 - a'^2, x] = \frac{i}{2} \frac{m\omega}{2\hbar} \left[\left(x + \frac{ip}{m\omega} \right)^2 - \left(x - \frac{ip}{m\omega} \right)^2, x \right]$$

$$= \frac{m\omega r}{4\hbar} \left[x^2 + \frac{2ixp}{m\omega} + \frac{p^2}{m^2\omega^2} - \left(x^2 - \frac{2ixp}{m\omega} + \frac{p^2}{m^2\omega^2} \right), x \right]$$

$$= \frac{m\omega r}{4\hbar} \left[\frac{2i}{m\omega} (xp + px), x \right] = \frac{ri}{2\hbar} [xp + px, x]$$

$$= \frac{ri}{2\hbar} (-2i\hbar x) = -ri^2 x = -rx \Rightarrow C = -r$$

$$[xp, x] = x[p, x] + [x, x]p = -i\hbar x$$

$$[px, x] = p[x, x] + [p, x]x = -i\hbar x$$

$$\Rightarrow e^A B e^{-A} = S^\dagger(r) x S(r) = e^{-rx}$$

(c) $Q_r \equiv e^r x$ $P_r \equiv e^{-r} p$ $a_r \equiv \sqrt{\frac{m\omega}{2\hbar}} \left(x e^r + \frac{p i e^{-r}}{m\omega} \right)$

$$[Q_r, P_r] = [e^r x, e^{-r} p] = e^r e^{-r} [x, p] = i\hbar$$

$\Rightarrow$ Grupo de conmutación canónica.

$$[a_r, a_r^\dagger] = \frac{m\omega}{2\hbar} \left[\left(x e^r + \frac{p i e^{-r}}{m\omega} \right), \left(x e^r - \frac{p i e^{-r}}{m\omega} \right) \right]$$

uso $a_r^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(x e^r - \frac{p i e^{-r}}{m\omega} \right)$

calculo auxiliar:

$$\begin{aligned} \bullet a_r a_r^\dagger &= \frac{m\omega}{2\hbar} \left(x e^r + \frac{p i e^{-r}}{m\omega} \right) \left(x e^r - \frac{p i e^{-r}}{m\omega} \right) = \\ &= \frac{m\omega}{2\hbar} \left(x^2 e^{2r} - \frac{i}{m\omega} x p + \frac{i}{m\omega} p x + \frac{p^2 e^{-2r}}{m^2 \omega^2} \right) \\ &\quad - \frac{i}{m\omega} (x p - p x) = -\frac{i}{m\omega} [x, p] = -\frac{i \cdot i \hbar}{m\omega} = \frac{\hbar}{m\omega} \end{aligned}$$

$$= \frac{m\omega}{2\hbar} \left(x^2 e^{2r} + \frac{\hbar}{m\omega} + \frac{p^2 e^{-2r}}{m^2 \omega^2} \right)$$

$$\begin{aligned} \bullet a_r^\dagger a_r &= \frac{m\omega}{2\hbar} \left(x e^r - \frac{p i e^{-r}}{m\omega} \right) \left(x e^r + \frac{p i e^{-r}}{m\omega} \right) \\ &= \frac{m\omega}{2\hbar} \left(x^2 e^{2r} + \frac{i}{m\omega} (x p - p x) + \frac{p^2 e^{-2r}}{m^2 \omega^2} \right) \\ &\quad = \frac{m\omega}{2\hbar} \left(x^2 e^{2r} - \frac{\hbar}{m\omega} + \frac{p^2 e^{-2r}}{m^2 \omega^2} \right) \end{aligned}$$

$$[a_r, a_r^\dagger] = a_r a_r^\dagger - a_r^\dagger a_r = 2 \cdot \frac{m\omega}{2\hbar} \cdot \frac{\hbar}{m\omega} = 1$$

a_r y $a_r^\dagger$ cumplen con las condiciones de conmutación de los operadores de creación y destrucción y Q_r y P_r son los análogos a los operadores de creación y destrucción posición y momento.

(d)

 ~~$S(\alpha)$~~

$$a_r |\alpha, r\rangle = \underbrace{S(r)}_{A=S(r)S^\dagger(r)} a_r S(r) |\alpha\rangle = S(r) S^\dagger(r) a_r S(r) |\alpha\rangle$$

$$= S(r) \cdot \sqrt{\frac{m\omega}{2\hbar}} S^\dagger(r) \left(x e^r + \frac{ip}{m\omega} e^{-r} \right) S(r) |\alpha\rangle$$

$$= S(r) \cdot \sqrt{\frac{m\omega}{2\hbar}} \left[\underbrace{e^r S^\dagger(r) x S(r)}_{e^{-r} x} + \frac{ip}{m\omega} \underbrace{e^{-r} S^\dagger(r) p S(r)}_{e^r p} \right] |\alpha\rangle$$

$$= \sqrt{\frac{m\omega}{2\hbar}} S(r) \left[e^r e^{-r} x + \frac{i}{m\omega} e^{-r} e^r p \right] |\alpha\rangle$$

$$= \sqrt{\frac{m\omega}{2\hbar}} S(r) \left[x + \frac{ip}{m\omega} \right] |\alpha\rangle$$

$$= S(r) \cdot a |\alpha\rangle \stackrel{A}{=} S(r) |\alpha\rangle$$

$$a |\alpha\rangle = \alpha |\alpha\rangle$$

$$= S(r) \alpha |\alpha\rangle = \alpha S(r) |\alpha\rangle = \alpha |\alpha, r\rangle$$

$\Rightarrow |\alpha, r\rangle$ es ~~autovalor~~ ^{autoestado} de a_r con autovalor α .

Esto implica que al hacer un cambio de variables tal que $x \rightarrow Q_r = e^r x$ y $p \rightarrow P_r = e^{-r} p$

los estados coherentes siguen la transformación

$|\alpha\rangle \rightarrow |\alpha, r\rangle = S(r) |\alpha\rangle$ y mantienen el

que sea ~~energía~~ ~~autovalor~~ ~~energía~~ ~~energía~~

del estado mismo autovalor.

(e)

$$\langle \alpha, r | x | \alpha, r \rangle = \langle \alpha | S^\dagger(r) x S(r) | \alpha \rangle = \langle \alpha | e^{-r} x | \alpha \rangle = e^{-r} \langle \alpha | x | \alpha \rangle$$

$$\langle \alpha, r | = \langle \alpha | S^\dagger(r)$$

$$\langle x \rangle = e^{-r} \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^*)$$

$$\langle \alpha, r | p | \alpha, r \rangle = \langle \alpha | S^\dagger(r) p S(r) | \alpha \rangle = e^r \langle \alpha | p | \alpha \rangle = e^r \langle \alpha | \frac{\hbar}{i} \frac{d}{dx} | \alpha \rangle$$

$$= \frac{e^r}{i} \sqrt{\frac{\hbar m \omega}{2}} (\alpha - \alpha^*)$$

$$\langle \alpha, r | x^2 | \alpha, r \rangle = \langle \alpha | S^\dagger(r) x^2 S(r) | \alpha \rangle = e^{-2r} \langle \alpha | x^2 | \alpha \rangle$$

$$= e^{-2r} \frac{\hbar}{2m\omega} (\alpha^2 + \alpha^{*2} + 2|\alpha|^2 + 1)$$

$$S^\dagger(r) x^2 S(r) = e^{-2r} x^2$$

$$\left[\frac{r}{2} (\alpha^{+2} - \alpha^{*2}) \right] = x \left[\frac{r}{2} (\alpha^{+2} - \alpha^{*2}) \right] + \left[\frac{r}{2} \alpha^{+2} - \alpha^{*2} \right] x = -2rx^2$$

$$\langle \alpha, r | p^2 | \alpha, r \rangle = e^{2r} \langle \alpha | p^2 | \alpha \rangle = \frac{-e^{2r}}{2} \frac{\hbar^2 m \omega}{2} (\alpha^2 + \alpha^{*2} - 2|\alpha|^2 - 1)$$

$$\Rightarrow \text{Var}(x) = \frac{\hbar}{2m\omega} e^{-2r} (\alpha^2 + \alpha^{*2} + 2|\alpha|^2 + 1) - \frac{\hbar}{2m\omega} e^{-2r} (\alpha^2 + \alpha^{*2} + 2|\alpha|^2)$$

$$= \frac{\hbar}{2m\omega} e^{-2r}$$

$$\text{Var}(p) = -e^{2r} \frac{\hbar^2 m \omega}{2} (\alpha^2 + \alpha^{*2} - 2|\alpha|^2 - 1)$$

$$- \left(-e^{2r} \frac{\hbar^2 m \omega}{2} \right) (\alpha^2 + \alpha^{*2} - 2|\alpha|^2)$$

$$= e^{2r} \frac{\hbar^2 m \omega}{2}$$

$$\Rightarrow \text{Var}(x) \text{Var}(p) = \frac{\hbar}{2m\omega} e^{-2r} e^{2r} \frac{\hbar^2 m \omega}{2} = \frac{\hbar^2}{4}$$

$\Rightarrow$ son los estados de mínima incertidumbre
lo cual implica que su función de
onda es de la forma de un paquete
gaussiano. los estados $|\alpha, r\rangle$ son una
traslación en el espacio de fases del estado
fundamental

OK

Donc se obtiennent les siguientes resultados:

$$\begin{aligned}\langle \alpha | x | \alpha \rangle &= \langle \alpha | \sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger) | \alpha \rangle = \sqrt{\frac{\hbar}{2m\omega}} \langle \alpha | a + a^\dagger | \alpha \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega}} (\langle \alpha | a | \alpha \rangle + \langle \alpha | a^\dagger | \alpha \rangle) \\ &= \sqrt{\frac{\hbar}{2m\omega}} (\alpha^* + \alpha)\end{aligned}$$

Calculo auxiliares:

- $\langle \alpha | a^\dagger | \alpha \rangle = \alpha^* \langle \alpha | \alpha \rangle = \alpha^*$
- $\langle \alpha | a^2 | \alpha \rangle = \alpha \langle \alpha | a | \alpha \rangle = \alpha^2$
- $\langle \alpha | a^{\dagger 2} | \alpha \rangle = \alpha^{*2}$
- $\langle \alpha | a^\dagger a | \alpha \rangle = \langle \alpha | a a^\dagger | \alpha \rangle = 1$

$a | \alpha \rangle = \alpha | \alpha \rangle$
 $\langle \alpha | a^\dagger = \langle \alpha | a^*$

- $\langle \alpha | a | \alpha \rangle = \alpha \langle \alpha | \alpha \rangle = \alpha$

$$\langle \alpha | p | \alpha \rangle = \frac{1}{i} \sqrt{\frac{\hbar m \omega}{2}} \langle \alpha | a - a^\dagger | \alpha \rangle = \sqrt{\frac{\hbar m \omega}{2}} \frac{\alpha - \alpha^*}{i}$$

$$\begin{aligned}\langle \alpha | x^2 | \alpha \rangle &= \frac{\hbar}{2m\omega} \langle \alpha | (a + a^\dagger)(a + a^\dagger) | \alpha \rangle = \frac{\hbar}{2m\omega} \langle \alpha | a^2 + a a^\dagger + a^\dagger a + a^{\dagger 2} | \alpha \rangle \\ &= \frac{\hbar}{2m\omega} \langle \alpha | a^2 + \underbrace{a a^\dagger + a^\dagger a}_{1} + a^{\dagger 2} | \alpha \rangle = \frac{\hbar}{2m\omega} \langle \alpha | a^2 + 2a^\dagger a + 1 + a^{\dagger 2} | \alpha \rangle \\ 1 &= [a^\dagger, a] = a^\dagger a - a a^\dagger \\ \Rightarrow a a^\dagger &= 1 + a^\dagger a\end{aligned}$$

$$= \frac{\hbar}{2m\omega} \left(\alpha^2 + 2|\alpha|^2 + 1 + \alpha^{*2} \right)$$

$$\begin{aligned}\langle \alpha | p^2 | \alpha \rangle &= \frac{\hbar m \omega}{2} \langle \alpha | (a - a^\dagger)(a - a^\dagger) | \alpha \rangle = \frac{\hbar m \omega}{2} \langle \alpha | a^2 - a a^\dagger - a^\dagger a + a^{\dagger 2} | \alpha \rangle \\ &= \frac{\hbar m \omega}{2} \langle \alpha | a^2 - 2a^\dagger a - 1 + a^{\dagger 2} | \alpha \rangle = \frac{\hbar m \omega}{2} (\alpha^2 - 2|\alpha|^2 - 1 + \alpha^{*2})\end{aligned}$$

(f)

$$\langle \alpha, r | H | \alpha, r \rangle = \langle \alpha | S^\dagger(r) H S(r) | \alpha \rangle = \langle \alpha | S^\dagger(r) \left(\frac{p^2}{2m} + \frac{m\omega^2}{2} x^2 \right) S(r) | \alpha \rangle$$

como Hamiltoniano de oscilador

$$\text{armónico } H = \frac{p^2}{2m} + \frac{m\omega^2}{2} x^2$$

$$= \frac{1}{2m} \langle \alpha | S^\dagger(r) p^2 S(r) | \alpha \rangle + \frac{m\omega^2}{2} \langle \alpha | S^\dagger(r) x^2 S(r) | \alpha \rangle$$

$$= \frac{1}{2m} \langle \alpha | e^{2r} p^2 | \alpha \rangle + \frac{m\omega^2}{2} \langle \alpha | e^{-2r} x^2 | \alpha \rangle$$

$$= \frac{e^{2r}}{2m} \langle \alpha | p^2 | \alpha \rangle + \frac{m\omega^2}{2} e^{-2r} \langle \alpha | x^2 | \alpha \rangle$$

$$= \frac{e^{2r}}{2m} \cdot \frac{1}{2m\omega} (d^2 + d^{*2} - 2|\alpha|^2 - 1) + e^{-2r} \frac{m\omega}{2} (d^2 + d^{*2} + 2|\alpha|^2 + 1)$$

$$= (d^2 + d^{*2}) \left(\frac{\hbar\omega}{4} e^{2r} + \frac{\hbar\omega}{4} e^{-2r} \right) +$$

$$+ \frac{\hbar m\omega}{4} (d^2 + d^{*2}) (-e^{2r} + e^{-2r}) + \frac{\hbar\omega}{4} (2|\alpha|^2 + 1) (e^{2r} + e^{-2r})$$

$$= \frac{i\hbar\omega}{2} (d^2 + d^{*2}) \sin(2r) + \frac{\hbar\omega}{2} (2|\alpha|^2 + 1) \cos(2r)$$

(g) Para el caso $|0, r\rangle$ tengo:

$$\langle x(t) \rangle = \langle 0, r | x \cos(\omega t) + \frac{p}{m\omega} \sin(\omega t) | 0, r \rangle$$

$$= \cos(\omega t) \langle 0, r | x | 0, r \rangle + \frac{\sin(\omega t)}{m\omega} \langle 0, r | p | 0, r \rangle$$

$$\stackrel{=0}{\uparrow} \text{ del inciso (e) } \langle x \rangle \Big|_{\alpha=0} = \langle p \rangle \Big|_{\alpha=0} = 0$$

$$\langle p(t) \rangle = 0 \text{ análogamente.}$$

$$\begin{aligned}
 \langle x(t)^2 \rangle &= \langle 0, r | (x \cos \omega t + \frac{p}{m\omega} \sin \omega t) (x \cos \omega t + \frac{p}{m\omega} \sin \omega t) | 0, r \rangle \\
 &= \langle 0, r | x^2 \cos^2 \omega t + \frac{p^2}{m^2 \omega^2} \sin^2 \omega t + (xp + px) \frac{\cos \omega t \sin \omega t}{m\omega} | 0, r \rangle \\
 &= \langle 0, r | x^2 | 0, r \rangle \cos^2 \omega t + \langle 0, r | p^2 | 0, r \rangle \frac{\sin^2 \omega t}{m^2 \omega^2} + \\
 &\quad + \langle 0, r | xp + px | 0, r \rangle \frac{\cos \omega t \sin \omega t}{m\omega}
 \end{aligned}$$

Calculo auxiliar:

$$\bullet \langle 0, r | x^2 | 0, r \rangle = \frac{\hbar}{2m\omega} e^{-2r}$$

↑
maiso(e) con $\alpha=0$

$$\bullet \langle 0, r | p^2 | 0, r \rangle = \frac{\hbar m \omega}{2} e^{2r}$$

$$\bullet \langle 0, r | xp + px | 0, r \rangle = \langle 0, r | xp + i\hbar | 0, r \rangle$$

$$i\hbar = [x, p] = xp - px \Rightarrow px = xp + i\hbar$$

$$= 2\langle 0, r | xp | 0, r \rangle + i\hbar = \frac{\hbar}{2} (1 + i)$$

$$\bullet \langle \alpha, r | xp | \alpha, r \rangle = \langle \alpha | s^\dagger(r) \frac{\hbar}{2} (a^2 - a^{\dagger 2} + a a^\dagger + a^\dagger a) s(r) | \alpha \rangle$$

$[a^\dagger, a] = 1$

$$= \frac{\hbar}{2} \langle \alpha | s^\dagger(r) (a^2 - a^{\dagger 2} + 1) s(r) | \alpha \rangle$$

$$= \langle \alpha | s^\dagger(r) x p s(r) | \alpha \rangle$$

$$= \frac{\hbar}{2} \langle \alpha | \underbrace{s^\dagger(r) x s(r)}_{e^{-r} x} \underbrace{s^\dagger(r) p s(r)}_{e^r p} | \alpha \rangle = \langle \alpha | xp | \alpha \rangle$$

$$= \frac{\hbar}{2} \langle \alpha | a^2 - a^{\dagger 2} + \underbrace{a^\dagger a - a a^\dagger}_{=[a^\dagger, a]=1} | \alpha \rangle$$

$$= \frac{\hbar}{2} (\alpha^2 - \alpha^{*2} + 1)$$

$$\langle x(t)^2 \rangle = \frac{\hbar}{2m\omega} \cos^2 \omega t + e^{2r} \frac{\hbar m \omega}{2} \frac{\sin^2 \omega t}{m^2 \omega^2} + \frac{\hbar}{m\omega} \cos \omega t \sin \omega t (\alpha^2 - \alpha^{*2} + 1)$$

$$\langle x(t) \rangle = \frac{\hbar}{2m\omega} \left(e^{-2r} \cos^2 \omega t + e^{2r} \sin^2 \omega t \right) + \frac{\hbar}{2m\omega} \sin(2\omega t) (1+i) \times \text{[scribble]}$$

$$\langle x^2(t) \rangle = \frac{\hbar}{2m\omega} [e$$

$$\langle p^2(t) \rangle = \langle 0, r | p^2 \cos^2 \omega t + m^2 \omega^2 x^2 \sin^2 \omega t + (xp + px) \frac{m\omega}{2} \sin(2\omega t) | 0, r \rangle$$

$$= \langle 0, r | p^2 | 0, r \rangle \cos^2 \omega t + m^2 \omega^2 \langle 0, r | x^2 | 0, r \rangle \sin^2 \omega t - \frac{m\omega}{2} \sin(2\omega t) \langle 0, r | xp + px | 0, r \rangle$$

$$= \frac{\hbar^2}{2} \cos^2 \omega t + m^2 \omega^2 \frac{\hbar^2}{2} \sin^2 \omega t - \frac{m\omega}{2} \sin(2\omega t) \hbar(1+i)$$

$$= \frac{\hbar^2 m \omega}{2} \left[e^{2r} \cos^2 \omega t + e^{-2r} \sin^2 \omega t - \sin(2\omega t)(1+i) \right]$$

$$\text{Var}[x(t)] = \frac{\hbar}{2m\omega} \left[e^{-2r} \cos^2 \omega t + e^{2r} \sin^2 \omega t + \sin(2\omega t)(1+i) \right]$$

$$\text{Var}[p(t)] = \frac{\hbar^2 m \omega}{2} \left[e^{2r} \cos^2 \omega t + e^{-2r} \sin^2 \omega t - \sin(2\omega t)(1+i) \right]$$

$$\text{Var}(x(t)) \text{Var}(p(t)) = \frac{\hbar^2}{4} \left[\cos^4 \omega t + \sin^4 \omega t - \sin^2(2\omega t)(2i^0) + \right.$$

$$\left. + e^{4r} \cos^2 \omega t \sin^2 \omega t (e^{-4r} + e^{4r}) + \sin(2\omega t)(1+i) \cos^2 \omega t (e^{2r} - e^{-2r}) - \sin(2\omega t)(1+i) \sin^2 \omega t (e^{2r} - e^{-2r}) \right]$$

$$\text{Var}[x(t)] = \frac{\hbar}{2m\omega} \left[e^{2r} \cos^2 \omega t + e^{-2r} \sin^2 \omega t - \sin(2\omega t)(1+i) \right]$$

$$\text{Var}(x(t)) \text{Var}(p(t)) = \frac{\hbar^2}{4} \left[\cos^4 \omega t + \sin^4 \omega t - \sin^2(2\omega t)(2i^0) + e^{4r} \cos^2 \omega t \sin^2 \omega t (e^{-4r} + e^{4r}) + \sin(2\omega t)(1+i) \cos^2 \omega t (e^{2r} - e^{-2r}) - \sin(2\omega t)(1+i) \sin^2 \omega t (e^{2r} - e^{-2r}) \right]$$

$$\text{Var}(x(t)) \text{Var}(p(t)) = \frac{\hbar^2}{4} \left[\cos^4 \omega t + \sin^4 \omega t - i \sin^2(2\omega t) (1 + i \cos(4r)) \right. \\ \left. + i \sin(2\omega t) \sin(2r) (1 + i) (\cos^2 \omega t - \sin^2 \omega t) \right]$$

$$= \frac{\hbar^2}{4} \left[\cos^4(\omega t) + \sin^4(\omega t) + \sin^2(2\omega t) (\cos(4r) - i) \right. \\ \left. + \sin(2\omega t) \sin(2r) \left(\frac{i-1}{2} \right) \cos(2\omega t) \right]$$

Si $t=0$

$$\text{Var}(x) \text{Var}(p) = \frac{\hbar^2}{4}$$

Si $t = \frac{\pi}{2\omega}$

$$\text{Var}(x) \text{Var}(p) = \frac{\hbar^2}{4} \left[\frac{1}{4} + \frac{1}{4} \right] = \frac{\hbar^2}{4}$$

Si $t = \frac{\pi}{\omega}$

$$\text{Var}(x) \text{Var}(p) = 1 \cdot \frac{\hbar^2}{4}$$

$\Rightarrow$ Para $t=0, \frac{\pi}{2\omega}$ y $t=\frac{\pi}{\omega}$ el producto de

las varianzas es mínima y esperamos que al tener incerteza mínima tenga la forma de una gaussiana.

Dado que el estado squeezed general $|a, r\rangle$ es una traslación en el espacio de fases de $|0, r\rangle$ espero que evolucione temporalmente de la misma forma con una diferente frecuencia.