

Momento angular

$$\left. \begin{array}{l}
 \text{momento angular orbital (OAM): } \vec{L} \\
 \text{" " de spin (SAM): } \vec{S}
 \end{array} \right\} \begin{array}{l}
 \vec{J} = \vec{L} + \vec{S} \\
 \text{Total angular} \\
 \text{momentum (TAM)} \\
 \text{MAT}
 \end{array}$$

Notación: mom. ang. arbitrario: \vec{J}

OAM

Usando las reglas de cuantización:

$$\left. \begin{array}{l}
 L_x = y p_z - z p_y \\
 L_y = z p_x - x p_z \\
 L_z = x p_y - y p_x
 \end{array} \right\} \vec{L} = \vec{R} \times \vec{P}$$

Evaluemos:

$$\begin{aligned}
 [L_x, L_y] &= [y p_z - z p_y, z p_x - x p_z] \\
 &= [y p_z, z p_x] - [y p_z, x p_z] - [z p_y, z p_x] + [z p_y, x p_z] \\
 &= [y p_z, z p_x] + [z p_y, x p_z] \\
 &= y [p_z, z] p_x + x [z, p_z] p_y \\
 &= -i\hbar y p_x + i\hbar x p_y = i\hbar (x p_y - y p_x) \\
 &= i\hbar L_z
 \end{aligned}$$

y análogamente

$$\begin{aligned}
 [L_y, L_z] &= i\hbar L_x & [L_i, L_j] &= i\hbar \epsilon_{ijk} L_k \\
 [L_z, L_x] &= i\hbar L_y
 \end{aligned}$$

H₀ [N³]

Nombre:

húsaes

Def. un mom. ang. general \vec{J} por las reglas de conmutación:

$$[J_i, J_j] = i\hbar \epsilon_{ijk} J_k$$

Introducimos el operador hermítico:

$$J^2 = J_x^2 + J_y^2 + J_z^2$$

Se muestra fácilmente que:

$$[J^2, J_i] = 0 \quad (*)$$

Comutan y \therefore son operadores compatibles.

En cambio, J_x, J_y, J_z no.

En problemas con potencial central tomaremos

$$\{H, L^2, L_z\} \text{ como CCOC.}$$

(*) Prueba:

$$\begin{aligned} [J^2, J_x] &= [J_x^2 + J_y^2 + J_z^2, J_x] = [J_y^2, J_x] + [J_z^2, J_x] \\ &= J_y [J_y, J_x] + [J_y, J_x] J_y + J_z [J_z, J_x] + [J_z, J_x] J_z \\ &= -i\hbar J_y J_z - i\hbar J_z J_y + i\hbar J_z J_y + i\hbar J_y J_z = 0 \end{aligned}$$

Ladder operators

$$J_{\pm} = J_x \pm i J_y$$

$$\begin{aligned} J_+ J_- &= (J_x + i J_y)(J_x - i J_y) \\ &= J_x^2 + J_y^2 + i J_y J_x - i J_x J_y \\ &= J_x^2 + J_y^2 - i [J_x, J_y] \end{aligned}$$

$$\blacksquare = J_x^2 + J_y^2 + \hbar J_z \quad (*)$$

$$J_- J_+ = J_x^2 + J_y^2 - \hbar J_z \quad (*)$$

$$\Rightarrow [J_+, J_-] = 2\hbar J_z$$

$$[J^2, J_{\pm}] = [J^2, J_x \pm i J_y] = [J^2, J_x] \pm i [J^2, J_y] = 0$$

$$\begin{aligned} [J_z, J_{\pm}] &= [J_z, J_x \pm i J_y] \\ &= [J_z, J_x] \pm i [J_z, J_y] \\ &= (i\hbar J_y \pm i(-i\hbar)J_x) \\ &= \blacksquare i\hbar J_y \blacksquare \pm \hbar J_x \end{aligned}$$

$$= \pm \hbar (J_x \pm i \blacksquare J_y) = \pm \hbar J_{\pm}$$

(*) usando $J^2 = J_x^2 + J_y^2 + J_z^2$

$$\begin{cases} J_+ J_- = J^2 - J_z^2 + \hbar J_z & \Rightarrow J^2 = J_+ J_- + J_z^2 - \hbar J_z \\ J_- J_+ = J^2 - J_z^2 - \hbar J_z & \Rightarrow J^2 = J_- J_+ + J_z^2 + \hbar J_z \end{cases}$$

$$\Rightarrow J^2 = \frac{1}{2} (J_+ J_- + J_- J_+) + J_z^2$$

Teo. Los autovalores de J^2 son positivos ($J^2|\lambda\rangle = \lambda|\lambda\rangle$)

$$\forall |\psi\rangle : \langle \psi | J^2 | \psi \rangle = \langle \psi | J_x^2 | \psi \rangle + \langle \psi | J_y^2 | \psi \rangle + \langle \psi | J_z^2 | \psi \rangle$$

$$= \| J_x | \psi \rangle \|^2 + \| J_y | \psi \rangle \|^2 + \| J_z | \psi \rangle \|^2 \geq 0$$

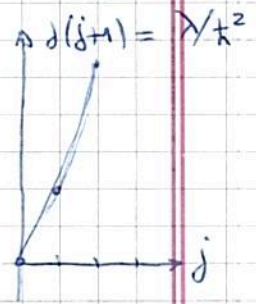
En particular si $|\psi\rangle = |\lambda\rangle$ es autoestado de $J^2|\lambda\rangle = \lambda|\lambda\rangle$

~~$$J^2|\lambda\rangle = \lambda|\lambda\rangle$$~~

$$\langle \lambda | J^2 | \lambda \rangle = \lambda \langle \lambda | \lambda \rangle = \lambda \geq 0$$

Podemos expresarlo : $\lambda = j(j+1)\hbar^2$ (con $j \geq 0$)

~~Autovalores de J_z : $m\hbar$~~



Problema : resolver simultáneamente:

$$\begin{cases} J^2 |j, m\rangle = j(j+1)\hbar^2 |j, m\rangle \\ J_z |j, m\rangle = m\hbar |j, m\rangle \end{cases}$$

¿Cuáles son los posibles valores de j, m ?

~~$$-j \leq m \leq j \quad (|m| \leq j)$$~~

Lemma 1

~~$$(i) \text{ si } m = j \Rightarrow J_- |j, j\rangle = 0$$~~

~~$$(ii) \text{ si } m > -j \Rightarrow J^2 J_- |j, m\rangle = j(j+1)\hbar^2 J_- |j, m\rangle$$~~

~~$$J_z J_- |j, m\rangle = (m-1)\hbar J_- |j, m\rangle$$~~

Autovalores de J^2 y J_z

- Si $|jm\rangle$ es autoestado común de J^2 y J_z :

$$\begin{cases} J^2 |jm\rangle = j(j+1)\hbar^2 |jm\rangle \\ J_z |jm\rangle = m\hbar |jm\rangle \end{cases}$$

Entonces $J_{\pm} |jm\rangle$ también lo es.

Dem.

$$J^2 J_{\pm} |jm\rangle = J_{\pm} J^2 |jm\rangle = j(j+1)\hbar^2 J_{\pm} |jm\rangle$$

\uparrow $[J^2, J_{\pm}] = 0$ \downarrow mismo autovalor

$$\begin{aligned} J_z J_{\pm} |jm\rangle &= (J_{\pm} J_z \pm \hbar J_{\pm}) |jm\rangle = \\ &\uparrow \\ &[J_z, J_{\pm}] = \pm \hbar J_{\pm} \\ &= J_{\pm} (J_z \pm \hbar \mathbb{1}) |jm\rangle \\ &= J_{\pm} (m\hbar \pm \hbar) |jm\rangle \\ &= \underbrace{(m \pm 1)\hbar}_{\text{cambia autovalor } \pm \hbar} J_{\pm} |jm\rangle \end{aligned}$$

autovalores
de J_z

$(m+2)\hbar$	$J_+^2 jm\rangle$
$(m+1)\hbar$	$J_+ jm\rangle$
$m\hbar$	$ jm\rangle$
$(m-1)\hbar$	$J_- jm\rangle$
$(m-2)\hbar$	$J_-^2 jm\rangle$

j fijo

• Existe un m_{\max} / $J_+ |j m_{\max}\rangle = 0$

(La componente J_z no puede exceder el J total)

~~Desarrollando~~ \Rightarrow

$$\begin{aligned} J^2 |j m_{\max}\rangle &= (J_- J_+ + J_z^2 + \hbar J_z) |j m_{\max}\rangle \\ &= (0 + m_{\max}^2 \hbar^2 + m_{\max} \hbar^2) |j m_{\max}\rangle \\ &= m_{\max} (m_{\max} + 1) \hbar^2 |j m_{\max}\rangle \\ &= j(j+1) \hbar^2 |j m_{\max}\rangle \end{aligned}$$

• También existe m_{\min} / $J_- |j m_{\min}\rangle = 0 \Rightarrow$

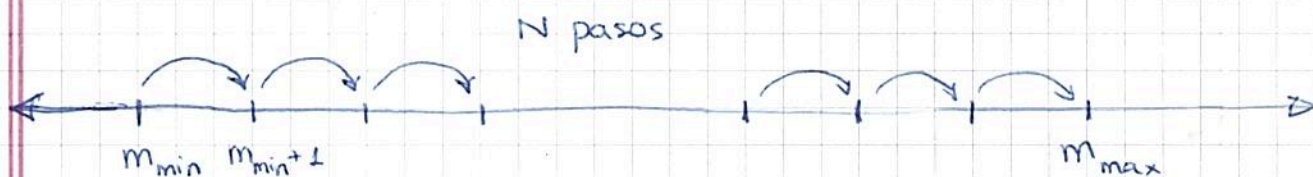
$$\begin{aligned} J^2 |j m_{\min}\rangle &= (J_+ J_- + J_z^2 - \hbar J_z) |j m_{\min}\rangle \\ &= m_{\min} (m_{\min} - 1) \hbar^2 |j m_{\min}\rangle \\ &= j(j+1) \hbar^2 |j m_{\min}\rangle \end{aligned}$$

Vemos que:

$$m_{\max} (m_{\max} + 1) = m_{\min} (m_{\min} - 1)$$

Si $m_{\max} = m_{\min} - 1$, se satisface, pero $m_{\max} < m_{\min}$ no sirve

Si $m_{\max} = -m_{\min}$, se satisface \checkmark



Tenemos $m_{\max} = m_{\min} + N = -m_{\max} + N$

$$\Rightarrow N = 2m_{\max} \Rightarrow m_{\max} = \frac{N}{2}$$

~~XXXXXXXXXX~~~~XX~~~~XX~~

- Además: $j = m_{\max}$

Conclusión sobre autovalores:

$$\begin{cases} j = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots \\ m = -j, -j+1, \dots, j-1, j \end{cases}$$

~~Autovectores comunes de J^2 y J_z~~

~~Autovectores de J^2~~ Normalization

$$\text{Si } J_+ |j, m\rangle = A_{jm} |j, m+1\rangle$$

Hallar A_{jm} que normalice $J_+ |j, m\rangle$

$$J_+ |j, m\rangle \propto |j, m+1\rangle$$

↳ está normalizado?

$$\begin{aligned} \|J_+ |j, m\rangle\|^2 &= \langle j, m | J_- J_+ |j, m\rangle \\ &= \langle j, m | J^2 - J_z^2 - \hbar J_z |j, m\rangle \\ &= j(j+1)\hbar^2 - m(m+1)\hbar^2 \end{aligned}$$

$$\Rightarrow |j, m+1\rangle = \frac{1}{\sqrt{j(j+1)\hbar^2 - m(m+1)\hbar^2}} J_+ |j, m\rangle$$

En síntesis:

$$J_z |j m\rangle = m \hbar |j m\rangle$$

$$J_{\pm} |j m\rangle = \hbar \sqrt{j(j+1) - m(m \pm 1)} |j, m \pm 1\rangle$$

$$J^2 |j m\rangle = j(j+1) \hbar^2 |j m\rangle$$

Matriz de J^2

Elementos de matriz: $\langle j' m' | J^2 | j m \rangle$

$$\langle j' m' | J^2 | j m \rangle = j(j+1) \hbar^2 \delta_{j'j} \delta_{m'm}$$

j'	m'	j	m	$= j(j+1) \hbar^2 \delta_{j'j} \delta_{m'm}$			
0	0	0	0	0	0	0	0
	0	0	0	0	0	0	0
1/2	1/2	1/2	1/2	3/4 \hbar^2			
	-1/2	1/2	-1/2	3/4 \hbar^2			
1	1	1	1		2 \hbar^2		
	0	1	0		2 \hbar^2		
	-1	1	-1			2 \hbar^2	
3/2	3/2	3/2	3/2				15/4 \hbar^2
	1/2	3/2	1/2				15/4 \hbar^2
	-1/2	3/2	-1/2				15/4 \hbar^2
	-3/2	3/2	-3/2				15/4 \hbar^2

$$\frac{3 \cdot 5}{2 \cdot 2} = \frac{15}{4}$$

En el pizarrón, borrar los valores y escribir:

$$\langle j' m' | J_z | j m \rangle = m \hbar \delta_{j'j} \delta_{m'm}$$

Elementos de matriz de J_x y J_y

$$\begin{cases} J_+ = J_x + i J_y \\ J_- = J_x - i J_y \end{cases} \quad \begin{cases} J_x = \frac{1}{2} (J_+ + J_-) \\ J_y = \frac{1}{2i} (J_+ - J_-) = \frac{i}{2} (J_- - J_+) \end{cases}$$

$$\begin{aligned} \langle j m | J_{\pm} | j' m' \rangle &= \langle j m | \hbar \sqrt{j'(j'+1) - m'(m' \pm 1)} | j', m' \pm 1 \rangle \\ &= \hbar \sqrt{j(j+1) - m'(m' \pm 1)} \delta_{jj'} \delta_{m, m' \pm 1} \end{aligned}$$

Entonces:

$$\begin{aligned} \langle j m | J_x | j' m' \rangle &= \frac{1}{2} \langle j m | J_+ | j' m' \rangle + \frac{1}{2} \langle j m | J_- | j' m' \rangle \\ &= \frac{\hbar \delta_{jj'}}{2} \left(\sqrt{j(j+1) - m'(m'+1)} \delta_{m, m'+1} + \sqrt{j(j+1) - m'(m'-1)} \delta_{m, m'-1} \right) \end{aligned}$$

$$\langle j m | J_y | j' m' \rangle = \frac{\hbar}{2i} \left(\sqrt{\dots} \delta_{m, m'+1} - \sqrt{\dots} \delta_{m, m'-1} \right) \delta_{jj'}$$

$$j = 1/2$$

$$(J_x)^{1/2} = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$(J_y)^{1/2} = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$(J_+)^{1/2} = \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$(J_-)^{1/2} = \frac{\hbar}{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$j = 1 \quad (J_+)^1 = \hbar \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix}$$

$$(J_-)^1 = \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix} \hbar$$

$$(J_x)^1 = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$(J_y)^1 = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}$$