

$$1. \quad (a) \quad \left(\frac{\partial S}{\partial T} \right)_x = \frac{C_x}{T} \Rightarrow S(T, x) = C_x \ln T + g(x)$$

$$\left(\frac{\partial S}{\partial x} \right)_T = - \left(\frac{\partial f}{\partial T} \right)_x = -\alpha x \Rightarrow g'(x) = -\alpha x \Rightarrow g(x) = -\frac{1}{2} \alpha x^2 + c$$

Maxwell:

$$dF = -S dT + f dx$$

$$\Rightarrow - \left(\frac{\partial S}{\partial x} \right)_T = \left(\frac{\partial f}{\partial T} \right)_x$$

$$\Rightarrow \text{Escribiendo } c = -C_x \ln \theta,$$

$$S(T, x) = C_x \ln \frac{T}{\theta} - \frac{1}{2} \alpha x^2$$

$$(b) \quad \left(\frac{\partial F}{\partial T} \right)_x = -S = \frac{1}{2} \alpha x^2 - C_x \ln \frac{T}{\theta}$$

$$\Rightarrow F(T, x) = \frac{1}{2} \alpha T x^2 - C_x [T(\ln T - 1) - T \ln \theta] + g(x)$$

$$= \frac{1}{2} \alpha T x^2 - C_x T \left(\ln \frac{T}{\theta} - 1 \right) + g(x)$$

$$\left(\frac{\partial F}{\partial x} \right)_T = f = \alpha T x \Rightarrow \alpha T x + g'(x) = \alpha T x \Rightarrow g'(x) = 0$$

$$\Rightarrow g(x) = c$$

Eligiendo origen de energías,

$$F(T, x) = T \left[\frac{1}{2} \alpha x^2 - C_x \left(\ln \frac{T}{T_i} - 1 \right) \right]$$

$$\Rightarrow U(T, x) = F + TS = C_x T = U(T, x)$$

$$(c) \quad S(T, x) - S(T_i, 0) = C_x \ln \frac{T}{T_i} - \frac{1}{2} \alpha x^2$$

Adiabático y cuasiestático $\Rightarrow \Delta S = 0$

$$\Rightarrow C_x \ln \frac{T}{T_i} = \frac{1}{2} \alpha x^2 \Rightarrow T = T_i e^{\frac{\alpha x^2}{2C_x}}$$

La temperatura aumenta, no importa si contraemos o expandimos el resorte. Esto es porque

$$dW = f dx = \frac{\alpha T}{2} dx^2$$

Si empezamos en $x=0$, $dx^2 > 0$

$$\text{siempre} \Rightarrow dW > 0 \Rightarrow dU > 0 \Rightarrow dT > 0$$

$\uparrow$ Adiabático $\uparrow$
 $U = C_x T$

(d) Adiabático: $Q = 0$
 Libre: $W = 0$ } $\Rightarrow \Delta U = 0 \Rightarrow \Delta T = 0$
 $\uparrow$
 $U = C_V T$

$$\Delta S = S(T, a) - S(T, x) = \frac{1}{2} \alpha x^2 = \Delta S$$

$\Delta S > 0$ porque el proceso es adiabático y no cuasiestático.

2. (a) Estados:
- 1 igual
 - 2 ventaja Nadal
 - 3 " Djokovic
 - 4 gana Nadal
 - 5 " Djokovic

$$M = \begin{matrix} & \begin{matrix} 1 & 2 & 3 & 4 & 5 \end{matrix} \\ \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \end{matrix} & \begin{pmatrix} q & p & q & 0 & 0 \\ q & 0 & 0 & p & 0 \\ p & 0 & 0 & 0 & q \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \end{matrix}$$

4 y 5 absorbentes $\Rightarrow$ No ergódico.

Desde cualquier estado transitorio se puede acceder a un absorbente

$\Rightarrow$ Absorbente

La matriz ya está en forma canónica.

$$Q = \begin{pmatrix} 0 & p & q \\ q & 0 & 0 \\ p & 0 & 0 \end{pmatrix}$$

$$R = \begin{pmatrix} 0 & 0 \\ p & 0 \\ 0 & q \end{pmatrix}$$

$$(b) \quad I - Q = \begin{pmatrix} 1 & -p & -q \\ -q & 1 & 0 \\ -p & 0 & 1 \end{pmatrix}$$

$$\Rightarrow \det(I - Q) = \begin{vmatrix} 1 & -p \\ -q & 1 \end{vmatrix} - q \begin{vmatrix} -q & 1 \\ -p & 0 \end{vmatrix}$$

$$= 1 - pq - pq = 1 - 2pq$$

$$\Rightarrow N = (I - Q)^{-1} = \frac{1}{1 - 2pq} \begin{pmatrix} 1 & p & q \\ q & 1 - pq & q^2 \\ p & p^2 & 1 - pq \end{pmatrix} = N$$

$$\Rightarrow \langle t \rangle = \frac{1}{1 - 2pq} (1 + p + q)$$

Empezamos

en el estado 1 (iguales)

Si $p=1$, $\langle t \rangle = 2$ ✓

(Si Nadal gana quiere el game termine en 2 puntos)

$$\text{Si } p = \frac{1}{2}, \quad \langle t \rangle = \frac{1}{1 - \frac{2}{4}} \times 2 = \frac{2}{1/2} = 4 \quad \checkmark$$

(si ninguno gana sigue el game tarde más en terminar)

$$(c) \quad B = NR = \frac{1}{1 - 2pq} \begin{pmatrix} 1 & p & q \\ q & 1 - pq & q^2 \\ p & p^2 & 1 - pq \end{pmatrix} \begin{pmatrix} 0 & 0 \\ p & 0 \\ 0 & q \end{pmatrix}$$

$$= \frac{1}{1 - 2pq} \begin{pmatrix} p^2 & q^2 \\ p(1 - pq) & q^3 \\ p^3 & q(1 - pq) \end{pmatrix}$$

$$P(\text{Nadal gana el game}) = \frac{p^2}{1 - 2pq}$$

Empezamos

en el estado

1 (igual)

$$\text{Si } p = 1, \quad P(\text{Nadal}) = 1 \quad \checkmark$$

$$\text{Si } p = \frac{1}{2}, \quad P(\text{Nadal}) = \frac{\frac{1}{4}}{1 - \frac{2}{4}} = \frac{\frac{1}{4}}{\frac{2}{4}} = \frac{1}{2} \quad \checkmark$$

3. (a) $E = \hbar \omega_n \pm \epsilon$

↑

Eligiendo origen energías

(b) $Z = Z_1^N$ $Z_1 = Z_{1, \text{tras}} Z_{1, \text{int}}$

$$Z_{1, \text{tras}} = \sum_{n=0}^{\infty} e^{-\beta \hbar \omega n} = \frac{1}{1 - e^{-\beta \hbar \omega}}$$

$$Z_{1, \text{int}} = e^{\beta \epsilon} + e^{-\beta \epsilon} = 2 \cosh(\beta \epsilon)$$

$$\Rightarrow Z(\beta, N) = \left[\frac{2 \cosh(\beta \epsilon)}{1 - e^{-\beta \hbar \omega}} \right]^N$$

(c) $\langle E_1 \rangle = - \frac{\partial}{\partial \beta} \log Z_1 = - \frac{\partial}{\partial \beta} \log Z_{1, \text{tras}} - \frac{\partial}{\partial \beta} \log Z_{1, \text{int}}$

$$= \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} - \epsilon \tanh(\beta \epsilon)$$

$$\Rightarrow \langle E_1 \rangle = \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1} - \epsilon \tanh(\beta \epsilon)$$

$$c = \frac{d\langle E \rangle}{dT} = - \frac{1}{kT^2} \frac{d\langle E \rangle}{d\beta} = - \frac{1}{kT^2} \left\{ - \frac{\hbar\omega}{(e^{\beta\hbar\omega} - 1)^2} \hbar\omega e^{\beta\hbar\omega} - \frac{e^2}{\cosh^2(\beta e)} \right\}$$

$$\begin{aligned} \tanh'(x) &= \\ &= 1 - \tanh^2 x \\ &= \frac{1}{\cosh^2 x} \end{aligned}$$

$$\Rightarrow c(T) = \frac{1}{kT^2} \left[(\hbar\omega)^2 \frac{e^{\frac{\hbar\omega}{kT}}}{(e^{\frac{\hbar\omega}{kT}} - 1)^2} + \frac{e^2}{\cosh^2\left(\frac{E}{kT}\right)} \right]$$

$$(d) \quad kT \gg \hbar\omega, \quad e \Rightarrow \beta\hbar\omega, \quad \beta e \ll 1$$

$$\Rightarrow \langle E \rangle \approx \hbar\omega \left[\frac{1}{\beta\hbar\omega} - \frac{1}{2} + \frac{\beta\hbar\omega}{12} \right] - \beta e^2$$

$$\approx \frac{1}{\beta} - \frac{\hbar\omega}{2} + \beta \left[\frac{(\hbar\omega)^2}{12} - e^2 \right]$$

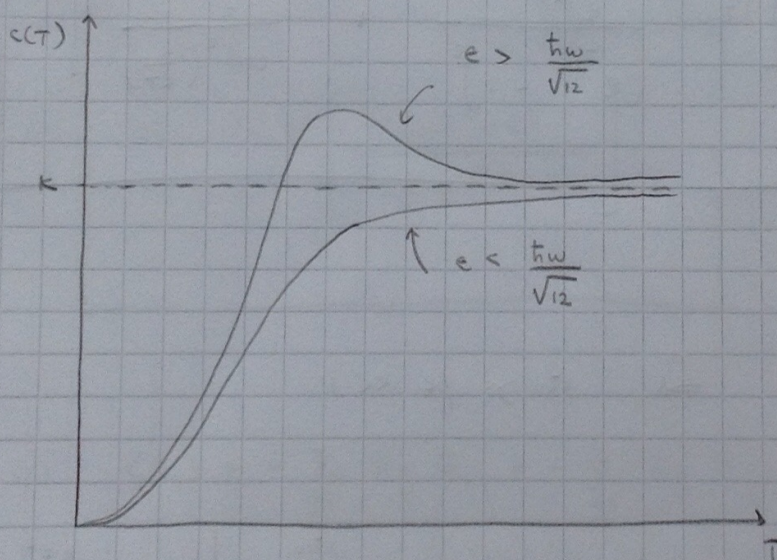
$$= kT - \frac{\hbar\omega}{2} + \frac{1}{kT} \left[\frac{(\hbar\omega)^2}{12} - e^2 \right] \approx \langle E \rangle$$

$$\Rightarrow c(T) \approx k + \frac{1}{\kappa T^2} \left[e^2 - \frac{(\hbar\omega)^2}{12} \right]$$

↑
Altes
temperaturas

Bajas temperaturas $\kappa T \ll \hbar\omega, \epsilon$

$$\Rightarrow c(T) \approx \frac{1}{\kappa T^2} \left[(\hbar\omega)^2 e^{-\frac{\hbar\omega}{\kappa T}} + (2\epsilon)^2 e^{-\frac{2\epsilon}{\kappa T}} \right]$$



$$(e) \quad f = \left(\frac{\partial F}{\partial N} \right)_T$$

$$F = -\kappa T \log Z = -N \kappa T \log Z_1(T)$$

$$\Rightarrow f = -\kappa T \log Z_1 = -\kappa T \log \left[\frac{2 \cosh(\beta\epsilon)}{1 - e^{-\beta\hbar\omega}} \right] = f$$