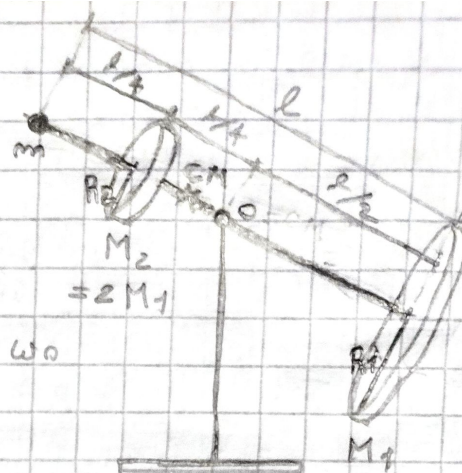


Entrega 5

Antes de agregar la masa m , hay balance de torques gravitatorios respecto del centro del eje O



$\underline{C.I.}: \theta(0) = \theta_0, \varphi(0) = \psi(0) = \dot{\theta}(0) = 0, \dot{\varphi}(0) = \omega_0$
 con $\Omega_3 = \Omega_0 \gg \gamma$, y $\dot{\varphi} \ll 1$

a) $I_1^{\circ} = I_2^{\circ} = I^{\circ} = ?$, $I_3^{\circ} = I_3 = ?$

Los momentos de inercia de los dos discos y de la masa m respecto de sus respectivos CM son:

$M_1) I_1^{CM_1} = I_2^{CM_1} = I^{CM_1} = \frac{M_1 R_1^2}{4}$, $I_3^{CM_1} = \frac{M_1 R_2^2}{2}$

$M_2) I^{CM_2} = \frac{M_2 R_2^2}{4}$, $I_3^{CM_2} = \frac{M_2 R_1^2}{2}$

$m) I^{CMm} = 0$, $I_3^{CMm} = 0$

Ahora los escribo en O usando el Teo de Steiner: (como O y los CM's estan en el eje principal \hat{z} , los I_3 van a ser iguales respecto de O y de los CM.)

$M_1) I^{O_1} = \frac{M_1 R_1^2}{4} - M_1 \left(\frac{l}{2}\right)^2 = M_1 \frac{R_1^2 - l^2}{4}$, $I_3^{O_1} = \frac{M_1 R_2^2}{2}$

$M_2) I^{O_2} = \frac{M_2 R_2^2}{4} + M_2 \left(\frac{l}{4}\right)^2 = M_2 \frac{R_2^2 + (l/2)^2}{4}$, $I_3^{O_2} = \frac{M_2 R_1^2}{2}$

$m) I^{Omm} = 0 + m \left(\frac{l}{2}\right)^2 = \frac{m l^2}{4}$, $I_3^{Omm} = 0$

Ahora, los momentos propios de inercia respecto de O van a ser la suma de los momentos de cada masa:

además $M_2 = 2M_1$

$$I = \frac{M_1}{4} \left(R_1^2 + 2R_2^2 - \frac{l^2}{2} \right) + \frac{m l^2}{4}$$

$$I_3 = \frac{M_1}{2} (R_1^2 + 2R_2^2)$$

b) $T(\varphi, \theta, \psi, \dot{\varphi}, \dot{\theta}, \dot{\psi}) = ?$ como O es punto fijo, T es solo de rotacion

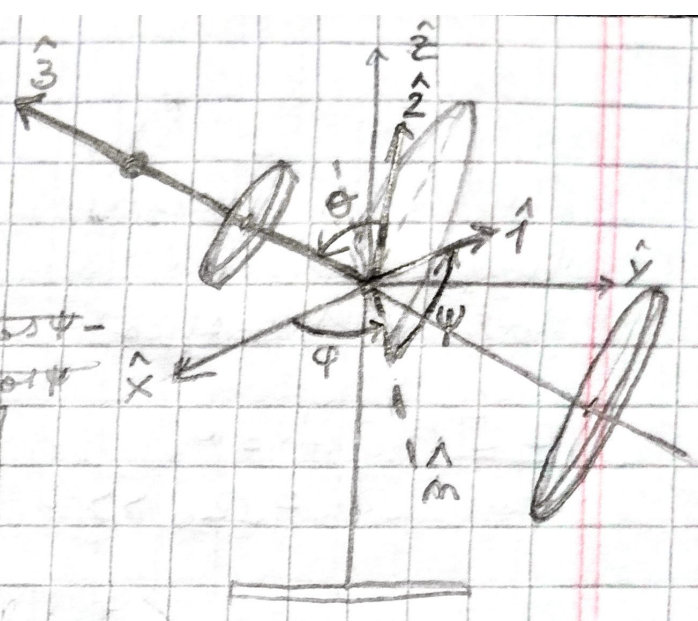
$$\Rightarrow T = T_{rot} = \frac{I}{2} (\Omega_1^2 + \Omega_2^2) + \frac{I_3}{2} \Omega_3^2 \quad (\vec{\Omega} = (\Omega_1, \Omega_2, \Omega_3))$$

ya que \underline{I}° es diagonal

$$\rightarrow \begin{cases} \Omega_1 = \dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi \\ \Omega_2 = \dot{\phi} \sin \theta \cos \psi - \dot{\theta} \sin \psi \\ \Omega_3 = \dot{\phi} \cos \theta + \dot{\psi} \end{cases}$$

$$\Rightarrow \Omega_1^2 + \Omega_2^2 = \dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2 + 2\dot{\phi}\dot{\theta} \sin \theta \sin \psi \cos \psi - 2\dot{\phi}\dot{\theta} \sin \theta \cos \psi \sin \psi$$

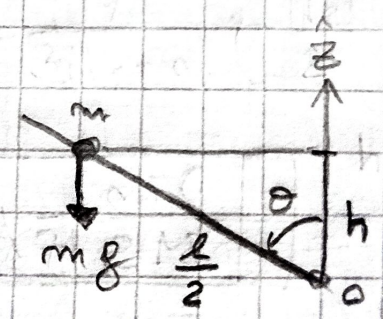
$$\Rightarrow T = \frac{I}{2} (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{I_3}{2} (\dot{\phi} \cos \theta + \dot{\psi})^2$$



Como el sistema estaba en equilibrio de torques antes de agregar la masa m , la fuerza gravitatoria va a estar sobre la misma.

V está dada por esta fuerza gravitatoria:

$$V = m g \frac{l}{2} \cos \theta$$



$$c) \mathcal{L} = \frac{I}{2} (\dot{\varphi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{I_3}{2} (\dot{\varphi} \cos \theta + \dot{\psi})^2 - mg \frac{l}{2} \cos \theta$$

EL

$$\varphi) \frac{\partial \mathcal{L}}{\partial \varphi} = 0 \rightarrow \varphi \text{ cíclico}$$

$$\Rightarrow I \dot{\varphi} \sin^2 \theta + I_3 (\dot{\varphi} \cos \theta + \dot{\psi}) \cos \theta \equiv P_\varphi \text{ de}$$

$$\psi) \frac{\partial \mathcal{L}}{\partial \psi} = 0 \rightarrow \psi \text{ cíclico}$$

$$\Rightarrow I_3 (\dot{\varphi} \cos \theta + \dot{\psi}) \equiv P_\psi \text{ de} \rightarrow P_\varphi = I \dot{\varphi} \sin^2 \theta + P_\psi \cos \theta$$

$$\Rightarrow \dot{\varphi} = \frac{P_\varphi - P_\psi \cos \theta}{I \sin^2 \theta} \rightarrow \text{derivado:}$$

$$\begin{aligned} \Rightarrow \ddot{\varphi} &= - \frac{2 P_\psi \dot{\theta} \sin \theta \cos \theta}{I \sin^4 \theta} - \frac{P_\psi \dot{\theta} (-\sin^2 \theta - 2 \cos^2 \theta) \sin \theta}{\sin^4 \theta} \\ &= \frac{P_\psi (\sin^2 \theta + 2 \cos^2 \theta) - 2 P_\psi \cos \theta \dot{\theta}}{I \sin^3 \theta} = I \ddot{\varphi} \sin^2 \theta + P_\psi \cos \theta \\ &= \frac{P_\psi \sin 2\theta + 2 P_\psi \cos^2 \theta \dot{\theta} - 2 I \dot{\varphi} \sin \theta \cos \theta - 2 P_\psi \cos^2 \theta \dot{\theta}}{I \sin^3 \theta} \\ &= \frac{P_\psi - 2 I \dot{\varphi} \cos \theta}{I \sin \theta} \dot{\theta} = \frac{I_3 (\dot{\varphi} \cos \theta + \dot{\psi}) - 2 I \dot{\varphi} \cos \theta}{I \sin \theta} \dot{\theta} \end{aligned}$$

$$\Rightarrow \ddot{\varphi} = \frac{(I_3 - 2I) \dot{\varphi} \cos \theta + I_3 \dot{\psi}}{I \sin \theta} \dot{\theta}$$

Para $\dot{\psi}$ derivos $P_\psi = I_3 (\dot{\varphi} \cos \theta + \dot{\psi})$:

$$\Rightarrow 0 = I_3 (\ddot{\varphi} \cos \theta - \dot{\varphi} \dot{\theta} \sin \theta + \ddot{\psi}) \Rightarrow \ddot{\psi} = \dot{\varphi} \dot{\theta} \sin \theta + \ddot{\varphi} \cos \theta$$

$$\Rightarrow \ddot{\psi} = \dot{\varphi} \dot{\theta} \sin \theta - \frac{(I_3 - 2I) \dot{\varphi} \cos^2 \theta + I_3 \dot{\psi} \cos \theta}{I \sin \theta} \dot{\theta}$$

$$e) \frac{d}{dt} [I \dot{\theta}]^{EL} = I \dot{\varphi}^2 \sin \theta \cos \theta - I_3 (\dot{\varphi} \cos \theta + \dot{\psi}) \dot{\varphi} \sin \theta + mg \frac{l}{2} \sin \theta$$

$$\Rightarrow \ddot{\theta} = \frac{I - I_3}{I} \dot{\varphi}^2 \sin \theta \cos \theta - \frac{I_3}{I} \dot{\psi} \dot{\varphi} \sin \theta + \frac{mg l}{2 I} \sin \theta$$

d) $\theta \approx \theta_0, \dot{\theta}, \dot{\varphi}, \dot{\psi} \ll \omega$ tomblén $\dot{\psi} = \Omega_3 \equiv \Omega_0 \text{ de } \gg 1$
 $\approx \frac{m g l}{2 I}$
 deb $\varepsilon = \theta - \theta_0 \ll \Rightarrow \theta = \theta_0 + \varepsilon \Rightarrow \dot{\theta} = \dot{\varepsilon}$

$\begin{cases} \sin \theta \approx \sin \theta_0 + \varepsilon \cos \theta_0 \\ \cos \theta \approx \cos \theta_0 - \varepsilon \sin \theta_0 \end{cases}$ cuando tenga $\dot{\theta}, \dot{\varphi}, \dot{\psi}$ o m por un seno o un coseno, solo va a sobrevivir el término sin ε , ya que el término con ε queda de 2do orden
 (Ej: $\dot{\varphi} \sin \theta \approx \dot{\varphi} \sin \theta_0 + \dot{\varphi} \varepsilon \cos \theta_0$
 $\approx \dot{\varphi} \sin \theta_0$)

$$3) I \ddot{\varphi} \sin \theta_0 = (I_3 - 2I) \dot{\varphi} \dot{\theta} \cos \theta + I_3 \ddot{\varphi} \dot{\theta} \stackrel{= \Omega_0}{\rightarrow 0}$$

$$\Rightarrow \boxed{\ddot{\varphi} - \frac{I_3 \Omega_0}{I \sin \theta_0} \dot{\theta} = 0} \quad (1)$$

$$4) \ddot{\psi} = \dot{\varphi} \dot{\theta} \sin \theta - \ddot{\varphi} \cos \theta = -\ddot{\varphi} \cos \theta_0$$

$$\Rightarrow \boxed{\ddot{\psi} + \frac{I_3 \Omega_0 \cos \theta_0}{I \sin \theta_0} \dot{\theta} = 0} \quad (2)$$

$$5) \ddot{\theta} = \frac{I - I_3}{I} \dot{\varphi} \dot{\theta} \cos \theta - \frac{I_3}{I} \dot{\varphi} \dot{\theta} \sin \theta + \frac{mgl}{2I} \sin \theta$$

$$\Rightarrow \boxed{\ddot{\theta} + \frac{I_3 \Omega_0 \sin \theta_0}{I} \dot{\varphi} = \frac{mgl}{2I} \sin \theta_0} \quad (3)$$

c) Plazo (3) - i (1):

$$I \ddot{\theta} - i I \dot{\varphi} \sin \theta_0 + I_3 \Omega_0 \dot{\varphi} \sin \theta_0 + i I_3 \Omega_0 \dot{\theta} = \frac{mgl}{2} \sin \theta_0$$

$$\Rightarrow I(\ddot{\theta} - i \dot{\varphi} \sin \theta_0) + I_3 \Omega_0 \left(\frac{i}{2} (\dot{\varphi} \sin \theta_0 + i \dot{\theta}) \right) = \frac{mgl}{2} \sin \theta_0$$

$$\Rightarrow \boxed{I(\ddot{\theta} - i \dot{\varphi} \sin \theta_0) + i I_3 \Omega_0 (\dot{\theta} - i \dot{\varphi} \sin \theta_0) = \frac{mgl}{2} \sin \theta_0}$$

$$\text{defino } \boxed{\lambda \equiv \dot{\theta} - i \dot{\varphi} \sin \theta_0} \Rightarrow \boxed{\dot{\lambda} = \ddot{\theta} - i \ddot{\varphi} \sin \theta_0} \Rightarrow \boxed{\ddot{\lambda} = \ddot{\theta} - i \ddot{\varphi} \sin \theta_0}$$

$$\Rightarrow \boxed{I \ddot{\lambda} + i I_3 \Omega_0 \dot{\lambda} = \frac{mgl}{2} \sin \theta_0}$$

La solución es $\lambda = \lambda_h + \lambda_p$

$$1) \text{ Para } \lambda_p \text{ hido } \boxed{\ddot{\lambda}_p = 0} \Rightarrow 0 + i I_3 \Omega_0 \dot{\lambda}_p = \frac{mgl}{2} \sin \theta_0$$

$$\Rightarrow \boxed{\dot{\lambda}_p = -i \frac{mgl/2}{I_3 \Omega_0} \sin \theta_0 = \text{cte}}$$

\Rightarrow Proposito:

$$\rightarrow \text{como } \boxed{\tilde{C} = C - i t \sin \theta_0 \in \mathbb{C}}$$

$$\lambda_p = -i \frac{mgl/2}{I_3 \Omega_0} \sin \theta_0 t + \tilde{C}$$

$\equiv \omega_p$: velocidad de precesión de Larmor
de tiempo pesado

$$\Rightarrow \boxed{\lambda_p = -i \omega_p \sin \theta_0 t + \tilde{C}}$$

$$\omega_p = \frac{mgl/2}{I_3 \Omega_0}$$

$$\tilde{C} = C - i t \sin \theta_0$$

→ •) La homogénea lo iguala a 0:

$$I \ddot{\lambda}_h + i I_3 \Omega_0 \dot{\lambda}_h = 0 \Rightarrow \ddot{\lambda}_h + i \frac{I_3}{I} \Omega_0 \dot{\lambda}_h = 0$$

ω_L : velocidad de precesión de Larmour

La solución es de la forma:

$$\lambda_h = a e^{i(\omega_L t - \varphi_0)} \text{sen } \theta_0$$

ó $\lambda_h = \lambda_0 e^{i \omega_L t}$ con $\lambda_0 = a e^{-i \varphi_0} \text{sen } \theta_0$

$$\Rightarrow \lambda = a \text{sen } \theta_0 e^{i(\omega_L t - \varphi_0)} - i \omega_p \text{sen } \theta_0 t + C - i f \text{sen } \theta_0$$

$$= \cos(\omega_L t - \varphi_0) + i \text{sen}(\omega_L t - \varphi_0) \text{ (Euler)}$$

$$\Rightarrow \lambda = \underbrace{a \text{sen } \theta_0 \cos(\omega_L t - \varphi_0)}_{=\vartheta} + C - i \underbrace{(\omega_p t + a \text{sen}(\omega_L t - \varphi_0) + f)}_{=\varphi} \text{sen } \theta_0$$

y queda $\lambda = \vartheta - i \varphi \text{sen } \theta_0$

$$\Rightarrow \begin{cases} \vartheta(t) = a \text{sen } \theta_0 \cos(\omega_L t - \varphi_0) + C \\ \varphi(t) = \omega_p t + a \text{sen}(\omega_L t - \varphi_0) + f \end{cases}$$

Falta hallar los dos $\{a, \varphi_0, C, f\}$ usando las

CI: $\varphi(0) = 0, \dot{\varphi}(0) = \omega_0$
 $\vartheta(0) = \vartheta_0, \dot{\vartheta}(0) = 0$

$$\vartheta(0) = a \text{sen } \theta_0 \cos(-\varphi_0) + C = \vartheta_0 \Rightarrow C = \vartheta_0 - a \text{sen } \theta_0 \cos \varphi_0$$

$$\varphi(0) = a \text{sen } \varphi_0 + f = 0 \Rightarrow f = -a \text{sen } \varphi_0$$

$$\dot{\vartheta}(0) = -a \text{sen } \theta_0 \omega_L \text{sen}(-\varphi_0) = 0 \Rightarrow a \omega_L \text{sen } \theta_0 \text{sen } \varphi_0 = 0$$

$$\dot{\varphi}(0) = \omega_p + a \omega_L \cos(-\varphi_0) = \omega_0 \Rightarrow a = \frac{\omega_0 - \omega_p}{\omega_L \cos \varphi_0}$$

$$\Rightarrow \frac{\omega_0 - \omega_p}{\omega_L \cos \varphi_0} \text{sen } \theta_0 \text{sen } \varphi_0 = \frac{\omega_0 - \omega_p}{\omega_L} \text{sen } \theta_0 \text{tg } \varphi_0 = 0$$

$$\Rightarrow \text{tg } \varphi_0 = 0 \Rightarrow \varphi_0 = m\pi \quad m \in \mathbb{Z} \Rightarrow f = 0$$

$$\Rightarrow a = \pm \frac{\omega_0 - \omega_p}{\omega_L}, \quad C = \vartheta_0 \mp a \text{sen } \theta_0 \Rightarrow \cos \varphi_0 = \pm 1$$

$$\Rightarrow \vartheta(t) = \pm \frac{\omega_0 - \omega_p}{\omega_L} \text{sen } \theta_0 \cos(\omega_L t - m\pi) + \vartheta_0 - \frac{\omega_0 - \omega_p}{\omega_L} \text{sen } \theta_0$$

$$\varphi(t) = \omega_p t \pm \frac{\omega_0 - \omega_p}{\omega_L} \text{sen}(\omega_L t - m\pi)$$

elijo $m=1 \Rightarrow \varphi_0 = \pi \rightarrow \begin{cases} \cos \varphi_0 = -1 \\ \cos(\omega_L - \varphi_0) = -\cos(\omega_L) \\ \sin(\omega_L - \varphi_0) = -\sin(\omega_L) \end{cases}$

$$\theta(t) = \frac{\omega_p - \omega_0}{\omega_L} \sin \theta_0 (1 - \cos(\omega_L t)) + \theta_0$$

$$\varphi(t) = \omega_p t - \frac{\omega_p - \omega_0}{\omega_L} \sin(\omega_L t)$$

$$\dot{\varphi}(t) = \omega_p - (\omega_p - \omega_0) \cos(\omega_L t)$$

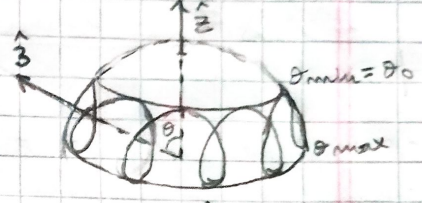
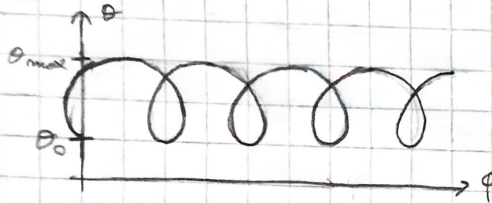
Ahora grafico la variación de θ (nutación) cuando el giroscopo gira en φ .

Observaciones:

- \rightarrow Si $\omega_p - \omega_0 > 0$ el giroscopo sube inicialmente, $\theta_0 = \theta_{min}$
- \rightarrow Si $\omega_p - \omega_0 < 0$ el giroscopo baja inicialmente, $\theta_0 = \theta_{max}$
- \rightarrow Si $\omega_p - \omega_0 = 0$ no hay nutación (ni sube ni baja)
- \rightarrow Si $|\omega_p - \omega_0| > \omega_p$ $\dot{\varphi}$ cambia de signo

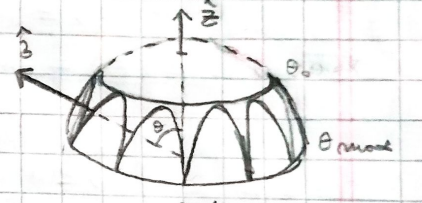
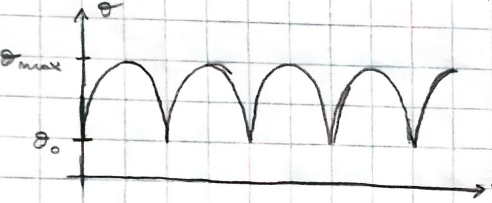
$\omega_0 = -\omega_p$:

$$\begin{cases} \omega_p + \omega_0 > 0 \\ |\omega_p + \omega_0| > \omega_p \end{cases}$$



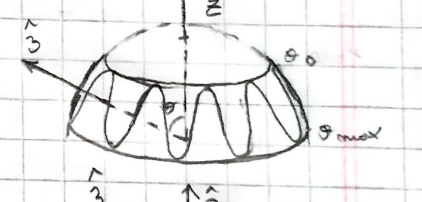
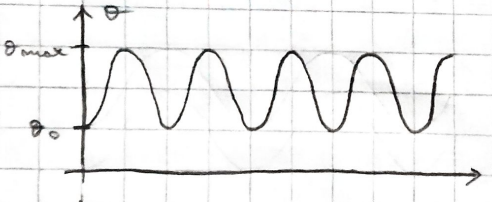
$\omega_0 = 0$:

$$\begin{cases} \omega_p - 0 > 0 \\ |\omega_p - 0| = \omega_p \end{cases}$$



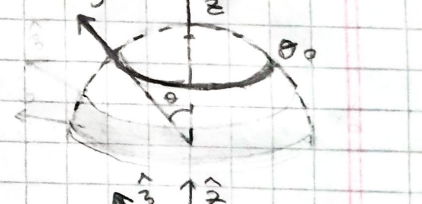
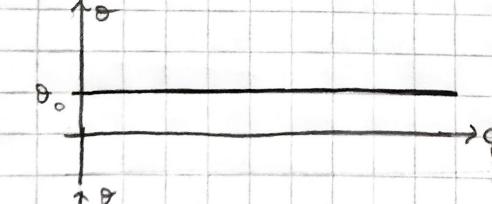
$\omega_0 = \frac{\omega_p}{2}$:

$$\begin{cases} \omega_p - \frac{\omega_p}{2} > 0 \\ |\omega_p - \frac{\omega_p}{2}| < \omega_p \end{cases}$$



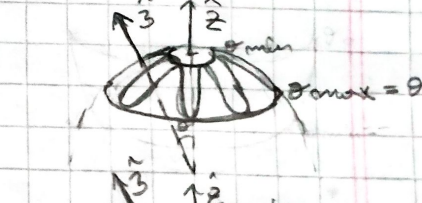
$\omega_0 = \omega_p$:

$$\begin{cases} \omega_p - \omega_p = 0 \\ |\omega_p - \omega_p| < \omega_p \end{cases}$$



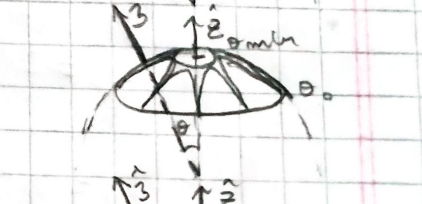
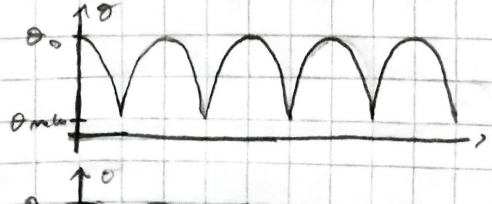
$\omega_0 = \frac{3}{2}\omega_p$:

$$\begin{cases} \omega_p - \frac{3}{2}\omega_p < 0 \\ |\omega_p - \frac{3}{2}\omega_p| < \omega_p \end{cases}$$



$\omega_0 = 2\omega_p$:

$$\begin{cases} \omega_p - 2\omega_p < 0 \\ |\omega_p - 2\omega_p| = \omega_p \end{cases}$$



$\omega_0 = 3\omega_p$:

$$\begin{cases} \omega_p - 3\omega_p < 0 \\ |\omega_p - 3\omega_p| > \omega_p \end{cases}$$

