

# Relatividad General



## \* Convenciones

i) Índices  $\begin{cases} \text{griegos: } \mu = 0, 1, 2, 3 \\ \text{latinos: } i = 1, 2, 3 \end{cases}$   
 espacio-temporales  
 espaciales

ii) Signatura de la métrica:

$$g_{\mu\nu} = \begin{pmatrix} 1 & & & \\ & -1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \quad (\text{"mostly minus"})$$

## Guía 1: Relatividad especial

$$g_{\mu\nu} = \eta_{\mu\nu} = \begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \rightarrow \text{métrica de Minkowski}$$

$\eta_{ij} = -\delta_{ij}$

$\eta_{00} = +1, \eta_{0i} = 0, \eta_{ii} = -1, \eta_{ij} = 0 (i \neq j)$

\* Transf. de Lorentz  $\Lambda^\mu_\nu$  tal que  $\Lambda^\mu_\nu \Lambda^\alpha_\beta \eta_{\mu\alpha} = \eta_{\nu\beta}$

$$\Lambda^\sigma_\nu \Lambda^\alpha_\beta \eta_{\sigma\alpha} = \sum_{\sigma=0}^3 \sum_{\alpha=0}^3 \Lambda^\sigma_\nu \Lambda^\alpha_\beta \eta_{\sigma\alpha} = \sum_{\sigma=0}^3 (\Lambda^\sigma_\nu \Lambda^0_\beta \eta_{\sigma 0} + \Lambda^\sigma_\nu \Lambda^1_\beta \eta_{\sigma 1} + \Lambda^\sigma_\nu \Lambda^2_\beta \eta_{\sigma 2} + \Lambda^\sigma_\nu \Lambda^3_\beta \eta_{\sigma 3})$$

\* Covariante  $V$ : objeto con 4 componentes  $V^\mu$  ;  $\mu = 0, 1, 2, 3$  asociadas a cierto  $S$

$$\Rightarrow \text{En } S': V'^\mu = \Lambda^\mu_\nu V^\nu$$

Ejemplo:  $x$ , con  $x^\mu = (ct, x, y, z)$

y una transf. de Lorentz manda  $x^\mu \mapsto x'^\mu = \Lambda^\mu_\nu x^\nu$

\* Transf. de Lorentz.  $\rightarrow$  rotaciones

y una Trans. de Lorentz manda  $x^\mu \mapsto x'^\mu = \Lambda^\mu_\nu x^\nu$

\* Trans de Lorentz:  $\begin{cases} \rightarrow \text{rotaciones} \\ \rightarrow \text{boosts} \end{cases}$

ej:  $ct' = \gamma (ct - \beta x) ; y' = y$   
 $x' = \gamma (x - \beta ct) ; z' = z$

$$\left. \begin{matrix} \Lambda^0_0 = \gamma \\ \Lambda^1_1 = \gamma \\ \Lambda^0_1 = -\beta\gamma/c \end{matrix} \right\} \leftarrow \Lambda^\mu_\nu = \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

\*  $\eta$  nos da un producto invariante de Lorentz:

$$V \cdot W := V^\mu W^\nu \eta_{\mu\nu} = V^0 W^0 - V^1 W^1 - V^2 W^2 - V^3 W^3 =$$

$$= V^0 W^0 - \vec{V} \cdot \vec{W}$$

Venimos que es invariante:

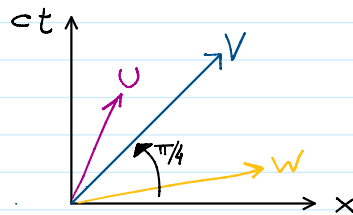
$$V'^\mu W'^\nu \eta_{\mu\nu} \stackrel{\text{def. de 4-vector}}{=} \Lambda^\mu_\beta V^\beta \Lambda^\nu_\alpha W^\alpha \eta_{\mu\nu}$$

$$= V^\beta W^\alpha (\Lambda^\mu_\beta \Lambda^\nu_\alpha \eta_{\mu\nu}) \stackrel{\text{def. } \Lambda}{=} V^\beta W^\alpha \eta_{\beta\alpha} = V \cdot W \quad \checkmark$$

\* Separamos en 3 casos según la norma de  $V$ ,  $\|V\|^2 = V \cdot V$ :

- i)  $V \cdot V > 0 \rightarrow V$  es time-like o tipo-tiempo
- ii)  $V \cdot V = 0 \rightarrow V$  es null-like o tipo-nulo o nulo
- iii)  $V \cdot V < 0 \rightarrow V$  es space-like o tipo-espacio

\* Diagrama espacio-temporal:



$$V \cdot V = (V^0)^2 - (V^1)^2 = 0$$

$$|V^0| = |V^1|$$

$V$  nulo

$U$  time-like

$W$  space-like

Obs:  $U \cdot U = (U^0)^2 - \vec{U}^2 > 0 \Rightarrow$  una buena descripción de  $U$

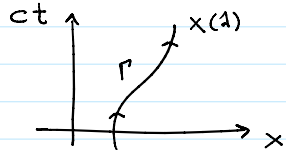
Obs:  $U \cdot U = (U^0)^2 - \vec{U}^2 > 0 \Rightarrow$  una buena descripción de  $U$

$$= \|U\|^2 (ch^2\beta - sh^2\beta) = \|U\|^2$$

es:  $U^0 = \|U\| ch\beta$   
 $U^x = \|U\| sh\beta$  *sho cosp*  
 $U^y = \|U\| sh\beta$  *sho ship*  
 $U^z = \|U\| sh\beta$  *sho*

ej: Para  $W$  intercambio  $ch\beta \leftrightarrow sh\beta \Rightarrow \|W\|^2 < 0$

Problema 1



$$\frac{dx^\mu(\lambda)}{d\lambda} = \gamma(v) (c \frac{dt}{d\lambda}, c \frac{dx}{d\lambda}, 0, v)$$

$$\frac{d(ct)}{d\lambda} = \gamma c \frac{dt}{d\lambda}$$

a) Veamos  $g'$  a  $\lambda = \tau$  (tiempo propio):  $ds^2 = c^2 dt^2 - dx^2 \stackrel{t=\tau \leftrightarrow d\tau=0}{=} c^2 d\tau^2$

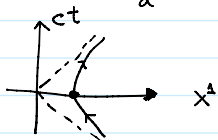
$$\Rightarrow ds^2|_{\tau} = c^2 dt^2 - dx^2|_{\tau} = c^2 \left(\frac{dt}{d\lambda} d\lambda\right)^2 - \left(\frac{dx}{d\lambda} d\lambda\right)^2 = d\lambda^2 \left(\frac{dct}{d\lambda} - \frac{dx}{d\lambda}\right)^2$$

$$= d\lambda^2 \left[ \gamma^2 c^2 ch^2 \frac{dt}{d\lambda} - \gamma^2 c^2 sh^2 \frac{dt}{d\lambda} - \gamma^2 v^2 \right] = d\lambda^2 \gamma^2 (c^2 - v^2) = d\lambda^2 \gamma^2 c^2 (1 - v^2/c^2) = c^2 d\lambda^2$$

$\lambda = \tau \checkmark$

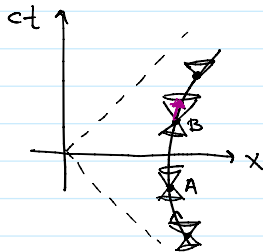
b)  $x^\mu(\lambda) = \gamma \frac{c^2}{a} (sh \frac{a\lambda}{c}, ch \frac{a\lambda}{c}, ct, \frac{v}{c} \cdot \frac{a\lambda}{c})$

$$(x^0)^2 - (x^1)^2 = \gamma^2 \frac{c^4}{a^2} (sh^2 \frac{a\lambda}{c} - ch^2 \frac{a\lambda}{c}) = -\gamma^2 \frac{c^4}{a^2} < 0$$



$$\frac{dt}{dx} = \frac{dt/d\lambda}{dx/d\lambda} = \frac{1}{\frac{v}{c} \frac{a\lambda}{c}} > 0 \quad (\text{ej, si } \lambda > 0)$$

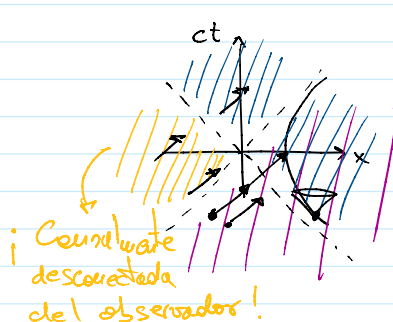
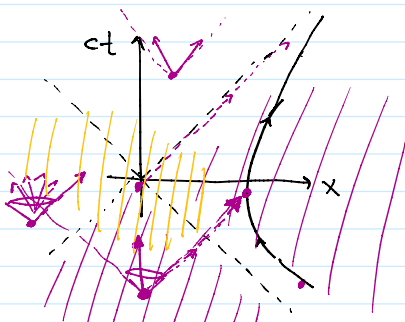
c)



$$\left\| \frac{dx^\mu}{d\lambda} \right\|^2 = c^2 > 0 \rightarrow \frac{dx^\mu}{d\lambda} \text{ es timelike}$$

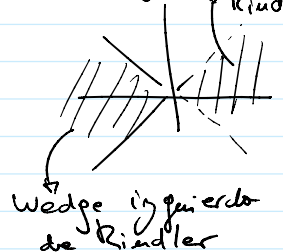
$$d(A,B) = \int_A^B ds = \int_A^B \sqrt{\left| \frac{dx^\mu}{d\lambda} \right|^2} d\lambda > 0$$

d)

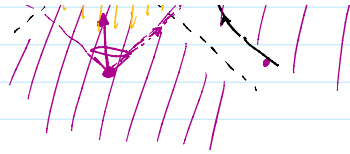


¡Coordenada descubierta del observador!

Wedge derecho de Rindler



Wedge izquierdo de Rindler



! Conmutate descometada del observador!

Wedge izquierda de Rindler

### Problema 3 (c=1)

a)  $U^\alpha = \frac{dx^\alpha}{d\tau} \rightarrow$  cuadrivelocidad  $\Rightarrow U^\alpha = \gamma(1, \vec{u})$  ya que  $\gamma(u)$

$$U^\alpha = \left( \frac{dt}{d\tau}, \frac{d\vec{x}}{d\tau} \right) = \left( \gamma, \frac{d\vec{x}}{dt} \frac{dt}{d\tau} \right) = \gamma(1, \vec{u})$$

$$dt^2 - d\vec{x}^2 = d\tau^2 = dt^2(1 - \frac{d\vec{x}^2}{dt^2}) = dt^2 \gamma^{-2} \rightarrow \frac{dt}{d\tau} = \gamma$$

$$a^\alpha = \frac{dU^\alpha}{d\tau} = \frac{dt}{d\tau} \frac{d}{dt} [\gamma(1, \vec{u})] = \gamma \left( \dot{\gamma}, \dot{\gamma} \vec{u} + \gamma \vec{a} \right)$$

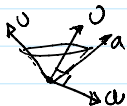
$$= \gamma \left( \vec{u} \cdot \vec{a} \gamma^3, \gamma \vec{a} + \vec{u} (\vec{u} \cdot \vec{a}) \gamma^3 \right)$$

b)  $U \cdot U = \gamma^2 (1 - \vec{a}^2) = 1$

c)  $U \cdot a = U^\alpha a^\beta \eta_{\alpha\beta} = U^0 a^0 - \vec{U} \cdot \vec{a} = \gamma^2 (\vec{a} \cdot \vec{a} \gamma^3 - \vec{u} \cdot (\gamma \vec{a} + \vec{u} (\vec{u} \cdot \vec{a}) \gamma^3))$

$$= \gamma^5 [\vec{a} \cdot \vec{a} - \sqrt{1 - \vec{a}^2} \vec{u} \cdot \vec{a}] = \gamma^5 (\vec{a} \cdot \vec{a}) [1 - \vec{a}^2 - \gamma^2] = 0 \Rightarrow a \cdot a < 0$$

a es tipo-espacio



$$\cos \alpha = \frac{U \cdot a}{\|U\| \|a\|}$$

d)  $a \cdot a = a^\alpha a^\beta \eta_{\alpha\beta} = (a^0)^2 - \vec{a}^2 = \dots = -\gamma^4 (\vec{a} \cdot \vec{a})^2 \left[ \gamma^2 + \frac{\vec{a}^2}{(\vec{a} \cdot \vec{a})^2} \right]$

$$\mu_{\parallel} = \frac{\vec{a} \cdot \vec{a} / |\vec{a}|}{\gamma^4 \vec{a}^2} = -\gamma^4 \vec{a}^2 \left[ 1 + \frac{\mu_{\parallel}^2}{1 - \vec{a}^2} \right] = -\frac{\gamma^6}{\gamma_{\parallel}^2} \vec{a}^2 \quad \text{con } \gamma_{\parallel} = [1 - (\vec{a}^2 \mu_{\parallel}^2)]^{-1/2}$$

$\vec{\alpha} := \frac{\gamma^3}{\gamma_{\parallel}} \vec{a}$  aceleración propia  $\rightarrow a \cdot a = -\vec{\alpha}^2$  invariante

obs: haciendo un boost con velocidad  $\vec{u}$ :  $\vec{\alpha} = \frac{\gamma^3}{\gamma_{\parallel}} \vec{a}' \gamma^{-3} = \frac{\vec{a}'}{\gamma_{\parallel}}$

y ahora si  $\mu_{\perp} = 0$ :  $\vec{\alpha} = \vec{a}' \Rightarrow \vec{\alpha}$  es la aceleración (instantánea) en el sistema del objeto (asumiendo  $\mu_{\perp} = 0$ )