

Más convenciones:  $\leftarrow$  índice arriba ( $A^\mu$ )  $\rightarrow$  "contravariante"  
 $\leftarrow$  índice abajo ( $A_\mu$ )  $\rightarrow$  "covariante"

\* Las relacionamos con la métrica:  $A_\mu = g_{\mu\nu} A^\nu \rightarrow$  automáticamente "subimos el índice":

$$A^\mu = g^{\mu\nu} A_\nu$$

↓  
 la inversa de la métrica  
 $(g^{\mu\nu} g_{\nu\alpha} = \delta^\mu_\alpha)$

Ej 8 y 9

Recordemos las ecs. de Maxwell:

$$\nabla \times \vec{E} + \partial_t \vec{B} = 0 \quad \nabla \times \vec{B} - \partial_t \vec{E} = \vec{j}$$

$$\nabla \cdot \vec{B} = 0 \quad \nabla \cdot \vec{E} = \rho$$

si defino  $\phi$  y  $\vec{A}$  tal que  $\vec{B} = \nabla \times \vec{A}$  y  $\vec{E} = -\nabla\phi - \dot{\vec{A}}$   
 puedo formar  $A_\alpha := (\phi, -\vec{A})$  ó  $A^\alpha = (\phi, \vec{A})$   
 $\hookrightarrow A_1 = -A^1, A_2 = -A^2, A_3 = -A^3$

Obs.:  $\vec{A}' = \vec{A} + \nabla\lambda$  "invariancia de gauge"

$$A'_\alpha = A_\alpha + \partial_\alpha \lambda = (\phi + \lambda, -\vec{A} + \nabla\lambda)$$

"Trans. de gauge"  $\phi' = \phi + \lambda$   $-\vec{A}' \Rightarrow \vec{A}' = \vec{A} - \nabla\lambda$

Def.  $F_{\alpha\beta} := \partial_\alpha A_\beta - \partial_\beta A_\alpha \Rightarrow F_{\alpha\beta} = \begin{pmatrix} 0 & -\partial_0 \vec{A} - \nabla\phi \\ -\vec{E} & -(\nabla \times \vec{A})^i & (\nabla \times \vec{A})^j \\ -F_{12} & 0 & -(\nabla \times \vec{A})^x \\ -F_{13} & -F_{23} & 0 \end{pmatrix} = \begin{pmatrix} 0 & -E & & \\ -E & 0 & -B^z & B^y \\ & B^z & 0 & -B^x \\ & -B^y & B^x & 0 \end{pmatrix}$

¡invariante de gauge!

$$\begin{aligned} \hookrightarrow F'_{\alpha\beta} &= \partial_\alpha A'_\beta - \partial_\beta A'_\alpha = \partial_\alpha (A_\beta + \partial_\beta \lambda) - \partial_\beta (A_\alpha + \partial_\alpha \lambda) \\ &= \partial_\alpha A_\beta - \partial_\beta A_\alpha + \partial_\alpha \partial_\beta \lambda - \partial_\beta \partial_\alpha \lambda = F_{\alpha\beta} \quad \checkmark \end{aligned}$$

(Ej.)  $F^{12} = g^{1\mu} g^{2\nu} F_{\mu\nu} = \eta^{1\mu} \eta^{2\nu} F_{\mu\nu} = +\delta_1^\mu \delta_2^\nu F_{\mu\nu} = F_{12} = -B^z$

F contravariante  $\Rightarrow F^{\alpha\beta} = g^{\alpha\mu} g^{\beta\nu} F_{\mu\nu} = g^{\alpha\mu} \left( \sum_{\nu=0}^3 g^{\beta\nu} F_{\mu\nu} \right)$

$$F^{\alpha\beta} = \begin{pmatrix} 0 & -E & & \\ E & 0 & -B^z & B^y \\ & B^z & 0 & -B^x \\ & -B^y & B^x & 0 \end{pmatrix}$$

$g = \eta$

Obs.:  $g^{\beta\nu} F_{\mu\nu} \stackrel{(?)}{=} g^{\beta\nu} F_{\nu\mu} = g^{\beta\nu} g_{\alpha\nu} F_\mu^\alpha = \delta_\mu^\beta F_\nu^\nu = F_\mu^\beta = g^{\beta\nu} F_{\mu\nu} \quad \checkmark$

o sea:  $A_\mu B^\mu = A^\mu B_\mu$  (da lo mismo cuál está arriba y cuál abajo si están formando contráctos)

Obs:  $F_{ij} = -\epsilon_{ijk} B^k$  con  $\epsilon_{23} = +1$ , o sea  $\epsilon_{132} = -1$

Las ecs. de Maxwell se reescriben como

$$\partial_\alpha F^{\alpha\beta} = j^\beta, \quad \partial_{[\alpha} F_{\beta\mu]} = 0$$

con  $j^\beta = (\rho, \vec{j})$

\* Fuerza de Lorentz:  $\vec{f} = q(\vec{E} + \vec{u} \times \vec{B})$

8 a) Cuadrifuerza:  $K^\alpha := q F^\alpha_\beta U^\beta = q F^{\alpha\beta} U_\beta$

Entonces, i)  $\alpha=0$ :  $K^0 = q F^0_\beta U^\beta = q \sum_{\beta\alpha} \frac{-\delta^0_\alpha E^i}{F^{0\alpha}} U^\beta$

$$= -q \sum_{\beta\alpha} \delta^0_\alpha E^i U^\beta = -q \sum_{\beta i} E^i U^\beta = +q \sum_{\beta} E^i U^\beta$$

$$= q E^i U^i = q \vec{E} \cdot \vec{U} = q \gamma \vec{E} \cdot \vec{u}$$

ii)  $\alpha=i$ :  $K^i = q F^i_\beta U^\beta = q \sum_{\alpha\beta} F^{i\alpha} U^\beta$

$$= q (F^{i0} U^0 - F^{ij} U^j) = q (E^i \gamma - (-\epsilon^{ijk} B^k) \gamma u^j)$$

$F^{i0} = E^i$   
 $F^{ij} = -\epsilon^{ijk} B^k$

$$= q \gamma (E^i + \epsilon^{ijk} u^j B^k) = q \gamma (\vec{E} + \vec{u} \times \vec{B})^i = \gamma f^i$$

$\vec{k} = \gamma \vec{f}$

Obs:  $K^\alpha = \frac{dp^\alpha}{dz}$

$$= \begin{cases} q \gamma \vec{E} \cdot \vec{u} = \frac{d(mu\gamma)}{dz} \Rightarrow q \vec{E} \cdot \vec{u} = \frac{d(mu\gamma)}{dz} = \frac{dI}{dt} \\ \gamma \vec{f} = \frac{d(\gamma \vec{u})}{dz} = \frac{dt}{dz} \frac{d(\gamma \vec{u})}{dt} = \gamma \frac{d(\gamma \vec{u})}{dt} \end{cases}$$

$\frac{d\vec{p}}{dt} = \vec{f}$

8 b) Como transforma  $\vec{f}$ ?  $K'^\mu = \Lambda^\mu_\nu K^\nu$

$\Lambda_{(ans)} = \begin{pmatrix} \gamma & -\gamma v^x & 0 & 0 \\ -\gamma v^x & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$

$$f'^i = \gamma^{-1} K'^i = \gamma^{-1} \Lambda^i_\mu K^\mu = \gamma^{-1} [\Lambda^i_0 K^0 + \Lambda^i_j K^j]$$

$$= \gamma^{-1} [-\gamma v^i q \gamma \vec{E} \cdot \vec{u} + \gamma \frac{K^i}{\gamma v^i}] = \gamma [v^i q \vec{E} \cdot \vec{u} + f^i]$$

$\Lambda$  boost con  $\vec{v} // \vec{f}$

\* o sea  $\vec{f}' = \gamma (\vec{f} + q \vec{v} (\vec{E} \cdot \vec{u}))$

8 c) Relación con  $\frac{d\vec{p}}{dt}$ :  $\frac{d\vec{p}}{dt} = \gamma^{-1} \frac{d\vec{p}}{dz} \rightarrow \frac{d\vec{p}}{dt} = \gamma \frac{d\vec{p}}{dz}$

Trans. con  $\vec{k} = \gamma \vec{f}$  igual

9) Ecuaciones de movimiento a partir de la acción

$$S = -m \int \sqrt{g_{\alpha\beta} dx^\alpha dx^\beta} - q \int A_\alpha dx^\alpha$$

$$= -m \int \sqrt{\sum_{\mu\nu} \frac{dx^\mu}{dt} \frac{dx^\nu}{dt}} dt - q \int A_\alpha \frac{dx^\alpha}{dt} dt$$

$$= -m \int \sqrt{1-u^2} dt - q \int (\phi - \vec{A} \cdot \vec{u}) dt \rightarrow m \text{ 3D}, q^i = x^i$$

$$E-L: \frac{d}{dt} \frac{\delta L}{\delta u_i} - \frac{\delta L}{\delta x_i} = 0 = \frac{d}{dt} \left[ \frac{-m}{2} \frac{(-2u^i)}{\sqrt{1-u^2}} + q A^i \right] - (-q \partial^i \phi + q \partial^i (\vec{A} \cdot \vec{u}))$$

$$0 = m \frac{d}{dt} (u^i) + q \frac{d}{dt} A^i + q \partial^i \phi - q \partial^i (\vec{A} \cdot \vec{u})$$

$$0 = \frac{d}{dt} (\vec{p}) + q (\partial_t \vec{A} + \underbrace{(\vec{u} \cdot \nabla) \vec{A}}_{\vec{B}}) + q \nabla \phi - q \nabla (\vec{A} \cdot \vec{u})$$

$$\nabla (\vec{u} \cdot \vec{A}) = (\vec{u} \cdot \nabla) \vec{A} + (\vec{A} \cdot \nabla) \vec{u} + \vec{u} \times (\nabla \times \vec{A}) + \vec{A} \times (\nabla \times \vec{u}) = (\vec{u} \cdot \nabla) \vec{A} + \vec{u} \times (\nabla \times \vec{A})$$

$$\nabla (\vec{u} \cdot \vec{A}) \Big|_{\partial u=0} = (\vec{u} \cdot \nabla) \vec{A} = \vec{u} \times (\nabla \times \vec{A})$$

$$0 = \frac{d\vec{p}}{dt} + q (\partial_t \vec{A} - \vec{u} \times \overbrace{(\nabla \times \vec{A})}^{\vec{B}}) + \nabla \phi = \frac{d\vec{p}}{dt} + q (-\vec{E} - \vec{u} \times \vec{B})$$

$$\boxed{\frac{d\vec{p}}{dt} = \vec{f}}$$

Alternativamente, de manera covariante:

$$S = -m \int \sqrt{\sum_{\mu\nu} \frac{dx^\mu}{dz} \frac{dx^\nu}{dz}} dz - q \int A_\alpha \frac{dx^\alpha}{dz} dz \quad ; \quad u \cdot u = 1$$

$$E-L: \frac{d}{dz} \frac{\delta L}{\delta u^\alpha} - \frac{\delta L}{\delta x^\alpha} = 0$$

$$\frac{d}{dz} \left[ \frac{-m}{2} \frac{2 U_\alpha}{\sqrt{u \cdot u}} - q A_\alpha \right] + q \partial_\alpha A_\beta U^\beta = 0$$

$$-m \frac{d}{dz} \left( \frac{U_\alpha}{\sqrt{u \cdot u}} \right) - q \frac{d}{dz} A_\alpha + q \partial_\alpha A_\beta U^\beta = 0$$

$$-m \left( \frac{dU_\alpha}{dz} \frac{1}{\sqrt{u \cdot u}} + \left(-\frac{1}{2}\right) U_\alpha \frac{2 U_\beta \frac{dU^\beta}{dz}}{(u \cdot u)^{3/2}} \right) - q U^\beta \partial_\beta A_\alpha + q \partial_\alpha A_\beta U^\beta = 0$$

$$-m \left( \frac{dU_\alpha}{\sqrt{u \cdot u}} - U_\alpha \frac{U_\beta dU^\beta}{(u \cdot u)^{3/2}} \right) + q \underbrace{(\partial_\alpha A_\beta - \partial_\beta A_\alpha)}_{F_{\alpha\beta}} U^\beta = 0$$

Y ahora,  $(z = \text{tiempo propio})$   
 $u \cdot u = 1$  :  $-m a_\alpha + q F_{\alpha\beta} U^\beta = 0 \Rightarrow m a^\alpha = q F^\alpha_{\beta} U^\beta =: K^\alpha$   
 $(a \cdot u = 0)$

$$u \cdot u = \left( \frac{dx^0}{dt'} \right)^2 - \left( \frac{dx^j}{dt'} \right)^2 = \left( \frac{dt}{dt'} \right)^2 - \left| \frac{dx^j}{dt} \right| \left( \frac{dt}{dt'} \right)^2 = \left( \frac{dt}{dt'} \right)^2 (1 - u^2) \stackrel{t'=z}{=} \gamma^2 (1 - u^2) = 1$$