### Física de Semiconductores

Lección 7

# Brillouin zone del grafeno





# Brillouin zone del grafeno I



# Brillouin zone del grafeno II



## Brillouin zone del grafeno



$$H = \begin{pmatrix} \frac{\hbar^{2}(\mathbf{k} - \mathbf{K})^{2}}{2m} - E \\ \frac{\hbar^{2}(\mathbf{k} - \mathbf{K})}{2m} - E \\ \frac{\hbar^{2}(\mathbf{k} - \mathbf{K}_{0})}{2m} - E \\ \frac{\hbar^{2}(\mathbf{k} - \mathbf{K}_{0})}{2m} & U_{\mathbf{K}_{1} - \mathbf{K}_{0}} & U_{\mathbf{K}_{2} - \mathbf{K}_{0}} & U_{\mathbf{K}_{3} - \mathbf{K}_{0}} & \cdots \\ U_{\mathbf{K}_{0} - \mathbf{K}_{1}} & \frac{\hbar^{2}(\mathbf{k} - \mathbf{K}_{1})}{2m} & U_{\mathbf{K}_{2} - \mathbf{K}_{1}} & U_{\mathbf{K}_{3} - \mathbf{K}_{1}} & \cdots \\ U_{\mathbf{K}_{0} - \mathbf{K}_{2}} & U_{\mathbf{K}_{1} - \mathbf{K}_{2}} & \frac{\hbar^{2}(\mathbf{k} - \mathbf{K}_{2})}{2m} & U_{\mathbf{K}_{3} - \mathbf{K}_{2}} & \cdots \\ U_{\mathbf{K}_{0} - \mathbf{K}_{3}} & U_{\mathbf{K}_{1} - \mathbf{K}_{3}} & U_{\mathbf{K}_{2} - \mathbf{K}_{3}} & \frac{\hbar^{2}(\mathbf{k} - \mathbf{K}_{3})}{2m} & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}$$

# Empty cell band structure

$$H = \begin{bmatrix} \frac{\hbar^2 \left( \mathbf{k} - \mathbf{K}_0 \right)}{2m} & 0 & 0 & 0 & \dots \\ 0 & \frac{\hbar^2 \left( \mathbf{k} - \mathbf{K}_1 \right)}{2m} & 0 & 0 & \dots \\ 0 & 0 & \frac{\hbar^2 \left( \mathbf{k} - \mathbf{K}_2 \right)}{2m} & 0 & \dots \\ 0 & 0 & 0 & \frac{\hbar^2 \left( \mathbf{k} - \mathbf{K}_3 \right)}{2m} & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots \end{bmatrix}$$

$$\begin{split} \mathbf{\Gamma} - \mathbf{M} \ & \mathbf{line} \qquad \mathbf{k} = \frac{2\pi}{a} \left( \frac{\delta}{\sqrt{3}}, 0, 0 \right) \quad \delta = 0 \to 1 \\ \mathbf{K}_{0} = 0; \quad E\left(\delta\right) = \frac{\hbar^{2}}{2m} \left( \frac{2\pi}{a} \right)^{2} \frac{\delta^{2}}{3} = E_{0} \frac{\delta^{2}}{3}; \quad E\left(0\right) = 0 \\ \mathbf{K}_{1} = \mathbf{b}_{1}; \quad E\left(\delta\right) = E_{0} \left[ \frac{\left(\delta - 1\right)^{2}}{3} + 1 \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{2} = \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[ \frac{\left(\delta - 1\right)^{2}}{3} + 1 \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{3} = -\mathbf{b}_{1}; \quad E\left(\delta\right) = E_{0} \left[ \frac{\left(\delta + 1\right)^{2}}{3} + 1 \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{4} = -\mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[ \frac{\left(\delta + 1\right)^{2}}{3} + 1 \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \end{split}$$

**F-M line (cont)** 
$$\mathbf{k} = \frac{2\pi}{a} \left( \frac{\delta}{\sqrt{3}}, 0, 0 \right) \quad \delta = 0 \to 1$$

$$\begin{split} \mathbf{K}_{5} &= \mathbf{b}_{1} + \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[\frac{\left(\delta - 2\right)^{2}}{3}\right]; \quad E\left(0\right) = \frac{4}{3}E_{0} \\ \mathbf{K}_{6} &= -\mathbf{b}_{1} - \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[\frac{\left(\delta + 2\right)^{2}}{3}\right]; \quad E\left(0\right) = \frac{4}{3}E_{0} \\ \mathbf{K}_{7} &= \mathbf{b}_{1} - \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[\frac{\delta^{2}}{3} + 4\right]; \quad E\left(0\right) = 4E_{0} \\ \mathbf{K}_{8} &= -\mathbf{b}_{1} + \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[\frac{\delta^{2}}{3} + 4\right]; \quad E\left(0\right) = 4E_{0} \end{split}$$

**Γ-M line (cont)**  $k = \frac{2\pi}{a} \left( \frac{\delta}{\sqrt{3}}, 0, 0 \right) \quad \delta = 0 \to 1$  $K_{9} = 2b_{1}; \quad E(\delta) = E_{0} \left| \frac{(\delta - 2)^{2}}{3} + 4 \right|; \quad E(0) = \frac{16}{3}E_{0}$  $K_{10} = 2b_2; \quad E(\delta) = E_0 \left| \frac{(\delta - 2)^2}{3} + 4 \right|; \quad E(0) = \frac{16}{3}E_0$  $m{K}_{11} = -2 m{b}_{1}; \quad E(\delta) = E_{0} \left| \frac{\left(\delta + 2\right)^{2}}{3} + 4 
ight|; \quad E(0) = \frac{16}{3} E_{0}$  $\mathbf{K}_{12} = -2\mathbf{b}_{2}; \quad E(\delta) = E_{0} \left| \frac{(\delta + 2)^{2}}{3} + 4 \right|; \quad E(0) = \frac{16}{3}E_{0}$ 

### Graphene empty cell Γ-M line



$$\begin{split} \mathbf{\Gamma} - \mathbf{K} \ \mathbf{line} \qquad & \mathbf{k} = \frac{2\pi}{a} \left( \frac{\lambda}{\sqrt{3}}, \frac{\lambda}{3}, 0 \right) \quad \delta = 0 \to 1 \\ \mathbf{K}_{0} = 0; \quad E\left(\lambda\right) = \frac{\hbar^{2}}{2m} \left( \frac{2\pi}{a} \right)^{2} \left[ \frac{\lambda^{2}}{3} + \frac{\lambda^{2}}{9} \right] = \frac{4}{9} E_{0} \lambda^{2}; \quad E\left(0\right) = 0 \\ \mathbf{K}_{1} = \mathbf{b}_{1}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda - 1\right)^{2}}{3} + \left( \frac{\lambda}{3} - 1 \right)^{2} \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{2} = \mathbf{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda - 1\right)^{2}}{3} + \left( \frac{\lambda}{3} + 1 \right)^{2} \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{3} = -\mathbf{b}_{1}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda + 1\right)^{2}}{3} + \left( \frac{\lambda}{3} + 1 \right)^{2} \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \\ \mathbf{K}_{4} = -\mathbf{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda + 1\right)^{2}}{3} + \left( \frac{\lambda}{3} - 1 \right)^{2} \right]; \quad E\left(0\right) = \frac{4}{3} E_{0} \end{split}$$

# **F-K line (cont)** $\mathbf{k} = \frac{2\pi}{a} \left( \frac{\lambda}{\sqrt{3}}, \frac{\lambda}{3}, 0 \right) \quad \delta = 0 \to 1$

$$\begin{split} \boldsymbol{K}_{5} &= \boldsymbol{b}_{1} + \boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[\frac{\left(\lambda - 2\right)^{2}}{3}\right]; \quad E\left(0\right) = \frac{4}{3}E_{0} \\ \boldsymbol{K}_{6} &= -\boldsymbol{b}_{1} - \boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[\frac{\left(\lambda + 2\right)^{2}}{3}\right]; \quad E\left(0\right) = \frac{4}{3}E_{0} \\ \boldsymbol{K}_{7} &= \boldsymbol{b}_{1} - \boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[\frac{\lambda}{3} - 2\right]^{2}; \quad E\left(0\right) = 4E_{0} \\ \boldsymbol{K}_{8} &= -\boldsymbol{b}_{1} + \boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[\frac{\lambda}{3} + 2\right]^{2}; \quad E\left(0\right) = 4E_{0} \end{split}$$

# **F-K line (cont)** $\mathbf{k} = \frac{2\pi}{a} \left( \frac{\lambda}{\sqrt{3}}, \frac{\lambda}{3}, 0 \right) \quad \delta = 0 \to 1$

$$\begin{split} \boldsymbol{K}_{9} &= 2\boldsymbol{b}_{1}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda - 2\right)^{2}}{3} + \left(\frac{\lambda}{3} - 2\right)^{2} \right]; \quad E\left(0\right) = \frac{16}{3} E_{0} \\ \boldsymbol{K}_{10} &= 2\boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda - 2\right)^{2}}{3} + \left(\frac{\lambda}{3} + 2\right)^{2} \right]; \quad E\left(0\right) = \frac{16}{3} E_{0} \\ \boldsymbol{K}_{11} &= -2\boldsymbol{b}_{1}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda + 2\right)^{2}}{3} + \left(\frac{\lambda}{3} + 2\right)^{2} \right]; \quad E\left(0\right) = \frac{16}{3} E_{0} \\ \boldsymbol{K}_{12} &= -2\boldsymbol{b}_{2}; \quad E\left(\lambda\right) = E_{0} \left[ \frac{\left(\lambda + 2\right)^{2}}{3} + \left(\frac{\lambda}{3} - 2\right)^{2} \right]; \quad E\left(0\right) = 4E_{0} \end{split}$$

# fcc unit cell



$oldsymbol{a}_1$	=	$\frac{a}{2}(\hat{\mathbf{j}} +$	$\hat{\mathbf{k}}$
$oldsymbol{a}_2$	=	$\frac{a}{2}(\hat{\mathbf{i}} +$	$\hat{\mathbf{k}}$
$oldsymbol{a}_{_3}$	—	$\frac{a}{2}(\hat{\mathbf{i}} +$	$\hat{\mathbf{j}}$



# Complex tight binding factors We define, to simplify the notation:

$$\begin{split} g_1 &\equiv \frac{1}{4} \bigg[ e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(4)}} \bigg] \\ g_2 &\equiv \frac{1}{4} \bigg[ e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(4)}} \bigg] \\ g_3 &\equiv \frac{1}{4} \bigg[ e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(4)}} \bigg] \\ g_4 &\equiv \frac{1}{4} \bigg[ e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}} - e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}} + e^{i \mathbf{k} \cdot \mathbf{n}_{AB}^{(4)}} \bigg] \end{split}$$

### H(sA,sB) matrix elements

$$\begin{split} H_{sA,sB}\left(\boldsymbol{k}\right) &= V_{ss\sigma}\sum_{m=1}^{N}e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{AB}^{(m)}}\\ H_{sB,sA}\left(\boldsymbol{k}\right) &= V_{ss\sigma}\sum_{m=1}^{N}e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{BA}^{(m)}} \end{split}$$

$$\begin{split} H_{sA,sB}\left(\boldsymbol{k}\right) &= V_{ss\sigma}\left(e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{AB}^{(1)}} + e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{AB}^{(2)}} + e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{AB}^{(3)}} + e^{i\boldsymbol{k}\cdot\boldsymbol{n}_{AB}^{(3)}}\right) = 4g_{1}V_{ss\sigma}\\ H_{sB,sA}\left(\boldsymbol{k}\right) &= H_{sB,sA}^{*}\left(\boldsymbol{k}\right) = 4g_{1}^{*}V_{ss\sigma} \end{split}$$

$$\begin{split} & \mathsf{H}(\mathsf{sA},\mathsf{pB}) \text{ matrix elements} \\ & H_{sA,p_{x}B}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{n} e^{i\mathbf{k}\cdot n_{AB}^{(m)}} \hat{n}_{AB,x}^{(m)} \\ &= \frac{V_{sp\sigma}}{\sqrt{3}} \left[ e^{i\mathbf{k}\cdot n_{AB}^{(1)}} + e^{i\mathbf{k}\cdot n_{AB}^{(2)}} + e^{i\mathbf{k}\cdot n_{AB}^{(3)}} \left(-1\right) + e^{i\mathbf{k}\cdot n_{AB}^{(4)}} \left(-1\right) \right] = \frac{4V_{sp\sigma}}{\sqrt{3}} g_2 \\ & H_{sA,p_{y}B}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{4} e^{i\mathbf{k}\cdot n_{AB}^{(m)}} n_{AB,y}^{(m)} = 4V_{sp\sigma} g_3 / \sqrt{3} \\ & H_{sA,p_{z}B}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{4} e^{i\mathbf{k}\cdot n_{AB}^{(m)}} n_{AB,z}^{(m)} = 4V_{sp\sigma} g_4 / \sqrt{3} \\ & H_{sB,p_{x}A}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{4} e^{i\mathbf{k}\cdot n_{BA}^{(m)}} n_{BA,x}^{(m)} = -4V_{sp\sigma} g_2 / \sqrt{3} \\ & H_{sB,p_{y}A}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{4} e^{i\mathbf{k}\cdot n_{BA}^{(m)}} n_{BA,y}^{(m)} = -4V_{sp\sigma} g_3 / \sqrt{3} \\ & H_{sB,p_{z}A}\left(\mathbf{k}\right) = V_{sp\sigma} \sum_{m=1}^{N} e^{i\mathbf{k}\cdot n_{BA}^{(m)}} n_{BA,z}^{(m)} = -V_{sp\sigma} g_4 / \sqrt{3} \end{split}$$

# Symmetrized elements in diamond

• Define

$$\begin{split} V_{ss} &= 4V_{ss\sigma} \\ V_{sp} &= 4V_{sp\sigma} / \sqrt{3} \\ V_{xx} &= \left(4V_{pp\sigma} / 3\right) + \left(8V_{pp\pi} / 3\right) \\ V_{xy} &= \left(4V_{pp\sigma} / 3\right) - \left(4V_{pp\pi} / 3\right) \end{split}$$

# Tight binding hamiltonian diamond

$$H = \begin{bmatrix} \varepsilon_s & 0 & 0 & 0 & V_{ss}g_1 & V_{sp}g_2 & V_{sp}g_3 & V_{sp}g_4 \\ 0 & \varepsilon_p & 0 & 0 & -V_{sp}g_2 & V_{xx}g_1 & V_{xy}g_4 & V_{xy}g_3 \\ 0 & 0 & \varepsilon_p & 0 & -V_{sp}g_3 & V_{xy}g_4 & V_{xx}g_1 & V_{xy}g_2 \\ 0 & 0 & 0 & \varepsilon_p & -V_{sp}g_4 & V_{xy}g_3 & V_{xy}g_2 & V_{xx}g_1 \\ V_{ss}g_1^* & -V_{sp}g_2^* & -V_{sp}g_3^* & -V_{sp}g_4^* & \varepsilon_s & 0 & 0 & 0 \\ V_{sp}g_2^* & V_{xx}g_1^* & V_{xy}g_4^* & V_{xy}g_3^* & 0 & \varepsilon_p & 0 & 0 \\ V_{sp}g_3^* & V_{xy}g_4^* & V_{xx}g_1^* & V_{xy}g_2^* & 0 & 0 & \varepsilon_p & 0 \\ V_{sp}g_4^* & V_{xy}g_3^* & V_{xy}g_2^* & V_{xx}g_1^* & 0 & 0 & 0 & \varepsilon_p \end{bmatrix}$$

#### **Diamond reciprocal lattice**

$$\begin{split} \mathbf{b}_{1} &= \frac{2\pi}{a} \left( \hat{\mathbf{i}} - \hat{\mathbf{j}} - \hat{\mathbf{k}} \right); \quad \mathbf{b}_{2} = \frac{2\pi}{a} \left( -\hat{\mathbf{i}} + \hat{\mathbf{j}} - \hat{\mathbf{k}} \right); \quad \mathbf{b}_{3} = \frac{2\pi}{a} \left( -\hat{\mathbf{i}} - \hat{\mathbf{j}} + \hat{\mathbf{k}} \right) \\ -\mathbf{b}_{1} &= \frac{2\pi}{a} \left( -\hat{\mathbf{i}} + \hat{\mathbf{j}} + \hat{\mathbf{k}} \right); \quad -\mathbf{b}_{2} = \frac{2\pi}{a} \left( \hat{\mathbf{i}} - \hat{\mathbf{j}} + \hat{\mathbf{k}} \right); \quad -\mathbf{b}_{3} = \frac{2\pi}{a} \left( \hat{\mathbf{i}} + \hat{\mathbf{j}} - \hat{\mathbf{k}} \right) \end{split}$$

$$\boldsymbol{b}_1 + \boldsymbol{b}_2 + \boldsymbol{b}_3 = \frac{2\pi}{a} \left( -\hat{\mathbf{i}} - \hat{\mathbf{j}} - \hat{\mathbf{k}} \right); \quad -\boldsymbol{b}_1 - \boldsymbol{b}_2 - \boldsymbol{b}_3 = \frac{2\pi}{a} \left( \hat{\mathbf{i}} + \hat{\mathbf{j}} + \hat{\mathbf{k}} \right)$$

But these 8 vectors do not define the BZ:

$$-\boldsymbol{b}_2 - \boldsymbol{b}_3 = \frac{4\pi}{a}\hat{\mathbf{i}}$$

#### Diamond BZ: a truncated octahedron





# Empty lattice band structure

• We consider the following reciprocal lattice vectors:

$$\begin{split} \mathbf{b}_{1} &= \frac{2\pi}{a} \big( \hat{\mathbf{i}} - \hat{\mathbf{j}} - \hat{\mathbf{k}} \big); \quad \mathbf{b}_{2} &= \frac{2\pi}{a} \big( -\hat{\mathbf{i}} + \hat{\mathbf{j}} - \hat{\mathbf{k}} \big); \quad \mathbf{b}_{3} &= \frac{2\pi}{a} \big( -\hat{\mathbf{i}} - \hat{\mathbf{j}} + \hat{\mathbf{k}} \big) \\ -\mathbf{b}_{1} &= \frac{2\pi}{a} \big( -\hat{\mathbf{i}} + \hat{\mathbf{j}} + \hat{\mathbf{k}} \big); \quad -\mathbf{b}_{2} &= \frac{2\pi}{a} \big( \hat{\mathbf{i}} - \hat{\mathbf{j}} + \hat{\mathbf{k}} \big); \quad -\mathbf{b}_{3} &= \frac{2\pi}{a} \big( \hat{\mathbf{i}} + \hat{\mathbf{j}} - \hat{\mathbf{k}} \big) \\ \mathbf{b}_{1} &+ \mathbf{b}_{2} + \mathbf{b}_{3} &= \frac{2\pi}{a} \big( -\hat{\mathbf{i}} - \hat{\mathbf{j}} - \hat{\mathbf{k}} \big); \quad -\mathbf{b}_{1} - \mathbf{b}_{2} - \mathbf{b}_{3} &= \frac{2\pi}{a} \big( \hat{\mathbf{i}} + \hat{\mathbf{j}} + \hat{\mathbf{k}} \big) \\ \mathbf{b}_{1} &+ \mathbf{b}_{2} &= -\frac{4\pi}{a} \hat{\mathbf{k}}; \quad -\mathbf{b}_{1} - \mathbf{b}_{2} = \frac{4\pi}{a} \hat{\mathbf{k}} \\ \mathbf{b}_{2} &+ \mathbf{b}_{3} &= -\frac{4\pi}{a} \hat{\mathbf{i}}; \quad -\mathbf{b}_{2} - \mathbf{b}_{3} = \frac{4\pi}{a} \hat{\mathbf{i}} \\ \mathbf{b}_{1} &+ \mathbf{b}_{3} &= -\frac{4\pi}{a} \hat{\mathbf{j}}; \quad -\mathbf{b}_{1} - \mathbf{b}_{3} = \frac{4\pi}{a} \hat{\mathbf{j}} \end{split}$$

$$\begin{split} & \mathbf{\Delta}\text{-line} \quad \mathbf{k} = \frac{2\pi}{2\pi} \left( \delta, 0, 0 \right) \\ & \mathbf{K}_{0} = 0; \quad E\left(\delta\right) = \frac{\hbar^{2}}{2m} \left( \frac{2\pi}{a} \right)^{2} \delta^{2} = E_{0} \delta^{2}; \quad E\left(0\right) = 0; \quad E\left(1\right) = E_{0} \\ & \mathbf{K}_{1} = \mathbf{b}_{1}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta - 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 2E_{0} \\ & \mathbf{K}_{2} = \mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta + 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 6E_{0} \\ & \mathbf{K}_{3} = \mathbf{b}_{3}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta + 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 6E_{0} \\ & \mathbf{K}_{4} = -\mathbf{b}_{1}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta + 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 6E_{0} \\ & \mathbf{K}_{5} = -\mathbf{b}_{2}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta - 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 2E_{0} \\ & \mathbf{K}_{6} = -\mathbf{b}_{3}; \quad E\left(\delta\right) = E_{0} \left[ \left(\delta - 1\right)^{2} + 2 \right]; \quad E\left(0\right) = 3E_{0}; \quad E\left(1\right) = 2E_{0} \end{split}$$

 $\Delta\text{-line (cont)} \quad k = \frac{2\pi}{a} \left( \delta, 0, 0 \right)$  $\boldsymbol{K}_{7} = \boldsymbol{b}_{1} + \boldsymbol{b}_{2} + \boldsymbol{b}_{3}; \quad E\left(\delta\right) = E_{0}\left|\left(\delta + 1\right)^{2} + 2\right|; E\left(0\right) = 3E_{0}; E\left(1\right) = 6E_{0}$  $\boldsymbol{K}_{8} = -\boldsymbol{b}_{1} - \boldsymbol{b}_{2} - \boldsymbol{b}_{3}; \quad E\left(\delta\right) = E_{0}\left|\left(\delta - 1\right)^{2} + 2\right|; E\left(0\right) = 3E_{0}; E\left(1\right) = 2E_{0}$  $\boldsymbol{K}_{9} = \boldsymbol{b}_{1} + \boldsymbol{b}_{2}; \quad E(\delta) = E_{0}(\delta^{2} + 4); \quad E(0) = 4E_{0}; \quad E(1) = 5E_{0}$  $K_{10} = -b_1 - b_2; \quad E(\delta) = E_0(\delta^2 + 4); \quad E(0) = 4E_0; \quad E(1) = 5E_0$  $\boldsymbol{K}_{11} = \boldsymbol{b}_{2} + \boldsymbol{b}_{3}; \quad E(\delta) = E_{0}(\delta + 2)^{2}; \quad E(0) = 4E_{0}; \quad E(1) = 9E_{0}$  $K_{12} = -b_2 - b_3; \quad E(\delta) = E_0(\delta - 2)^2; \quad E(0) = 4E_0; \quad E(1) = E_0$  $\boldsymbol{K}_{13} = \boldsymbol{b}_{1} + \boldsymbol{b}_{3}; \quad E\left(\delta\right) = E_{0}\left(\delta^{2} + 4\right); \quad E\left(0\right) = 4E_{0}; \quad E\left(1\right) = 15E_{0}$  $K_{14} = -b_1 - b_3; \quad E(\delta) = E_0(\delta^2 + 4); \quad E(0) = 4E_0; \quad E(1) = 15E_0$ 

$$\begin{array}{ll} \label{eq:constraint} \mathbf{\Lambda}\text{-line} \quad & \mathbf{k} = \frac{\pi}{a} \Big(\lambda, \lambda, \lambda\Big) = \frac{2\pi}{a} \bigg(\frac{\lambda}{2}, \frac{\lambda}{2}, \frac{\lambda}{2}\bigg) \\ & \mathbf{K}_{0} = 0; \quad E \Big(\lambda\Big) = \frac{\hbar^{2}}{2m} \bigg(\frac{2\pi}{a}\bigg)^{2} \bigg[ \bigg(\frac{\lambda}{2}\bigg)^{2} + \bigg(\frac{\lambda}{2}\bigg)^{2} + \bigg(\frac{\lambda}{2}\bigg)^{2} \bigg] = \frac{3}{4} E_{0} \lambda^{2}; \quad E \Big(0\Big) = 0 \\ & \mathbf{K}_{1} = \mathbf{b}_{1}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ \big(\lambda/2 - 1\big)^{2} + 2\big(\lambda/2 + 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \\ & \mathbf{K}_{2} = \mathbf{b}_{2}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ 2\big(\lambda/2 + 1\big)^{2} + \big(\lambda/2 - 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \\ & \mathbf{K}_{3} = \mathbf{b}_{3}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ 2\big(\lambda/2 + 1\big)^{2} + \big(\lambda/2 - 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \\ & \mathbf{K}_{4} = -\mathbf{b}_{1}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ \big(\lambda/2 + 1\big)^{2} + 2\big(\lambda/2 - 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \\ & \mathbf{K}_{5} = -\mathbf{b}_{2}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ \big(\lambda/2 + 1\big)^{2} + 2\big(\lambda/2 - 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \\ & \mathbf{K}_{6} = -\mathbf{b}_{3}; \quad E \Big(\lambda\Big) = E_{0} \bigg[ \big(\lambda/2 + 1\big)^{2} + 2\big(\lambda/2 - 1\big)^{2} \bigg]; \quad E \Big(0\Big) = 3E_{0} \end{aligned}$$

$$\begin{array}{l} \mathbf{\Lambda\text{-line (cont)}} \ \ \mathbf{k} = \frac{\pi}{\left[a}(\lambda,\lambda,\lambda)\right] = \frac{2\pi}{\left[a}\left(\frac{\lambda}{2},\frac{\lambda}{2},\frac{\lambda}{2}\right)\\ \mathbf{K}_{7} = \mathbf{b}_{1} + \mathbf{b}_{2} + \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[3(\lambda/2+1)^{2}\right]; \quad E(0) = 3E_{0}\\ \mathbf{K}_{8} = -\mathbf{b}_{1} - \mathbf{b}_{2} - \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[(\lambda/2-1)^{2}\right]; \quad E(0) = 3E_{0}\\ \mathbf{K}_{9} = \mathbf{b}_{1} + \mathbf{b}_{2}; \quad E(\lambda) = E_{0}\left[2(\lambda/2)^{2} + (\lambda/2+2)^{2}\right]; \quad E(0) = 4E_{0}\\ \mathbf{K}_{10} = -\mathbf{b}_{1} - \mathbf{b}_{2}; \quad E(\lambda) = E_{0}\left[2(\lambda/2)^{2} + (\lambda/2-2)^{2}\right]; \quad E(0) = 4E_{0}\\ \mathbf{K}_{11} = \mathbf{b}_{2} + \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[(\lambda/2+2)^{2} + 2(\lambda/2)^{2}\right]; \quad E(0) = 4E_{0}\\ \mathbf{K}_{12} = -\mathbf{b}_{2} - \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[(\lambda/2+2)^{2} + 2(\lambda/2)^{2}\right]; \quad E(0) = 4E_{0}\\ \mathbf{K}_{13} = \mathbf{b}_{1} + \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[(\lambda/2+2)^{2} + 2(\lambda/2)^{2}\right]; \quad E(0) = 4E_{0}\\ \mathbf{K}_{14} = -\mathbf{b}_{1} - \mathbf{b}_{3}; \quad E(\lambda) = E_{0}\left[(\lambda/2-2)^{2} + 2(\lambda/2)^{2}\right]; \quad E(0) = 4E_{0}\\ \end{array}$$

#### Solution comparison



## Harrison's match (PRB 20, 2420 (1979)

 Harrison suggested matching the TB band structure to the empty cell eigenvalues and use these to determine universal tight binding parameters.



#### TB Hamiltonian k = 0 H =

$arepsilon_s$	0	0	0	$V_{_{ss}}$	0	0	0
0	$arepsilon_p$	0	0	0	$V_{_{\! X\!X}}$	0	0
0	0	$arepsilon_p$	0	0	0	$V_{_{\! X\!X}}$	0
0	0	0	$arepsilon_p$	0	0	0	$V_{_{X\!X}}$
$V_{ss}$	0	0	0	$arepsilon_s$	0	0	0
0	$V_{_{\!X\!X}}$	0	0	0	${\mathcal E}_p^{}$	0	0
0	0	$V_{_{\!X\!X}}$	0	0	0	${\mathcal E}_p^{}$	0
0	0	0	$V_{_{\! X\!X}}$	0	0	0	$arepsilon_p$

#### s-states



$$\begin{split} E &= \varepsilon_s \pm \left| V_{ss} \right| \\ &2 \left| V_{ss} \right| = -8 V_{ss\sigma} = 3 E_0 \\ V_{ss\sigma} &= -\frac{3}{8} E_0 = -\frac{3}{8} \frac{\hbar^2}{2m} \left( \frac{2\pi}{a} \right)^2 = -\frac{3}{4} \frac{\hbar^2}{m} \frac{\pi^2}{a^2} = -\left( \frac{9\pi^2}{64} \right) \frac{\hbar^2}{md^2} \end{split}$$

# p-states $H = \begin{pmatrix} \varepsilon_p & V_{xx} \\ V_{xx} & \varepsilon_p \end{pmatrix}$

$$E = \varepsilon_p \pm \left| V_{xx} \right|$$
$$2 \left| V_{xx} \right| = 2 \left[ \left( 4V_{pp\sigma} / 3 \right) + \left( 8V_{pp\pi} / 3 \right) \right] = 4E_0 - 3E_0 = E_0$$

#### Tight-binding $\Delta$ direction



#### TB Hamiltonian $\Delta$ direction



#### **TB Hamiltonian X-point**



#### s-p states at X

$$H = \left( \begin{array}{cc} \varepsilon_{s} & iV_{sp} \\ -iV_{sp} & \varepsilon_{p} \end{array} \right)$$

$$E = \frac{\left(\varepsilon_{s} + \varepsilon_{p}\right) \pm \sqrt{\left(\varepsilon_{s} + \varepsilon_{p}\right)^{2} - 4\left(\varepsilon_{s}\varepsilon_{p} - V_{sp}^{2}\right)}}{2}$$

#### Pure p states at X

$$H = \left( egin{array}{cc} arepsilon_p & i V_{xy} \ -i V_{xy} & arepsilon_p \ -i V_{xy} & arepsilon_p \end{array} 
ight)$$

$$\begin{split} E &= \varepsilon_p \pm \left| V_{xy} \right| = \varepsilon_p \pm \left| \left( 4V_{pp\sigma} / 3 \right) - \left( 4V_{pp\pi} / 3 \right) \right| \\ p\left( 0 \right) - p\left( X \right) &= \left[ -\left( 4V_{pp\sigma} / 3 \right) - \left( 8V_{pp\pi} / 3 \right) \right] + \left[ \left( 4V_{pp\sigma} / 3 \right) - \left( 4V_{pp\pi} / 3 \right) \right] \\ &= -\frac{12}{3} V_{pp\pi} = -4V_{pp\pi} = 3E_0 - 2E_0 = E_0 = \frac{\hbar^2}{2m} \left( \frac{2\pi}{a} \right)^2 = \frac{\hbar^2}{md^2} \left( \frac{3}{8} \pi^2 \right) \\ \Rightarrow V_{pp\pi} = -\left( \frac{3}{32} \pi^2 \right) \frac{\hbar^2}{md^2} \end{split}$$

# **TB** universal parameters

$$V_{_{ll'm}}=\eta_{_{ll'm}}rac{\hbar^2}{md^2}$$

TB parameters	$\eta_{_{ss\sigma}}$	$\eta_{_{sp\sigma}}$	${\eta}_{_{pp\sigma}}$	${\eta}_{_{pp\pi}}$
Ideal	$9\pi^{2}/64$	$3\sqrt{15}\pi^2/64$	$21\pi^{2}/64$	$-3\pi^{2}/32$
Ideal	1.39	1.79	3.24	-0.93
Fit Ge	-1.40	1.84	3.24	-0.81
Universal	-1.32	1.42	2.22	-0.63

#### Lattice parameters



# Bonding in Si, Ge, Sn

