

Física de Semiconductores

Lección 14

Time reversal symmetry I

- Time dependent Schrödinger equation

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{nk}(\mathbf{r}, t) = i\hbar \frac{\partial \psi_{nk}(\mathbf{r}, t)}{\partial t}$$

- Reverse time

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{nk}(\mathbf{r}, -t) = -i\hbar \frac{\partial \psi_{nk}(\mathbf{r}, -t)}{\partial t}$$

- Take complex conjugate:

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{nk}^*(\mathbf{r}, -t) = i\hbar \frac{\partial \psi_{nk}^*(\mathbf{r}, -t)}{\partial t}$$

Time reversal symmetry II

- Choose $\psi_{n\mathbf{k}}^* (\mathbf{r}, -t) = \psi_{n\mathbf{k}}^* (\mathbf{r}) e^{-iE_n(\mathbf{k})/\hbar}$

- Then

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{n\mathbf{k}}^* (\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}^* (\mathbf{r})$$

- This means that

$$\psi_{n\mathbf{k}}^* (\mathbf{r}), \psi_{n\mathbf{k}} (\mathbf{r})$$

- have the same energy eigenvalue. Are they different Bloch functions?

Time reversal symmetry III

- For any Bloch function

$$\psi_{n\mathbf{k}}(\mathbf{r} + \mathbf{R}) = e^{i\mathbf{k}\cdot\mathbf{R}} \psi_{n\mathbf{k}}(\mathbf{r})$$

- Then

$$\psi_{n\mathbf{k}}^*(\mathbf{r} + \mathbf{R}) = e^{-i\mathbf{k}\cdot\mathbf{R}} \psi_{n\mathbf{k}}^*(\mathbf{r})$$

- This means that

$$\psi_{n\mathbf{k}}^*(\mathbf{r})$$

- is a Bloch function with wavevector $-\mathbf{k}$ and energy $E_n(\mathbf{k})$

$$\psi_{n,-\mathbf{k}}(\mathbf{r}) = \psi_{n\mathbf{k}}^*(\mathbf{r}) \Rightarrow E_n(-\mathbf{k}) = E_n(\mathbf{k})$$

Spatial inversion symmetry I

- If system has inversion symmetry $V(\mathbf{r}) = V(-\mathbf{r})$

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{n\mathbf{k}}(\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}(\mathbf{r})$$

- Invert \mathbf{r} :

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(-\mathbf{r}) \right] \psi_{n\mathbf{k}}(-\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}(-\mathbf{r})$$

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{n\mathbf{k}}(-\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}(-\mathbf{r})$$

Time-reversed Berry connection

- Let \mathcal{T} be time-reversal operator. Then

$$\begin{aligned}\mathcal{T}\mathcal{A}_n(\mathbf{k}) &= i \int d\mathbf{r} \left[\mathcal{T}u_{n\mathbf{k}}^*(\mathbf{r}) \right] \frac{\partial}{\partial \mathbf{k}} \left[\mathcal{T}u_{n\mathbf{k}}(\mathbf{r}) \right] \\ &= i \int d\mathbf{r} u_{n-\mathbf{k}}(\mathbf{r}) \frac{\partial}{\partial \mathbf{k}} u_{n-\mathbf{k}}^*(\mathbf{r}) = -i \int d\mathbf{r} u_{n-\mathbf{k}}^*(\mathbf{r}) \frac{\partial}{\partial \mathbf{k}} u_{n-\mathbf{k}}(\mathbf{r}) \\ &= i \int d\mathbf{r} u_{n-\mathbf{k}}^*(\mathbf{r}) \frac{\partial}{\partial(-\mathbf{k})} u_{n-\mathbf{k}}(\mathbf{r}) = \mathcal{A}_n(-\mathbf{k})\end{aligned}$$

- But if system has time-reversal symmetry

$$\mathcal{T}\mathcal{A}_n(\mathbf{k}) = \mathcal{A}_n(\mathbf{k}) + \nabla\chi(\mathbf{k})$$

- SO

$$\mathcal{A}_n(-\mathbf{k}) = \mathcal{A}_n(\mathbf{k}) + \nabla\chi(\mathbf{k})$$

Time-reversed Berry curvature

$$\Omega(\mathbf{k}) = i \int d\mathbf{r} \frac{\partial u_{n\mathbf{k}}^*(\mathbf{r})}{\partial \mathbf{k}} \times \frac{\partial u_{n\mathbf{k}}(\mathbf{r})}{\partial \mathbf{k}}$$

$$\mathcal{T}\Omega(\mathbf{k}) = i \int d\mathbf{r} \frac{\partial \mathcal{T}u_{n\mathbf{k}}^*(\mathbf{r})}{\partial \mathbf{k}} \times \frac{\partial \mathcal{T}u_{n\mathbf{k}}(\mathbf{r})}{\partial \mathbf{k}}$$

$$= i \int d\mathbf{r} \frac{\partial u_{n-\mathbf{k}}(\mathbf{r})}{\partial \mathbf{k}} \times \frac{\partial u_{n-\mathbf{k}}^*(\mathbf{r})}{\partial \mathbf{k}} = -i \int d\mathbf{r} \frac{\partial u_{n-\mathbf{k}}^*(\mathbf{r})}{\partial \mathbf{k}} \times \frac{\partial u_{n-\mathbf{k}}(\mathbf{r})}{\partial \mathbf{k}}$$

$$= -i \int d\mathbf{r} \frac{\partial u_{n-\mathbf{k}}^*(\mathbf{r})}{\partial(-\mathbf{k})} \times \frac{\partial u_{n-\mathbf{k}}(\mathbf{r})}{\partial(-\mathbf{k})} = -\Omega(-\mathbf{k})$$

- So if there is time reversal symmetry:

$$\Omega(\mathbf{k}) = -\Omega(-\mathbf{k})$$

Point group symmetry

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(S\mathbf{r}) \right] \psi_{n\mathbf{k}}(S\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}(S\mathbf{r})$$

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \right] \psi_{n\mathbf{k}}(S\mathbf{r}) = E_n(\mathbf{k}) \psi_{n\mathbf{k}}(S\mathbf{r})$$

- Therefore $\psi_{n\mathbf{k}}(S\mathbf{r})$ is an eigenfunction with eigenvalue $E_n(\mathbf{k})$ Now

$$\psi_{n\mathbf{k}}(S(\mathbf{r} + \mathbf{R})) = \psi_{n\mathbf{k}}(S\mathbf{r} + S\mathbf{R})$$

- But since $\psi_{n\mathbf{k}}(S\mathbf{r})$ is a Bloch function and $S\mathbf{R}$ is a lattice vector,

Point symmetry (II)

$$\begin{aligned}\psi_{n\mathbf{k}}(S\mathbf{r} + S\mathbf{R}) &= e^{i\mathbf{k}\cdot S\mathbf{R}} \psi_{n\mathbf{k}}(S\mathbf{r}) = \\ &= e^{iS^{-1}S\mathbf{k}\cdot S\mathbf{R}} \psi_{n\mathbf{k}}(S\mathbf{r}) = e^{iS^{-1}\mathbf{k}\cdot \mathbf{R}} \psi_{n\mathbf{k}}(S\mathbf{r})\end{aligned}$$

- So that $\psi_{n\mathbf{k}}(S\mathbf{r})$ is a Bloch function with energy $E_n(\mathbf{k})$ and wave vector $S^{-1}\mathbf{k}$

$$\psi_{n\mathbf{k}}(S\mathbf{r}) = \psi_{nS^{-1}\mathbf{k}}(\mathbf{r})$$

- So

$$E_n(\mathbf{k}) = E_n(S^{-1}\mathbf{k})$$

Inversion Berry connection

- If the point symmetry is the inversion \mathcal{J}

$$\mathcal{J}u_{n\mathbf{k}}(\mathbf{r}) = u_{n-\mathbf{k}}(\mathbf{r})$$

- One can then show that

$$\mathcal{J}\mathcal{A}(\mathbf{k}) = \mathcal{A}(-\mathbf{k})$$

$$\mathcal{J}\mathcal{O}(\mathbf{k}) = \mathcal{O}(-\mathbf{k})$$

- If the system is invariant under inversion

$$\mathcal{J}\mathcal{A}(\mathbf{k}) = \mathcal{A}(\mathbf{k}) + \nabla\chi(\mathbf{k}) = \mathcal{A}(-\mathbf{k})$$

$$\mathcal{J}\mathcal{O}(\mathbf{k}) = \mathcal{O}(\mathbf{k}) = \mathcal{O}(-\mathbf{k})$$

Chern number I

- If q_m is not zero

$$\begin{aligned} 4\pi q_m &= \oiint (\nabla \times \mathbf{A}) \cdot d\mathbf{S} = \int_{\partial D_1} (\mathbf{A}_1 - \mathbf{A}_2) \cdot d\mathbf{l} \\ &= \int_{\partial D_1} \nabla (2q_m \phi) \cdot d\mathbf{l} = 2q_m \int_{\partial D_1} \nabla (\phi) \cdot d\mathbf{l} = 2q_m (2\pi) \end{aligned}$$

- The integral

$$C = \frac{1}{2\pi q_m} \oiint (\nabla \times \mathbf{A}) \cdot d\mathbf{S}$$

- is then an integer.

Chern number II

- By analogy we define the Chern number as

$$C = \frac{1}{2\pi} \oint \Omega(\mathbf{k}) \cdot d\mathbf{S}$$

- If we identify the “surface” with the BZ, and consider that the integral of an odd function of \mathbf{k} over the BZ is zero, then we conclude that the Chern number is zero if there is time inversion symmetry.

Chern number III

- If there is both time reversal and inversion symmetry

$$\Omega(\mathbf{k}) = -\Omega(-\mathbf{k}) \quad \text{TR}$$

$$\Omega(\mathbf{k}) = \Omega(-\mathbf{k}) \quad \text{Inversion}$$

$$\Omega(\mathbf{k}) = 0$$

Effect on magnetic field

- If we move an electron in a closed loop in the presence of a magnetic field, the electron picks up a phase

$$\phi = \frac{e}{c\hbar} \oint \mathbf{A} \cdot d\mathbf{l} = \frac{e}{c\hbar} \iint \mathbf{B} \cdot d\mathbf{S}$$

- Under time reversal the phase changes sign, so

$$\mathcal{T} \iint \mathbf{B} \cdot d\mathbf{S} = - \iint \mathbf{B} \cdot d\mathbf{S}$$

$$\Rightarrow \mathcal{T}\mathbf{B} = -\mathbf{B}$$

Phase with time reversal symmetry

- If a system has time-reversal symmetry, we tentatively say

$$e^{i\phi} = e^{-i\phi} \Rightarrow \phi = -\phi \Rightarrow \phi = 0 \Rightarrow B = 0$$

- But true condition is

$$\phi = -\phi + 2m\pi$$

- So that to preserve time-reversal symmetry we only need $\phi = m\pi$, so we can have

$$B = \frac{c\hbar}{e} \sum_i m_i \pi \delta(r - r_i)$$

Application to Berry curvature

- Under time reversal symmetry

$$\Omega(\mathbf{k}) = \sum_i m_i \pi \delta(k - k_i)$$

Dirac point

- Consider a hamiltonian

$$H = v_F \begin{pmatrix} 0 & k_x - ik_y \\ k_x + ik_y & 0 \end{pmatrix} = v_F k \begin{pmatrix} 0 & e^{-i\phi} \\ e^{i\phi} & 0 \end{pmatrix}$$

- Eigenvalues: $\pm v_F k$

- Eigenvectors: $u_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\phi} \\ 1 \end{pmatrix}; \quad u_- = \frac{1}{\sqrt{2}} \begin{pmatrix} -e^{-i\phi} \\ 1 \end{pmatrix}$

Berry connection at Dirac point

- The Berry connection in polar coordinates has components k and ϕ :

$$\mathcal{A}_k^+ = i \left\langle u_+ \left| \frac{\partial}{\partial k} \right| u_+ \right\rangle = 0$$

$$\mathcal{A}_\phi^+ = \frac{i}{k} \left\langle u_+ \left| \frac{\partial}{\partial \phi} \right| u_+ \right\rangle = 0$$

$$= \frac{i}{2k} \begin{pmatrix} e^{i\phi} & 1 \end{pmatrix} \frac{\partial}{\partial \phi} \begin{pmatrix} e^{-i\phi} \\ 1 \end{pmatrix} = \frac{i}{2k} \begin{pmatrix} e^{i\phi} & 1 \end{pmatrix} \begin{pmatrix} -ie^{-i\phi} \\ 0 \end{pmatrix} = \frac{1}{2k}$$

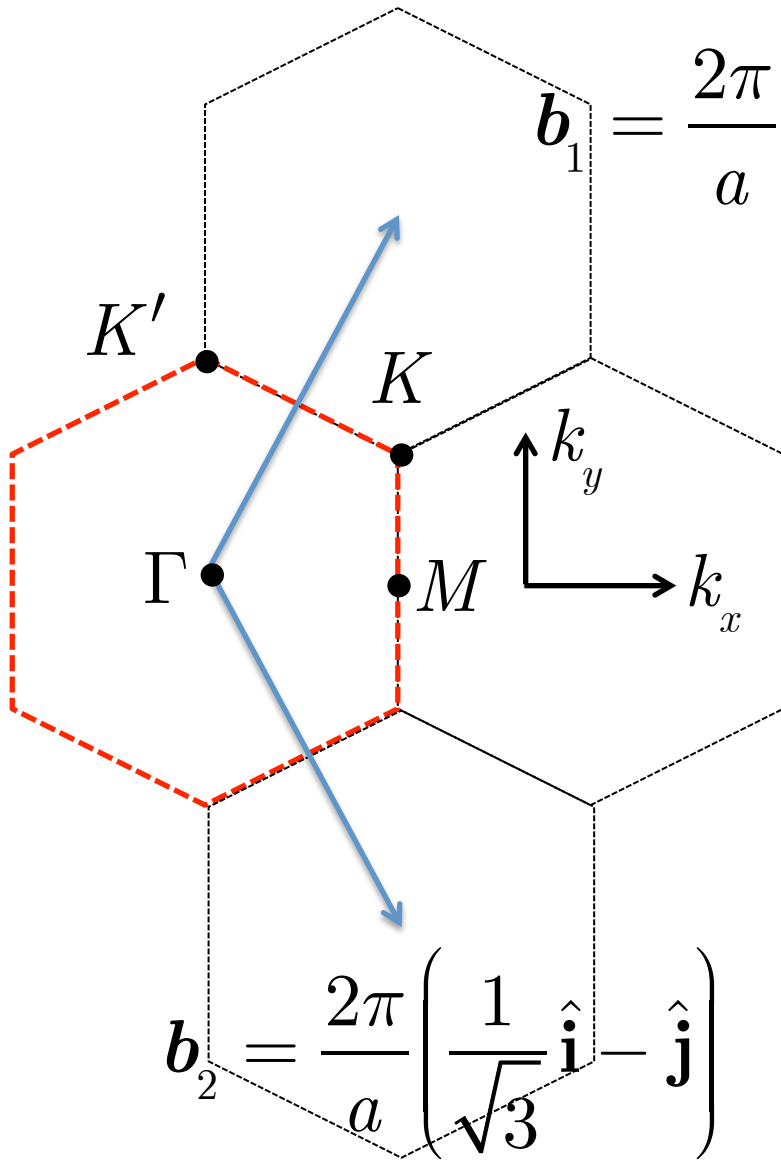
Berry flux

$$\iint \boldsymbol{\Omega}(\mathbf{k}) \cdot d\mathbf{S} = \oint \mathcal{A} \cdot d\mathbf{l}$$

- Taking a circle around $\mathbf{k}=0$

$$\oint \mathcal{A} \cdot d\mathbf{l} = \int_0^{2\pi} \frac{1}{2k} k d\phi = \pi$$

Brillouin zone del grafeno



$$\mathbf{b}_1 = \frac{2\pi}{a} \left(\frac{1}{\sqrt{3}} \hat{\mathbf{i}} + \hat{\mathbf{j}} \right)$$

$$M = \frac{1}{2} (\mathbf{b}_1 + \mathbf{b}_2) = \frac{2\pi}{a\sqrt{3}} \hat{\mathbf{i}}$$

$$K = \frac{2}{3} \mathbf{b}_1 + \frac{1}{3} \mathbf{b}_2 = \frac{2\pi}{a} \left(\frac{1}{\sqrt{3}} \hat{\mathbf{i}} + \frac{1}{3} \hat{\mathbf{j}} \right)$$

$$K' = \frac{1}{3} \mathbf{b}_1 - \frac{1}{3} \mathbf{b}_2 = \frac{2\pi}{a} \left(\frac{2}{3} \hat{\mathbf{j}} \right)$$

Graphene

- The hamiltonian matrix for graphene is

$$H = \begin{pmatrix} H_{AA} & H_{AB} \\ H_{BA} & H_{BB} \end{pmatrix}$$

$$H_{AA} = H_{BB} = \begin{pmatrix} \varepsilon_s & 0 & 0 & 0 \\ 0 & \varepsilon_{p\sigma} & 0 & 0 \\ 0 & 0 & \varepsilon_{p\sigma} & 0 \\ 0 & 0 & 0 & \varepsilon_{p\pi} \end{pmatrix}$$

$$H_{AB} = H_{BA}^*$$

$$\begin{pmatrix} V_{ss\sigma} (g_1 + g_2 + g_3) & V_{sp\sigma} \left[g_1 - \frac{1}{2}(g_2 + g_3) \right] & \frac{\sqrt{3}}{2} V_{sp\sigma} (g_2 - g_3) & 0 \\ -V_{sp\sigma} \left[g_1 - \frac{1}{2}(g_2 + g_3) \right] & g_1 V_{pp\sigma} + (g_2 + g_3) \left[\frac{3}{4} V_{pp\pi} + \frac{1}{4} V_{pp\sigma} \right] & \frac{\sqrt{3}}{4} (V_{pp\sigma} - V_{pp\pi}) (g_3 - g_2) & 0 \\ -\frac{\sqrt{3}}{2} V_{sp\sigma} (g_2 - g_3) & \frac{\sqrt{3}}{4} (V_{pp\sigma} - V_{pp\pi}) (g_3 - g_2) & g_1 V_{pp\pi} + (g_2 + g_3) \left[\frac{1}{4} V_{pp\pi} + \frac{3}{4} V_{pp\sigma} \right] & 0 \\ 0 & 0 & 0 & (g_1 + g_2 + g_3) V_p \end{pmatrix}$$

The $p\pi$ submatrix

$$H = \begin{pmatrix} \varepsilon_{p\pi} & (g_1 + g_2 + g_3) V_{pp\pi} \\ (g_1 + g_2 + g_3)^* V_{pp\pi} & \varepsilon_{p\pi} \end{pmatrix}$$

$$g_1 \equiv e^{i\mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}}$$

$$g_2 \equiv e^{i\mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}}$$

$$g_3 \equiv e^{i\mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}}$$

$$\mathbf{n}_{AB}^{(1)} = -\mathbf{n}_{BA}^{(1)} = (a/\sqrt{3})\hat{\mathbf{i}}$$

$$\mathbf{n}_{AB}^{(2)} = -\mathbf{n}_{BA}^{(2)} = (a/\sqrt{3})\left[-(1/2)\hat{\mathbf{i}} + (\sqrt{3}/2)\hat{\mathbf{j}}\right]$$

$$\mathbf{n}_{AB}^{(3)} = -\mathbf{n}_{BA}^{(3)} = (a/\sqrt{3})\left[-(1/2)\hat{\mathbf{i}} - (\sqrt{3}/2)\hat{\mathbf{j}}\right]$$

Pauli matrix formulation

$$H = \begin{pmatrix} \varepsilon_{p\pi} & (g_1 + g_2 + g_3) V_{pp\pi} \\ (g_1 + g_2 + g_3)^* V_{pp\pi} & \varepsilon_{p\pi} \end{pmatrix}$$

$$= \begin{pmatrix} H_{11}(\mathbf{k}) & H_{12}(\mathbf{k}) \\ H_{21}(\mathbf{k}) & H_{22}(\mathbf{k}) \end{pmatrix}$$

We can write this as

$$g_1 + g_2 + g_3 =$$

$$= H_0(\mathbf{k}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + H_x(\mathbf{k}) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + H_y(\mathbf{k}) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + H_z(\mathbf{k}) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

The $p\pi$ submatrix II

Then

$$H_{11}(\mathbf{k}) = H_0(\mathbf{k}) + H_z(\mathbf{k}); \quad H_{22}(\mathbf{k}) = H_0(\mathbf{k}) - H_z(\mathbf{k})$$

$$H_{12}(\mathbf{k}) = H_x(\mathbf{k}) - iH_y(\mathbf{k}); \quad H_{21}(\mathbf{k}) = H_x(\mathbf{k}) + iH_y(\mathbf{k})$$

$$H_0(\mathbf{k}) = \frac{H_{11}(\mathbf{k}) + H_{22}(\mathbf{k})}{2} = \frac{\varepsilon_{pA} + \varepsilon_{pB}}{2} = \varepsilon_p$$

$$H_z(\mathbf{k}) = \frac{H_{11}(\mathbf{k}) - H_{22}(\mathbf{k})}{2} = \frac{\varepsilon_{pA} - \varepsilon_{pB}}{2} = \Delta_p$$

$$H_x(\mathbf{k}) = \frac{H_{12}(\mathbf{k}) + H_{21}(\mathbf{k})}{2}$$

$$= \text{Re}[H_{12}(\mathbf{k})] = V_{pp\pi} \left[\cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}) + \cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}) + \cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}) \right]$$

$$H_y(\mathbf{k}) = i \frac{H_{12}(\mathbf{k}) - H_{21}(\mathbf{k})}{2} = -\text{Im}[H_{12}(\mathbf{k})]$$

$$= -V_{pp\pi} \left[\sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}) + \sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}) + \sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}) \right]$$

The $\rho\pi$ submatrix III

Expanding

$$\begin{aligned} H_x(\mathbf{k}) &= V_{pp\pi} \left[\cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}) + \cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}) + \cos(\mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}) \right] = \\ &= V_{pp\pi} \left[\cos(k_x a / \sqrt{3}) + \cos\left(-\frac{1}{2} k_x a / \sqrt{3} + \frac{1}{2} k_y a\right) + \cos\left(-\frac{1}{2} k_x a / \sqrt{3} - \frac{1}{2} k_y a\right) \right] \\ &= V_{pp\pi} \left[\cos(k_x a / \sqrt{3}) + 2 \cos\left(\frac{1}{2} k_x a / \sqrt{3}\right) \cos\left[\frac{1}{2} k_y a\right] \right] \\ H_y(\mathbf{k}) &= -V_{pp\pi} \left[\sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(1)}) + \sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(2)}) + \sin(\mathbf{k} \cdot \mathbf{n}_{AB}^{(3)}) \right] \\ &= -V_{pp\pi} \left[\sin(k_x a / \sqrt{3}) + \sin\left(-\frac{1}{2} k_x a / \sqrt{3} + \frac{1}{2} k_y a\right) + \sin\left(-\frac{1}{2} k_x a / \sqrt{3} - \frac{1}{2} k_y a\right) \right] \\ H_y(\mathbf{k}) &= -V_{pp\pi} \left[\sin(k_x a / \sqrt{3}) - 2 \sin\left(\frac{1}{2} k_x a / \sqrt{3}\right) \cos\left[\frac{1}{2} k_y a\right] \right] \end{aligned}$$

Near K (I)

- We propose $K = k_x = \frac{2\pi}{a\sqrt{3}} + \delta k_x; k_y = \frac{2\pi}{3a} + \delta k_y$

$$\begin{aligned}
 H_x(\mathbf{k}) &= V_{pp\pi} \left[\cos(k_x a/\sqrt{3}) + 2 \cos\left(\frac{1}{2} k_x a/\sqrt{3}\right) \cos\left[\frac{1}{2} k_y a\right] \right] \\
 &= V_{pp\pi} \left[\cos\left(\left\{\frac{2\pi}{a\sqrt{3}} + \delta k_x\right\} a/\sqrt{3}\right) + 2 \cos\left(\frac{1}{2} \left\{\frac{2\pi}{a\sqrt{3}} + \delta k_x\right\} a/\sqrt{3}\right) \cos\left[\frac{1}{2} \left\{\frac{2\pi}{3a} + \delta k_y\right\} a\right] \right] \\
 &= V_{pp\pi} \left[\cos\left(\frac{2\pi}{3} + \delta k_x a/\sqrt{3}\right) + 2 \cos\left(\frac{\pi}{3} + \frac{1}{2} \delta k_x a/\sqrt{3}\right) \cos\left(\frac{\pi}{3} + \frac{\delta k_y a}{2}\right) \right] \\
 &= V_{pp\pi} \left[\cos\left(\frac{2\pi}{3}\right) \cos\left(\delta k_x a/\sqrt{3}\right) - \sin\left(\frac{2\pi}{3}\right) \sin\left(\delta k_x a/\sqrt{3}\right) + \right. \\
 &\quad \left. + 2 \left\{ \cos\left(\frac{\pi}{3}\right) \cos\left(\frac{1}{2} \delta k_x a/\sqrt{3}\right) - \sin\left(\frac{\pi}{3}\right) \sin\left(\frac{1}{2} \delta k_x a/\sqrt{3}\right) \right\} \cos\left(\frac{\pi}{3} + \frac{\delta k_y a}{2}\right) \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2}\right) \cos\left(\delta k_x a/\sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\delta k_x a/\sqrt{3}\right) + \right. \\
 &\quad \left. 2 \left\{ \frac{1}{2} \cos\left(\frac{1}{2} \delta k_x a/\sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{1}{2} \delta k_x a/\sqrt{3}\right) \right\} \left[\cos\left(\frac{\pi}{3}\right) \cos\left(\frac{\delta k_y a}{2}\right) - \sin\left(\frac{\pi}{3}\right) \sin\left(\frac{\delta k_y a}{2}\right) \right] \right]
 \end{aligned}$$

Near K (II)

- We propose $K = k_x = \frac{2\pi}{a\sqrt{3}} + \delta k_x; k_y = \frac{2\pi}{3a} + \delta k_y$

$$\begin{aligned}
 &= V_{pp\pi} \left[\cos\left(\frac{2\pi}{3}\right) \cos(\delta k_x a / \sqrt{3}) - \sin\left(\frac{2\pi}{3}\right) \sin(\delta k_x a / \sqrt{3}) + 2 \left\{ \cos\left(\frac{\pi}{3}\right) \cos\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) - \sin\left(\frac{\pi}{3}\right) \sin\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) \right\} \cos\left(\frac{\pi}{3} + \frac{\delta k_y a}{2}\right) \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) \cos(\delta k_x a / \sqrt{3}) - \frac{\sqrt{3}}{2} \sin(\delta k_x a / \sqrt{3}) + 2 \left\{ \frac{1}{2} \cos\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) \right\} \left\{ \cos\left(\frac{\pi}{3}\right) \cos\left(\frac{\delta k_y a}{2}\right) - \sin\left(\frac{\pi}{3}\right) \sin\left(\frac{\delta k_y a}{2}\right) \right\} \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) \cos(\delta k_x a / \sqrt{3}) - \frac{\sqrt{3}}{2} \sin(\delta k_x a / \sqrt{3}) + 2 \left\{ \frac{1}{2} \cos\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) \right\} \left\{ \frac{1}{2} \cos\left(\frac{\delta k_y a}{2}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{\delta k_y a}{2}\right) \right\} \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) - \frac{\sqrt{3}}{2} (\delta k_x a / \sqrt{3}) + \left\{ 1 - \sqrt{3} \left(\frac{1}{2} \delta k_x a / \sqrt{3} \right) \right\} \left\{ \frac{1}{2} - \frac{\sqrt{3}}{4} \delta k_y a \right\} \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) - \frac{\delta k_x a}{2} + \left\{ \frac{1}{2} - \frac{\sqrt{3}}{4} \delta k_y a - \frac{1}{2} \left(\frac{1}{2} \delta k_x a \right) \right\} \right] \\
 &= V_{pp\pi} \left[-\frac{\delta k_x a}{2} + \left\{ -\frac{\sqrt{3}}{4} \delta k_y a - \frac{\delta k_x a}{4} \right\} \right] = V_{pp\pi} \left[-\frac{3\delta k_x a}{4} - \frac{\sqrt{3}}{4} \delta k_y a \right] = -\frac{V_{pp\pi} a}{4} [3\delta k_x + \sqrt{3}\delta k_y] = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x + \frac{\delta k_y}{\sqrt{3}} \right]
 \end{aligned}$$

Axis rotation K point

- We have
$$H_x(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x + \frac{\delta k_y}{\sqrt{3}} \right]$$

$$H_y(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[-\frac{\delta k_x}{\sqrt{3}} + \delta k_y \right]$$

- Define
$$\begin{pmatrix} q_x \\ q_y \end{pmatrix} = \begin{pmatrix} \frac{\sqrt{3}}{2} & \frac{1}{2} \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} \end{pmatrix} \begin{pmatrix} \delta k_x \\ \delta k_y \end{pmatrix}$$

$$H_x(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x + \frac{\delta k_y}{\sqrt{3}} \right] = -\frac{3V_{pp\pi} a}{4} \left(\frac{2q_x}{\sqrt{3}} \right) = -\frac{\sqrt{3}V_{pp\pi} a}{2} q_x$$

$$H_y(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[-\frac{\delta k_x}{\sqrt{3}} + \delta k_y \right] = -\frac{3V_{pp\pi} a}{4} \left(\frac{2q_y}{\sqrt{3}} \right) = -\frac{\sqrt{3}V_{pp\pi} a}{2} q_y$$

K' (II)

• We propose $K = k_x = \frac{2\pi}{a\sqrt{3}} + \delta k_x; k_y = -\frac{2\pi}{3a} + \delta k_y$

$$\begin{aligned}
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) \cos(\delta k_x a / \sqrt{3}) - \frac{\sqrt{3}}{2} \sin(\delta k_x a / \sqrt{3}) \right. \\
 &+ 2 \left. \left\{ \frac{1}{2} \cos\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) \right\} \left[\cos\left(\frac{\pi}{3}\right) \cos\left(\frac{\delta k_y a}{2}\right) + \sin\left(\frac{\pi}{3}\right) \sin\left(\frac{\delta k_y a}{2}\right) \right] \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) \cos(\delta k_x a / \sqrt{3}) - \frac{\sqrt{3}}{2} \sin(\delta k_x a / \sqrt{3}) \right. \\
 &+ 2 \left. \left\{ \frac{1}{2} \cos\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) - \frac{\sqrt{3}}{2} \sin\left(\frac{1}{2} \delta k_x a / \sqrt{3}\right) \right\} \left[\frac{1}{2} \cos\left(\frac{\delta k_y a}{2}\right) + \frac{\sqrt{3}}{2} \sin\left(\frac{\delta k_y a}{2}\right) \right] \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) - \frac{\sqrt{3}}{2} (\delta k_x a / \sqrt{3}) + \left\{ 1 - \sqrt{3} \left(\frac{1}{2} \delta k_x a / \sqrt{3} \right) \right\} \left[\frac{1}{2} + \frac{\sqrt{3}}{4} \delta k_y a \right] \right] \\
 &= V_{pp\pi} \left[\left(-\frac{1}{2} \right) - \frac{\delta k_x a}{2} + \left\{ \frac{1}{2} + \frac{\sqrt{3}}{4} \delta k_y a - \frac{1}{2} \left(\frac{1}{2} \delta k_x a \right) \right\} \right] \\
 &= V_{pp\pi} \left[-\frac{\delta k_x a}{2} + \left\{ +\frac{\sqrt{3}}{4} \delta k_y a - \frac{\delta k_x a}{4} \right\} \right] = V_{pp\pi} \left[-\frac{3\delta k_x a}{4} + \frac{\sqrt{3}}{4} \delta k_y a \right] = -\frac{V_{pp\pi} a}{4} [3\delta k_x - \sqrt{3}\delta k_y] = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x - \frac{\delta k_y}{\sqrt{3}} \right]
 \end{aligned}$$

Point K' (I)

- We propose $K = k_x = \frac{2\pi}{a\sqrt{3}} + \delta k_x; k_y = -\frac{2\pi}{3a} + \delta k_y$

$$\begin{aligned}
 H_y(\mathbf{k}) &= -V_{pp\pi} \left[\sin\left(k_x a/\sqrt{3}\right) - 2\sin\left(\frac{1}{2}k_x a/\sqrt{3}\right)\cos\left[\frac{1}{2}k_y a\right] \right] \\
 &= -V_{pp\pi} \left[\sin\left[\left\{\frac{2\pi}{a\sqrt{3}} + \delta k_x\right\}a/\sqrt{3}\right] - 2\sin\left[\frac{1}{2}\left\{\frac{2\pi}{a\sqrt{3}} + \delta k_x\right\}a/\sqrt{3}\right]\cos\left[\frac{1}{2}\left\{-\frac{2\pi}{3a} + \delta k_y\right\}a\right] \right] \\
 &= -V_{pp\pi} \left[\sin\left(\frac{2\pi}{3} + \frac{\delta k_x a}{\sqrt{3}}\right) - 2\sin\left(\frac{\pi}{3} + \frac{\delta k_x a}{2\sqrt{3}}\right)\cos\left(-\frac{\pi}{3} + \frac{\delta k_y a}{2}\right) \right] \\
 &= -V_{pp\pi} \left[\sin\left(\frac{2\pi}{3}\right)\cos\left(\frac{\delta k_x a}{\sqrt{3}}\right) + \cos\left(\frac{2\pi}{3}\right)\sin\left(\frac{\delta k_x a}{\sqrt{3}}\right) \right. \\
 &\quad \left. - 2\left[\sin\left(\frac{\pi}{3}\right)\cos\left(\frac{\delta k_x a}{2\sqrt{3}}\right) + \cos\left(\frac{\pi}{3}\right)\sin\left(\frac{\delta k_x a}{2\sqrt{3}}\right)\right]\left[\cos\left(-\frac{\pi}{3}\right)\cos\left(\frac{\delta k_y a}{2}\right) - \sin\left(-\frac{\pi}{3}\right)\sin\left(\frac{\delta k_y a}{2}\right)\right] \right] \\
 &= -V_{pp\pi} \left[\frac{\sqrt{3}}{2}\cos\left(\frac{\delta k_x a}{\sqrt{3}}\right) + \left(-\frac{1}{2}\right)\sin\left(\frac{\delta k_x a}{\sqrt{3}}\right) - 2\left[\frac{\sqrt{3}}{2}\cos\left(\frac{\delta k_x a}{2\sqrt{3}}\right) + \left(\frac{1}{2}\right)\sin\left(\frac{\delta k_x a}{2\sqrt{3}}\right)\right]\left[\left(\frac{1}{2}\right)\cos\left(\frac{\delta k_y a}{2}\right) + \left(\frac{\sqrt{3}}{2}\right)\sin\left(\frac{\delta k_y a}{2}\right)\right] \right]
 \end{aligned}$$

Alternative K' (II)

- We propose $K = k_x = \frac{2\pi}{a\sqrt{3}} + \delta k_x; k_y = -\frac{2\pi}{3a} + \delta k_y$

$$\begin{aligned}
 &= -V_{pp\pi} \left[\sin\left(\frac{2\pi}{3}\right) \cos\left(\frac{\delta k_x a}{\sqrt{3}}\right) + \cos\left(\frac{2\pi}{3}\right) \sin\left(\frac{\delta k_x a}{\sqrt{3}}\right) - 2 \left[\sin\left(\frac{\pi}{3}\right) \cos\left(\frac{\delta k_x a}{2\sqrt{3}}\right) + \cos\left(\frac{\pi}{3}\right) \sin\left(\frac{\delta k_x a}{2\sqrt{3}}\right) \right] \left[\cos\left(-\frac{\pi}{3}\right) \cos\left(\frac{\delta k_y a}{2}\right) - \sin\left(-\frac{\pi}{3}\right) \sin\left(\frac{\delta k_y a}{2}\right) \right] \right] \\
 &= -V_{pp\pi} \left[\frac{\sqrt{3}}{2} \cos\left(\frac{\delta k_x a}{\sqrt{3}}\right) + \left(-\frac{1}{2}\right) \sin\left(\frac{\delta k_x a}{\sqrt{3}}\right) - 2 \left[\frac{\sqrt{3}}{2} \cos\left(\frac{\delta k_x a}{2\sqrt{3}}\right) + \left(\frac{1}{2}\right) \sin\left(\frac{\delta k_x a}{2\sqrt{3}}\right) \right] \left[\left(\frac{1}{2}\right) \cos\left(\frac{\delta k_y a}{2}\right) + \left(\frac{\sqrt{3}}{2}\right) \sin\left(\frac{\delta k_y a}{2}\right) \right] \right] \\
 &= -V_{pp\pi} \left[\frac{\sqrt{3}}{2} + \left(-\frac{1}{2}\right) \left(\frac{\delta k_x a}{\sqrt{3}}\right) - \left[\sqrt{3} + \left(\frac{\delta k_x a}{2\sqrt{3}}\right) \right] \left[\left(\frac{1}{2}\right) + \left(\frac{\sqrt{3}}{2}\right) \left(\frac{\delta k_y a}{2}\right) \right] \right] \\
 &= -V_{pp\pi} \left[\frac{\sqrt{3}}{2} + \left(-\frac{1}{2}\right) \left(\frac{\delta k_x a}{\sqrt{3}}\right) - \left\{ \left(\frac{\sqrt{3}}{2}\right) + \left(\frac{3}{2}\right) \left(\frac{\delta k_y a}{2}\right) \right\} + \left(\frac{\delta k_x a}{4\sqrt{3}}\right) \right] = -V_{pp\pi} \left[\frac{\sqrt{3}}{2} - \left(\frac{\delta k_x a}{2\sqrt{3}}\right) - \left(\frac{\sqrt{3}}{2}\right) - \left(\frac{3}{2}\right) \left(\frac{\delta k_y a}{2}\right) - \left(\frac{\delta k_x a}{4\sqrt{3}}\right) \right] \\
 &= -V_{pp\pi} \left[-\left(\frac{3\delta k_x a}{4\sqrt{3}}\right) - \left(\frac{3}{2}\right) \left(\frac{\delta k_y a}{2}\right) \right] = -\frac{3V_{pp\pi} a}{4} \left[-\frac{\delta k_x}{\sqrt{3}} - \delta k_y \right]
 \end{aligned}$$

Axis rotation K' point

- We have

$$H_x(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x - \frac{\delta k_y}{\sqrt{3}} \right]$$

$$H_y(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[-\frac{\delta k_x}{\sqrt{3}} - \delta k_y \right]$$

- Define

$$\begin{pmatrix} q'_x \\ q'_y \end{pmatrix} = \begin{pmatrix} \frac{\sqrt{3}}{2} & -\frac{1}{2} \\ \frac{1}{2} & \frac{\sqrt{3}}{2} \end{pmatrix} \begin{pmatrix} \delta k_x \\ \delta k_y \end{pmatrix}$$

$$H_x(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[\delta k_x - \frac{\delta k_y}{\sqrt{3}} \right] = -\frac{3V_{pp\pi} a}{4} \left(\frac{2q'_x}{\sqrt{3}} \right) = -\frac{\sqrt{3}V_{pp\pi} a}{2} q'_x$$

$$H_y(\mathbf{k}) = -\frac{3V_{pp\pi} a}{4} \left[-\frac{\delta k_x}{\sqrt{3}} - \delta k_y \right] = -\frac{3V_{pp\pi} a}{4} \left(-\frac{2q'_y}{\sqrt{3}} \right) = -\frac{\sqrt{3}V_{pp\pi} a}{2} (-q'_y)$$

Near Dirac point hamiltonian

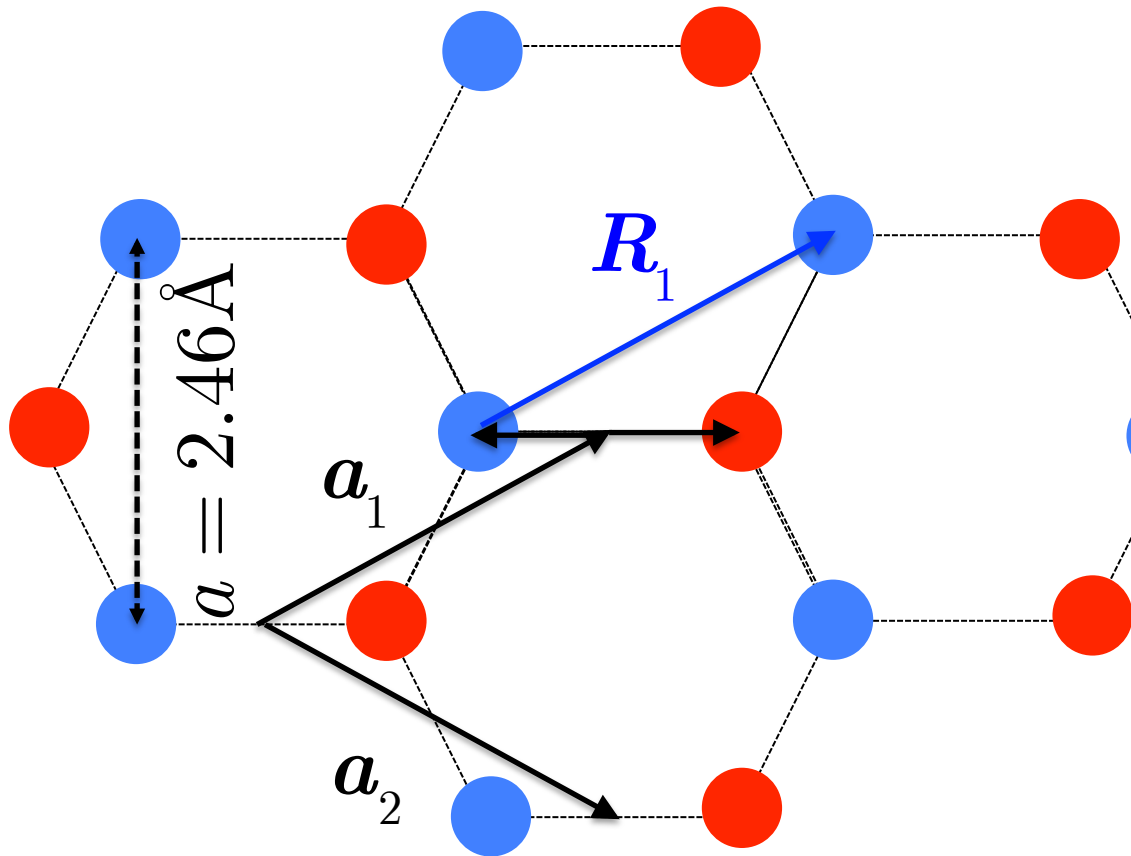
- For the K point:

$$H_K(\mathbf{q}) = \begin{pmatrix} \varepsilon_{pA} & 0 \\ 0 & \varepsilon_{pB} \end{pmatrix} - \frac{\sqrt{3}V_{pp\pi}a}{2}q_x \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} - \frac{\sqrt{3}V_{pp\pi}a}{2}q_y \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
$$= \varepsilon_p I - \frac{\sqrt{3}V_{pp\pi}a}{2}(\sigma_x q_x + \sigma_y q_y) + \sigma_z \Delta_p$$

- For the K' point:

$$H_{K'}(\mathbf{q}) = \begin{pmatrix} \varepsilon_{pA} & 0 \\ 0 & \varepsilon_{pB} \end{pmatrix} - \frac{\sqrt{3}V_{pp\pi}a}{2}q_x \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \frac{\sqrt{3}V_{pp\pi}a}{2}q_y \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
$$= \varepsilon_p I - \frac{\sqrt{3}V_{pp\pi}a}{2}(\sigma_x q_x - \sigma_y q_y) + \sigma_z \Delta_p$$

Graphene second neighbors



$$R_1 = \frac{a}{2} (\sqrt{3}\hat{i} + \hat{j})$$

$$R_2 = a\hat{j}$$

$$R_3 = \frac{a}{2} (-\sqrt{3}\hat{i} + \hat{j})$$

$$R_4 = \frac{a}{2} (-\sqrt{3}\hat{i} - \hat{j}) = -R_1$$

$$R_5 = -a\hat{j} = -R_2$$

$$R_6 = \frac{a}{2} (\sqrt{3}\hat{i} - \hat{j}) = -R_3$$

Second neighbor H

- Second neighbors belong to the same sublattice, so this will only add diagonal elements:

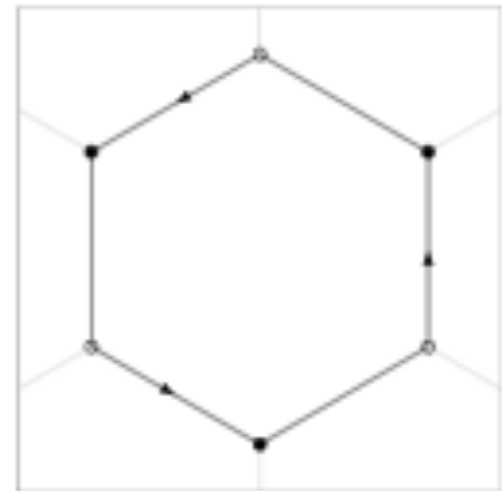
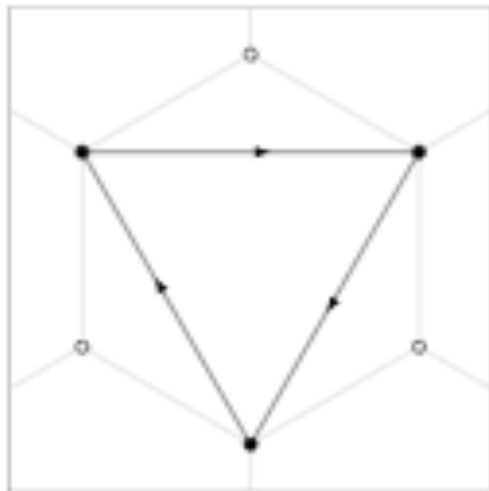
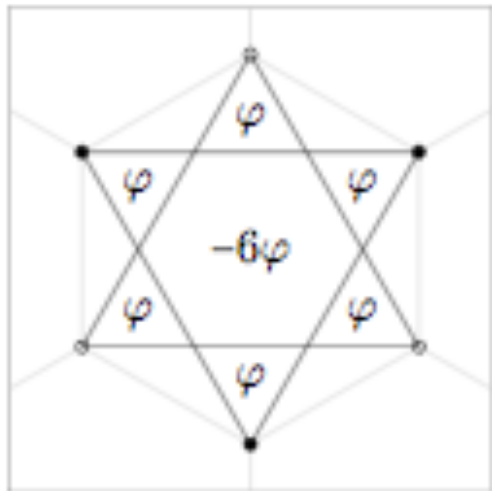
$$\begin{aligned} H_{jj} &= \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot(\mathbf{R})} \int d\mathbf{r} \phi_{j,0}^*(\mathbf{r}) H \phi_{j,\mathbf{R}}(\mathbf{r}) \\ &= \sum_k e^{i\mathbf{k}\cdot\mathbf{R}_k} \int d\mathbf{r} \phi_{j,0}^*(\mathbf{r}) H \phi_{j,\mathbf{nn}_k}(\mathbf{r}) \end{aligned}$$

Matrix element

- Let us now assume:

$$\int d\mathbf{r} \phi_{A,0}^*(\mathbf{r}) H \phi_{A,\mathbf{R}_k}(\mathbf{r}) = \begin{cases} V_2 e^{i\alpha} & \text{for } i = 1, 3, 4 \\ V_2 e^{-i\alpha} & \text{for } i = 2, 4, 6 \end{cases}$$
$$\int d\mathbf{r} \phi_{B,0}^*(\mathbf{r}) H \phi_{B,\mathbf{R}_k}(\mathbf{r}) = \begin{cases} V_2 e^{-i\alpha} & \text{for } i = 1, 3, 4 \\ V_2 e^{+i\alpha} & \text{for } i = 2, 4, 6 \end{cases}$$

How do we get this?



arXiv:1310.0255v2

H(k)

- Then

$$\begin{aligned} H_{AA} &= V_2 \left[e^{i(\mathbf{k} \cdot \mathbf{R}_1 + \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_3 + \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_5 + \alpha)} \right] \\ &\quad + V_2 \left[e^{i(\mathbf{k} \cdot \mathbf{R}_2 - \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_4 - \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_6 - \alpha)} \right] \\ &= 2V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 + \alpha) \right] \end{aligned}$$

$$\begin{aligned} H_{BB} &= V_2 \left[e^{i(\mathbf{k} \cdot \mathbf{R}_1 - \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_3 - \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_5 - \alpha)} \right] \\ &\quad + V_2 \left[e^{i(\mathbf{k} \cdot \mathbf{R}_2 + \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_4 + \alpha)} + e^{i(\mathbf{k} \cdot \mathbf{R}_6 + \alpha)} \right] \\ &= 2V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 - \alpha) \right] \end{aligned}$$

New elements

- This second neighbor interaction contributes an additional

$$\begin{aligned} H_0(\mathbf{k}) &= \frac{H_{AA}(\mathbf{k}) + H_{BB}(\mathbf{k})}{2} = \\ &= V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 + \alpha) \right] + \\ &V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 - \alpha) \right] \\ &= 2V_2 \cos \alpha \left[\cos(\mathbf{k} \cdot \mathbf{R}_1) + \cos(\mathbf{k} \cdot \mathbf{R}_3) + \cos(\mathbf{k} \cdot \mathbf{R}_5) \right] \end{aligned}$$

$$\begin{aligned} H_z(\mathbf{k}) &= \frac{H_{AA}(\mathbf{k}) - H_{BB}(\mathbf{k})}{2} = \\ &= V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 + \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 + \alpha) \right] \\ &- V_2 \left[\cos(\mathbf{k} \cdot \mathbf{R}_1 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_3 - \alpha) + \cos(\mathbf{k} \cdot \mathbf{R}_5 - \alpha) \right] \\ &= 2V_2 \sin \alpha \left[\sin(\mathbf{k} \cdot \mathbf{R}_1) + \sin(\mathbf{k} \cdot \mathbf{R}_3) + \sin(\mathbf{k} \cdot \mathbf{R}_5) \right] \end{aligned}$$

At K (I)

$$\mathbf{R}_1 = \frac{a}{2}(\sqrt{3}\hat{\mathbf{i}} + \hat{\mathbf{j}})$$

$$\mathbf{R}_3 = \frac{a}{2}(-\sqrt{3}\hat{\mathbf{i}} + \hat{\mathbf{j}})$$

$$\mathbf{R}_5 = -a\hat{\mathbf{j}} = -\mathbf{R}_2$$

• We have $K = k_x = \frac{2\pi}{a\sqrt{3}}; k_y = \frac{2\pi}{3a}$

$$H_0(\mathbf{k}) = 2V_2 \cos \alpha \left[\cos(\mathbf{k} \cdot \mathbf{R}_1) + \cos(\mathbf{k} \cdot \mathbf{R}_3) + \cos(\mathbf{k} \cdot \mathbf{R}_5) \right]$$

$$= 2V_2 \cos \alpha \left[\cos \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{\sqrt{3}a}{2} \right) + \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \cos \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{-\sqrt{3}a}{2} \right) + \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \cos \left[+ \left(\frac{2\pi}{3a} \right) (-a) \right] \right]$$

$$= 2V_2 \cos \alpha \left[\cos \left[\pi + \frac{\pi}{3} \right] + \cos \left[-\pi + \frac{\pi}{3} \right] + \cos \left[-\frac{2\pi}{3} \right] \right]$$

$$= 2V_2 \cos \alpha \left[\cos \left[\frac{4\pi}{3} \right] + \cos \left[-\frac{2\pi}{3} \right] + \cos \left[-\frac{2\pi}{3} \right] \right]$$

$$= 2V_2 \cos \alpha \left[\cos \left[\frac{4\pi}{3} \right] + \cos \left[-\frac{2\pi}{3} \right] + \cos \left[-\frac{2\pi}{3} \right] \right]$$

$$-3V_2 \cos \alpha$$

At K (II)

• We have
$$K = k_x = \frac{2\pi}{a\sqrt{3}}; k_y = \frac{2\pi}{3a}$$

$$\begin{aligned} H_z(\mathbf{k}) &= 2V_2 \sin \alpha \left[\sin(\mathbf{k} \cdot \mathbf{R}_1) + \sin(\mathbf{k} \cdot \mathbf{R}_3) + \sin(\mathbf{k} \cdot \mathbf{R}_5) \right] \\ &= 2V_2 \sin \alpha \left[\sin \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{\sqrt{3}a}{2} \right) + \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \sin \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{-\sqrt{3}a}{2} \right) + \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \sin \left[+ \left(\frac{2\pi}{3a} \right) (-a) \right] \right] \\ &= 2V_2 \sin \alpha \left[\sin \left[\pi + \frac{\pi}{3} \right] + \sin \left[-\pi + \frac{\pi}{3} \right] + \sin \left[-\frac{2\pi}{3} \right] \right] \\ &= 2V_2 \cos \alpha \left[\sin \left[\frac{4\pi}{3} \right] + \sin \left[-\frac{2\pi}{3} \right] + \sin \left[-\frac{2\pi}{3} \right] \right] \\ &= -3\sqrt{3}V_2 \sin \alpha \end{aligned}$$

At K' (II)

• We have
$$K' = k_x = \frac{2\pi}{a\sqrt{3}}; k_y = -\frac{2\pi}{3a}$$

$$\begin{aligned} H_z(\mathbf{k}) &= 2V_2 \sin \alpha \left[\sin(\mathbf{k} \cdot \mathbf{R}_1) + \sin(\mathbf{k} \cdot \mathbf{R}_3) + \sin(\mathbf{k} \cdot \mathbf{R}_5) \right] \\ &= 2V_2 \sin \alpha \left[\sin \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{\sqrt{3}a}{2} \right) - \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \sin \left[\left(\frac{2\pi}{a\sqrt{3}} \right) \left(\frac{-\sqrt{3}a}{2} \right) - \left(\frac{2\pi}{3a} \right) \frac{a}{2} \right] + \sin \left[+ \left(-\frac{2\pi}{3a} \right) (-a) \right] \right] \\ &= 2V_2 \sin \alpha \left[\sin \left[\pi - \frac{\pi}{3} \right] + \sin \left[-\pi - \frac{\pi}{3} \right] + \sin \left[\frac{2\pi}{3} \right] \right] \\ &= 2V_2 \cos \alpha \left[\sin \left[\frac{2\pi}{3} \right] + \sin \left[-\frac{4\pi}{3} \right] + \sin \left[\frac{2\pi}{3} \right] \right] \\ &= 3\sqrt{3}V_2 \sin \alpha \end{aligned}$$

Hamiltonian matrix

- The whole hamiltonian is

$$H = \begin{pmatrix} H_0 + H_z & H_x - iH_y \\ H_x + iH_y & H_0 - H_z \end{pmatrix} \quad \begin{aligned} H_z &= \Delta_p - 3\sqrt{3}V_2 \sin \alpha & (K) \\ H_z &= \Delta_p + 3\sqrt{3}V_2 \sin \alpha & (K') \end{aligned}$$

- with

$$\varepsilon_{\pm} = H_0 \pm \sqrt{H_x^2 + H_y^2 + H_z^2} = H_0 \pm H$$

- and

$$u_- = \begin{pmatrix} -\sin \frac{\theta}{2} \\ e^{i\varphi} \cos \frac{\theta}{2} \end{pmatrix} \quad \begin{aligned} H_x &= H \sin \theta \cos \varphi \\ H_y &= H \sin \theta \sin \varphi \\ H_z &= H \cos \theta \end{aligned}$$