

La clase pasada vimos:

Pozos cuánticos dobles y parabólicos

Aproximación de masa efectiva: impureza hidrogenoide

En esta clase veremos:

Nanoestructuras: cables y puntos cuánticos

Aproximación de masa efectiva en heteroestructuras

Dopaje remoto

Pozos cuánticos dobles y parabólicos

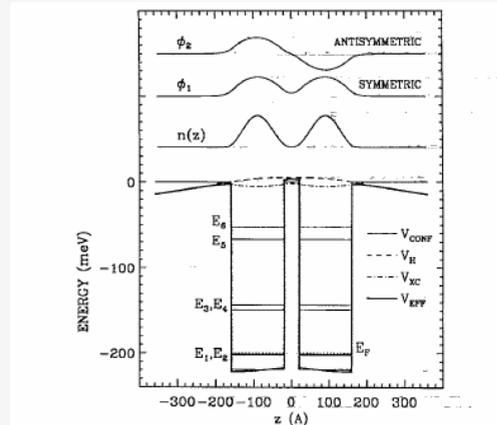
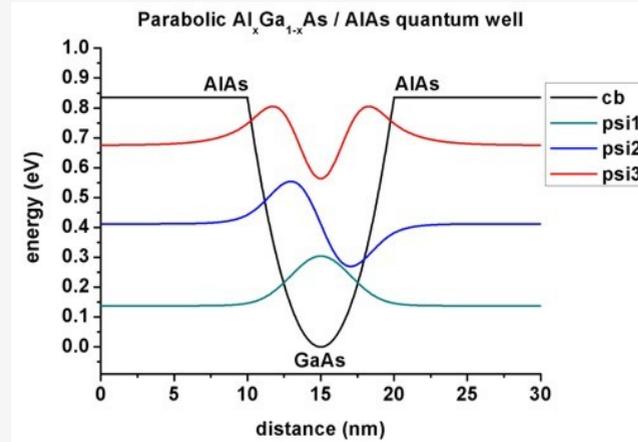
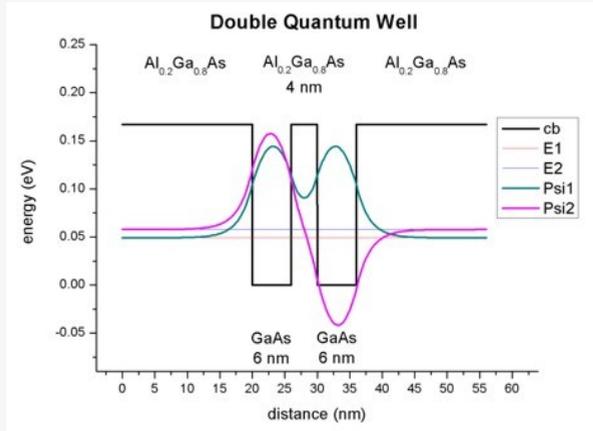


FIG. 1. The typical double quantum well structure ($N_S = 10^{11} \text{ cm}^{-2}$) with the self-consistent local-density-approximation calculated energy levels and wave functions for the symmetric-antisymmetric levels. The bare (V_{conf}), Hartree (V_H), and exchange-correlation (V_{xc}) contributions to the effective potential (V_{EFF}) are shown. (The SAS gap $\Delta_{\text{SAS}} = |E_2 - E_1|$.)

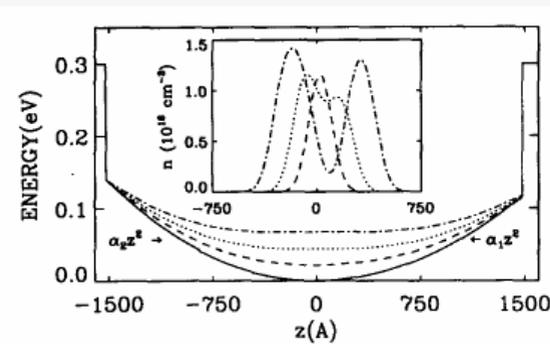
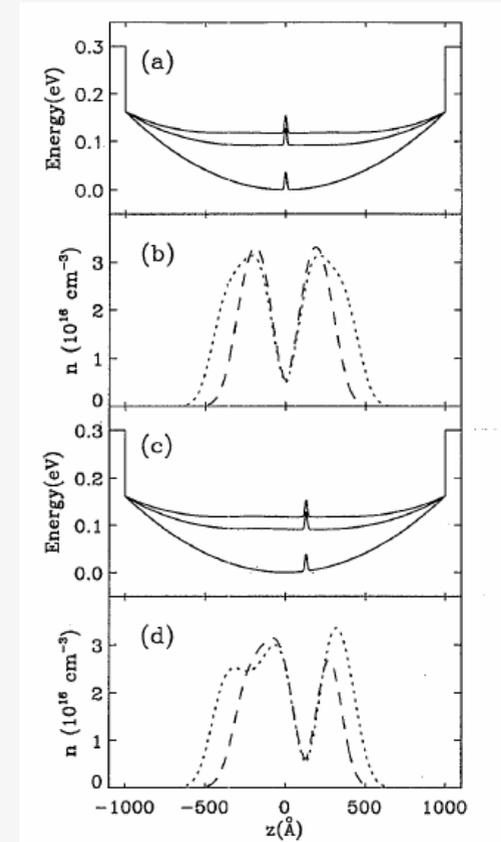
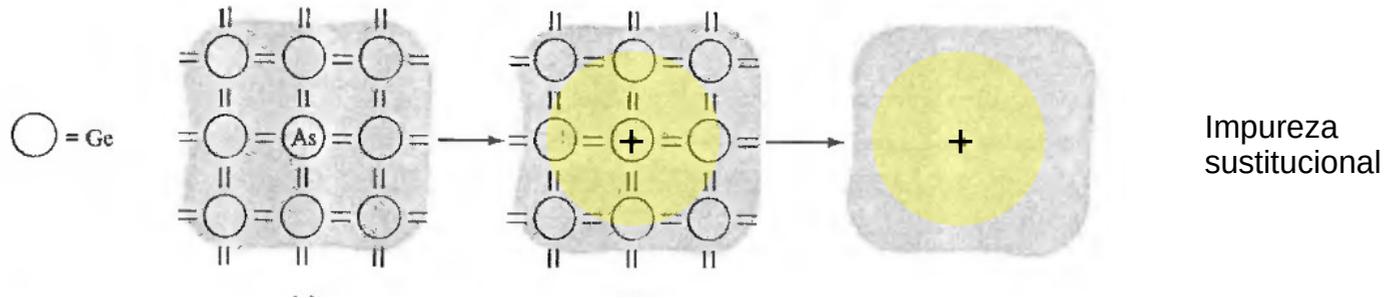


Fig. 1. Model bare potential of the asymmetric parabolic well (solid line), and calculated self-consistent potentials with an in-plane magnetic field of $B = 5.8 \text{ T}$, for $N_s = 2.4$ (dash line), 4.7 (dotted line), and $7.5 \times 10^{10} \text{ cm}^{-2}$ (dash-dot line); inset: corresponding calculated self-consistent densities.



REPASO

Aproximación de masa efectiva: impureza hidrogenoide



Ecuación de Schrödinger para el electrón donado al material:

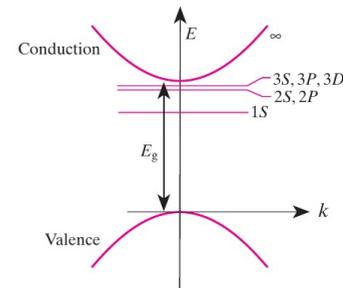
$$[\hat{H}_{\text{per}} + V_{\text{imp}}(\mathbf{R})]\psi(\mathbf{R}) = E\psi(\mathbf{R})$$

Aproximaciones: (1) una sola banda (2) Integrar k cerca del mínimo

$$\psi(x) \approx \phi_{n0}(x) \int_{-\pi/a}^{\pi/a} \tilde{\chi}(k) \exp(ikx) \frac{dk}{2\pi} = \phi_{n0}(x) \chi(x)$$

Y la función envolvente satisface:

$$\left[\varepsilon_n \left(-i \frac{d}{dx} \right) + V_{\text{imp}}(x) \right] \chi(x) = E \chi(x).$$



Confinamiento adicional:

Nanoestructuras - cables y puntos cuánticos

Confinamiento adicional: Nanoestructuras - cables y puntos cuánticos

Estamos usando:

THE PHYSICS OF
LOW-DIMENSIONAL
SEMICONDUCTORS

AN INTRODUCTION

JOHN H. DAVIES

Glasgow University

Otro libro especializado:

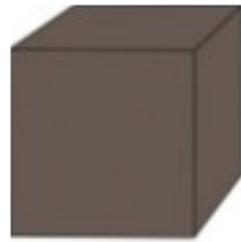
**QUANTUM WELLS, WIRES
AND DOTS**

Theoretical and Computational
Physics of Semiconductor
Nanostructures

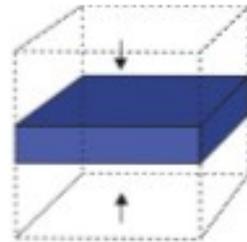
Paul Harrison

The University of Leeds, UK

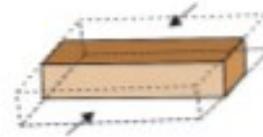
Confinamiento adicional: Nanoestructuras - cables y puntos cuánticos



Bulk material



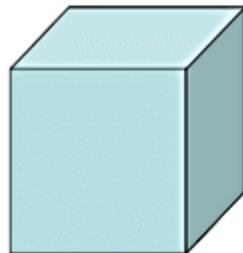
Quantum well



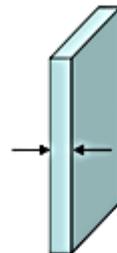
Quantum wire



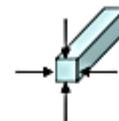
Quantum dot



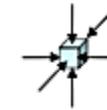
Bulk
3D



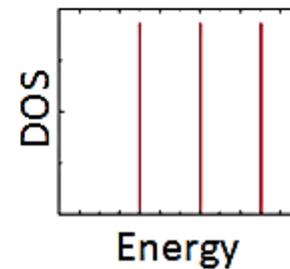
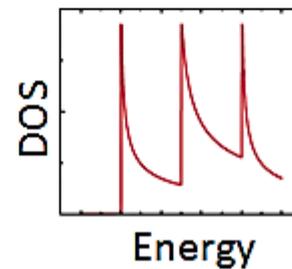
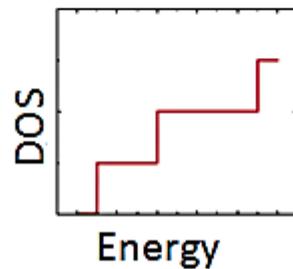
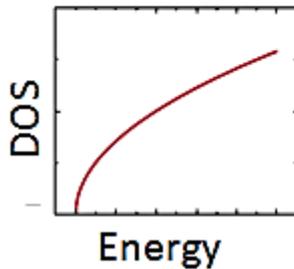
Quantum well
2D



Quantum wire
1D



Quantum dot
0D



Cables cuánticos

Quantum Wires

To make the transition from a 2D electron gas (quantum well) to a 1D electron gas (quantum wire), the electrons should be confined in two directions and only 1 degree of freedom remain. The **x direction** remains the **only one** for **free-electron propagation**.

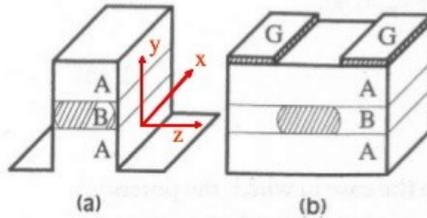


Figure 3.11. Two of the simplest examples of structures providing electron confinement in two dimensions: case (a) uses an etching technique, while case (b) is based on the split-gate technique.

Nanowhiskers o nanorods

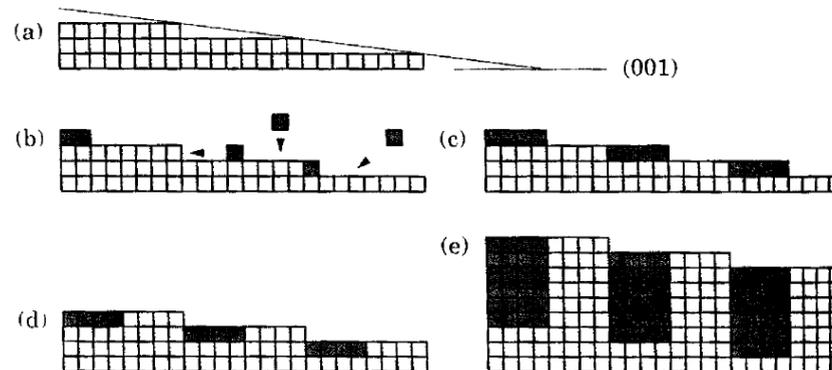
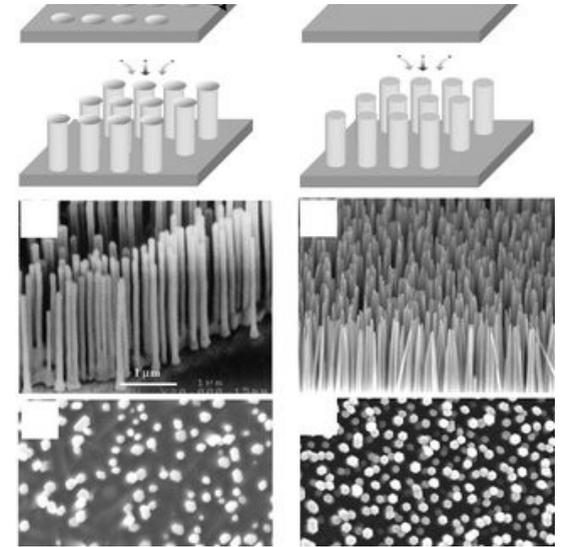
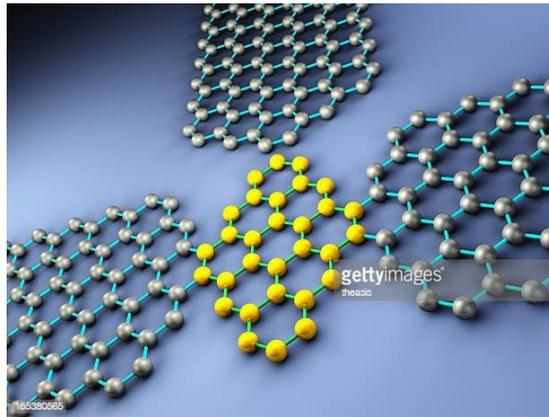
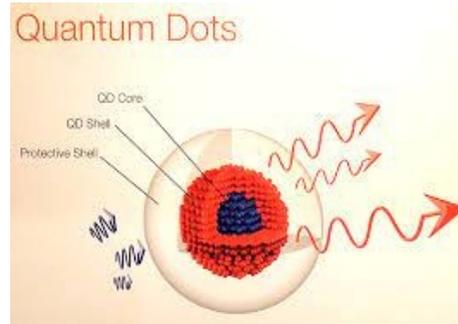
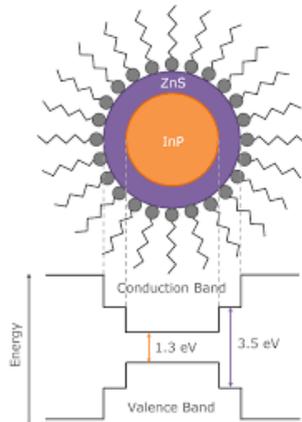


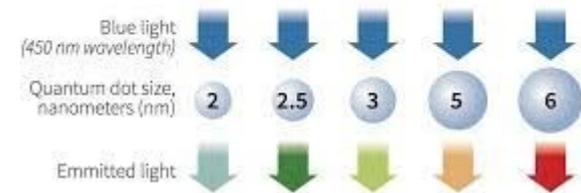
FIGURE 3.18. Growth of wires on a vicinal substrate. (a) Substrate before growth, showing steps and terraces on the vicinal (off-axis) surface. (b) Flux of AIAs applied, with motion of atoms over surface to steps where they stick. (c) Completion of a half-monolayer of AIAs. (d) Growth of a further half-monolayer of GaAs. (e) Growth of many layers, showing development of superlattice.

Puntos cuánticos (de químicos)

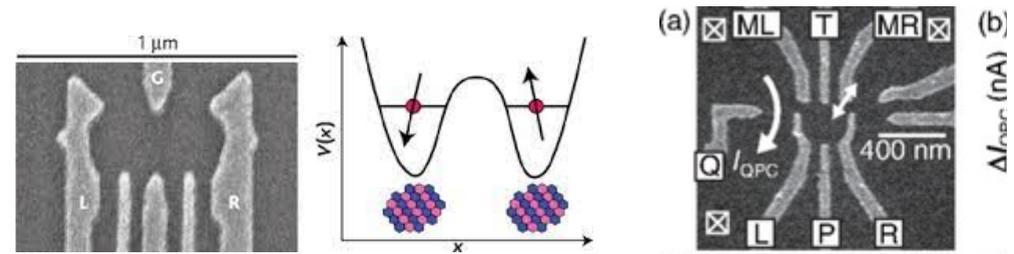
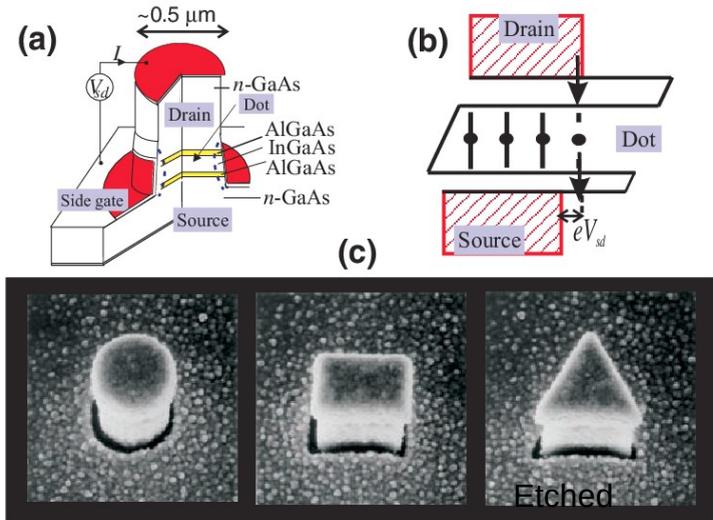


What are quantum dots

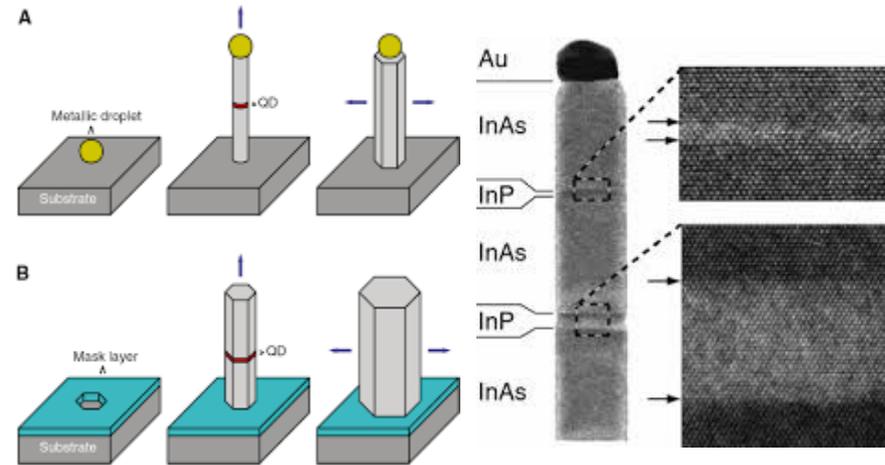
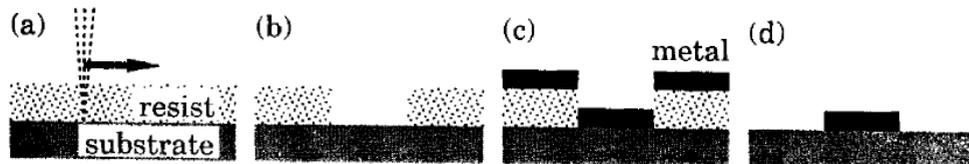
Tiny man-made crystals that have the ability to convert a spectrum of light into different colours



Puntos cuánticos (de físicos)



Con voltage gates

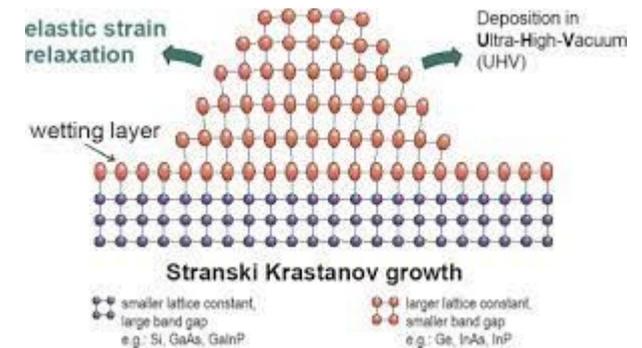
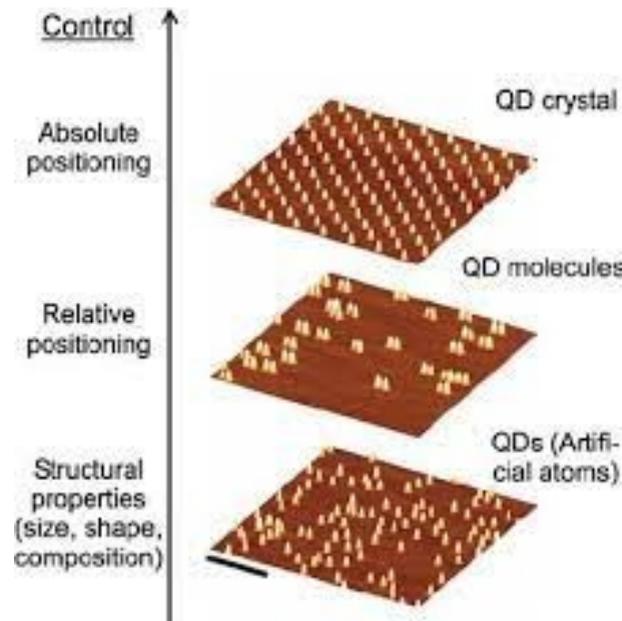
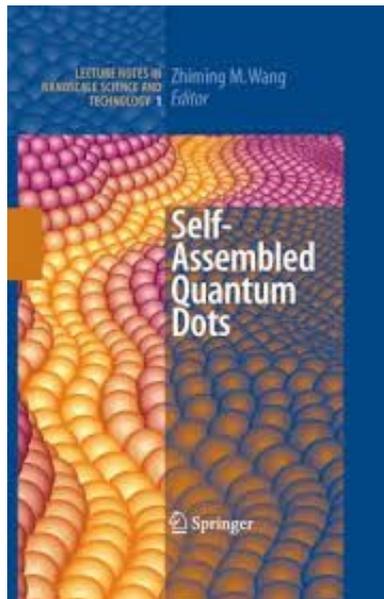
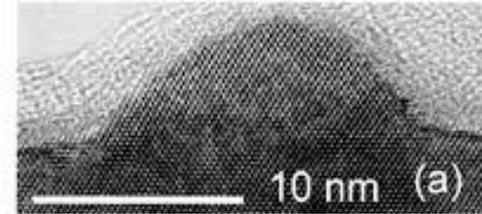
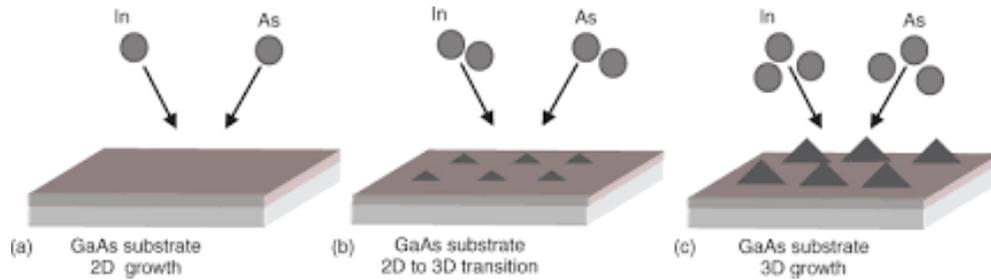


Dentro de nanorods

FIGURE 3.16. Production of a metal gate by electron-beam lithography and lift-off. (a) Exposure of resist by electron beam, (b) development of resist, (c) deposition of metal, and (d) lift-off to leave metal gate.

Puntos cuánticos (de físicos)

Self-assembled quantum dots



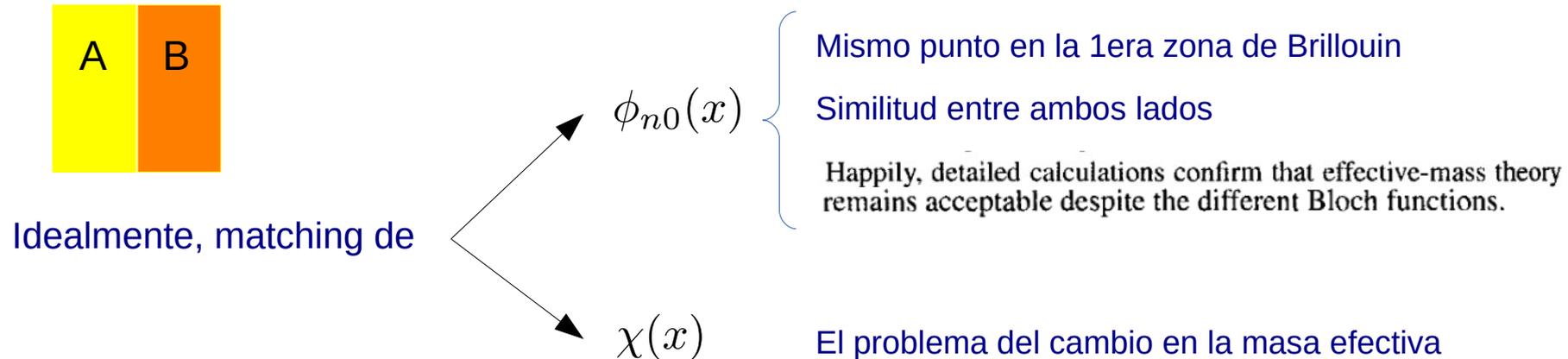
Aproximación de masa efectiva en heteroestructuras

Aproximación de masa efectiva en hetero-estructuras

Habíamos visto que en presencia de un potencial externo (sumado al del cristal) la función de onda de electrón se podría factorizar:

$$\psi(x) \approx \phi_{n0}(x) \int_{-\pi/a}^{\pi/a} \tilde{\chi}(k) \exp(ikx) \frac{dk}{2\pi} = \phi_{n0}(x) \chi(x).$$

¿Qué pasa en la interfaz entre dos materiales diferentes en una heteroestructura?



Aproximación de masa efectiva en hetero-estructuras

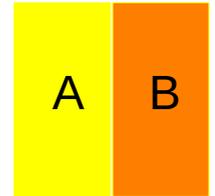
Envelope function $\chi(x)$

$$\text{----- } E_c^B$$

$$\text{----- } E_c^A$$

$$\Delta E_c = E_c^B - E_c^A$$

$$\left\{ \begin{array}{l} \left(E_c^A - \frac{\hbar^2}{2m_A m_0} \frac{d^2}{dz^2} \right) \chi(z) = E \chi(z) \\ \left(E_c^B - \frac{\hbar^2}{2m_B m_0} \frac{d^2}{dz^2} \right) \chi(z) = E \chi(z) \end{array} \right.$$

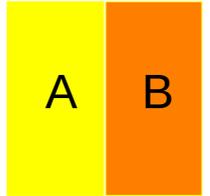


$$\chi(0_A) = \chi(0_B), \quad \frac{d\chi(z)}{dz} \Big|_{z=0_A} = \frac{d\chi(z)}{dz} \Big|_{z=0_B}$$

$$\chi(0_A) = \chi(0_B), \quad \frac{1}{m_A} \frac{d\chi(z)}{dz} \Big|_{z=0_A} = \frac{1}{m_B} \frac{d\chi(z)}{dz} \Big|_{z=0_B}$$

Aproximación de masa efectiva en hetero-estructuras

Envelope function $\chi(x)$



$$-\frac{\hbar^2}{2m_0 m(z)} \frac{d^2 \chi}{dz^2} + V(z)\chi(z) = E\chi(z). \quad (3.20)$$

This is not Hermitian (or of Sturm–Liouville form) if $m(z)$ varies, and many of the

Versión más correcta:

$$-\frac{\hbar^2}{2m_0} \frac{d}{dz} \left[\frac{1}{m(z)} \frac{d\chi}{dz} \right] + V(z)\chi(z) = E\chi(z)$$

Aproximación de masa efectiva en hetero-estructuras

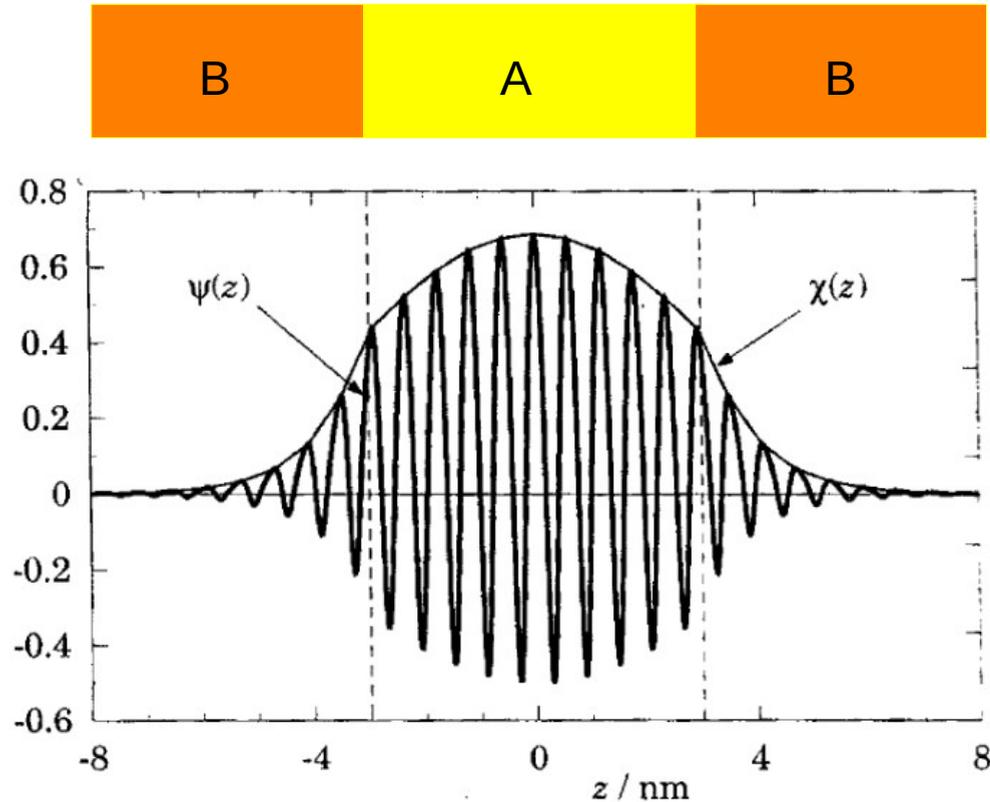


FIGURE 3.22. Wave function for the lowest state in a 6 nm quantum well in a heterostructure, including the Bloch functions. The thin curve is an approximate envelope function joining the peaks of the full wave function. [Redrawn from Burt (1994).]

Modulation doping - dopaje remoto

Modulation doping - dopaje remoto

WIKIPEDIA

Modulation doping

Modulation doping is a technique for fabricating semiconductors such that the free charge carriers are spatially separated from the donors. Because this eliminates scattering from the donors, modulation-doped semiconductors have very high carrier mobilities.

History

Modulation doping was conceived in Bell Labs in 1977 following a conversation between Horst Störmer and Ray Dingle,^[1] and implemented shortly afterwards by Arthur Gossard. In 1977, Störmer and Dan Tsui used a modulation-doped wafer to discover the fractional quantum Hall effect.

Modulation doping - dopaje remoto

Dopaje: agregar átomos que ceden electrones o huecos al material.

Problema: los iones que quedan son centros de scattering que reducen la mobility

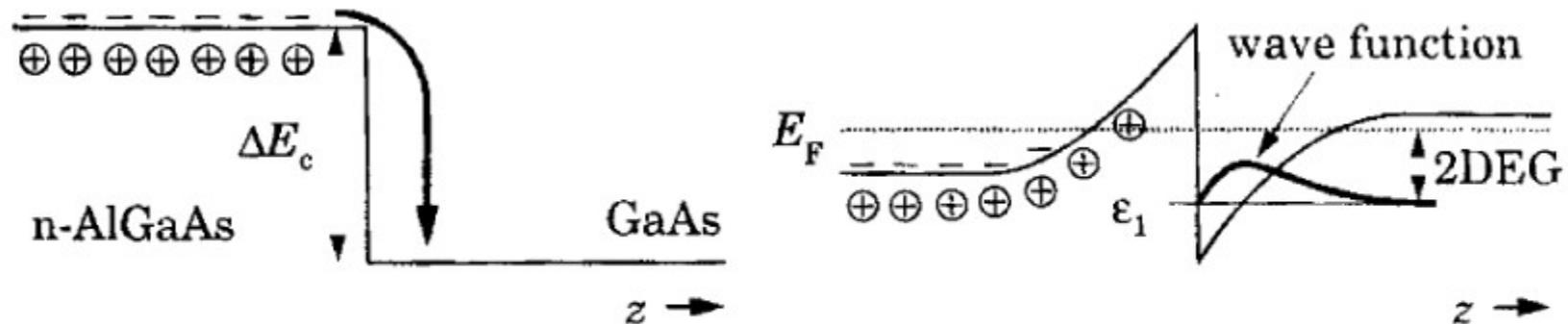


FIGURE 3.9. Conduction band around a heterojunction between n-AlGaAs and undoped GaAs, showing how electrons are separated from their donors to form a two-dimensional electron gas.

Los electrones migran de una región a otra y evitan las colisiones con los iones (cargados atractivamente)

Resumen de la clase 7

Cable y puntos cuánticos

Aproximación de masa efectiva en heteroestructuras

Dopaje remoto